COSMOLOGICAL FUZZY DARK MATTER SIMULATIONS

Differences to cold dark matter

Simon May simon.may@mpa-garching.mpg.de

Max-Planck-Institut für Astrophysik

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Outline

Introduction

- Motivation and theoretical background
- The fuzzy dark matter equations
- Impact of fuzzy dark matter on structure formation
 - #1: Initial conditions
 - #2: Dynamics

Numerical methods and challenges

Simulations and results

- Dark matter power spectrum
- Halo mass function
- Dark matter halo profiles

Summary

What is "fuzzy dark matter"?

- F(C)DM, BECDM, ULDM, ELBDM, ψDM, quantum-wave DM, (ultra-light) axion(-like) DM (ULA, ALP)...
- New extremely light scalar particle ($m \approx 10^{-22} \, \mathrm{eV!}$)
- Non-thermal production mechanism (thus not ultra-hot)
- Aggregations of bosons can form a **Bose–Einstein condensate**
- Quantum effects counteract gravity at small scales (uncertainty principle), erase structure
- Tiny mass
 - \Rightarrow large de Broglie wavelength ($\lambda \sim 1/m$)
 - \Rightarrow macroscopic quantum effects on kpc scales

Motivation for fuzzy dark matter

Particle physics perspective:

- Original concept strong CP problem: Why doesn't QCD violate CP symmetry?
- Solved by Peccei–Quinn U(1) symmetry and (pseudo-)scalar field (axion!) Peccei and Quinn (1977)!
- Fuzzy dark matter is **not** the QCD axion, but axion-like particles are a common feature of early-universe theories

Astrophysics perspective:

- Small-scale challenges (cusp-core, missing satellites, ...)
- ightarrow Ultra-light scalars: WIMP alternative, could improve this



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- ightarrow Ultra-light scalars: WIMP alternative, could improve this

wrong LSS
$$\leftarrow | \qquad | \qquad | \qquad | \qquad \rightarrow CDM$$
-like
10⁻²³ eV 10⁻²² eV 10⁻²¹ eV
No sign of (WIMP) CDM

The fuzzy dark matter equations

Schrödinger-Poisson system

Add a scalar field to the Einstein–Hilbert action of general relativity

$$S = \frac{1}{\hbar c^2} \int \mathrm{d}^4 x \sqrt{-g} \Biggl(\frac{1}{2} g^{\mu\nu} (\partial_\mu \phi) (\partial_\nu \phi) - \frac{1}{2} \frac{m^2 c^2}{\hbar^2} \phi^2 - \frac{\lambda}{\hbar^2 c^2} \phi^4 \Biggr)$$

Non-relativistic limit yields the Schrödinger equation

$$i\hbar \Bigl(\partial_t \psi + \frac{3}{2} H \psi \Bigr) = -\frac{\hbar^2}{2m} \nabla^2 \psi + m \Phi \psi$$

Mean field approximation: interpretation as single macroscopic wave function of BE condensate with density ρ = m|ψ|²
 "FDM equations": non-linear Schrödinger–Poisson system

$$\begin{split} i\hbar\partial_t\psi_{\rm c} &= -\frac{\hbar^2}{2ma^2}\nabla_{\rm c}^2\psi_{\rm c} + \frac{m}{a}\Phi_{\rm c}\psi_{\rm c} \\ \nabla_{\rm c}^2\Phi_{\rm c} &= 4\pi Gm(|\psi_{\rm c}|^2 - \langle|\psi_{\rm c}|^2\rangle) \end{split}$$

Only a single scale, determined by $\frac{\hbar}{m}$ (\rightarrow wavelength)

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Impact of fuzzy dark matter – #1: Initial conditions



(Only CDM initial conditions for now ightarrow comparability)

Seeds of structure suppressed below quantum wavelength

Impact of fuzzy dark matter – #2: Dynamics



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Impact of fuzzy dark matter – #2: Dynamics



Quantum fluctuations, interference patterns

Small-scale (sub-)structure suppressed

Mocz, Fialkov, et al. (2020)

0.5 Mp

Using pseudo-spectral methods to simulate fuzzy dark matter Symmetrized split-step Fourier $i\hbar\partial_t\psi = -\frac{\hbar^2}{2ma^2}\nabla^2\psi + \frac{m}{a}\Phi\psi$

- Symmetrized split-step Fourier method ("kick-drift-kick")
- Algorithm:

$$\begin{array}{ll} \psi \leftarrow e^{-i\frac{m}{\hbar}\frac{\Delta t}{2}\Phi}\psi & \mbox{kick} \\ \psi \leftarrow \mathrm{IFFT}\Big(e^{-i\frac{\hbar}{m}\frac{\Delta t}{2}k^2}\,\mathrm{FFT}(\psi)\Big) & \mbox{drift} \\ \Phi \leftarrow \mathrm{IFFT}\Big(-\frac{1}{k^2}\,\mathrm{FFT}\big(4\pi Gm(|\psi|^2 - \langle |\psi|^2\rangle)\big)\Big) & \mbox{update potential} \\ \psi \leftarrow e^{-i\frac{m}{\hbar}\frac{\Delta t}{2}\Phi}\psi & \mbox{kick} \end{array}$$

Choice of time step: $\Delta t < \min\left(\frac{4}{9\pi}\frac{m}{\hbar}a^2\Delta x^2\right)$, $2\pi\frac{\hbar}{m}a\frac{1}{|\Phi_{\max}|}$

- "Exact" solution
- Automatic conservation of mass
- Can adapt existing particle–mesh (PM) code
- Simple implementation

 $\nabla^2 \Phi = 4\pi Gm(|\psi|^2 - \langle |\psi|^2 \rangle)$

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Correspondence of CDM and FDM initial conditions

- Cold dark matter (CDM): Collisionless fluid, described by a phase space distribution function $f(\vec{x}, \vec{v})$
- Can construct a wave function ψ from a distribution function f:

$$\psi(\vec{x}) \sim \sum_{\vec{v}} \sqrt{f(\vec{x},\vec{v})} e^{i^m/\!\!/\hbar \vec{x}\cdot\vec{v} + R_{\vec{v}}}$$

For "cold"/single-stream distribution function:

$$\psi = \sqrt{\frac{\rho}{m}} e^{i\alpha}$$
$$\vec{v} = \frac{\hbar}{m} \nabla \alpha$$

 Grid discretization implies a maximum velocity which can be represented

Mocz, Lancaster, et al. (2018), Mocz, Fialkov, et al. (2020)

$$v < \frac{\hbar}{m} \frac{\pi}{\Delta x}$$

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Why is it hard to simulate fuzzy dark matter? Computational challenges

- Both large scales and small (kpc-scale) de Broglie wavelength must be resolved for correct evolution
- High velocities require high resolution even in low-density regions
 Velocity criterion: v < ^ħ/_m π/_{Δx}
- Time step criterion: $\Delta t \sim \Delta x^2$ (seems to be approach-independent)
- (Tooling: Hydrodynamics codes are designed for *N*-body simulations)

Schive, Chiueh, and Broadhurst (2014)





First results: Dark matter power spectrum $10 M_{pc/h}$ box simulations



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First results: Dark matter power spectrum 10 Mpc/h box simulations



Agreement on large scales
Delayed structure formation
Suppression on small scales

First results: Halo mass function 10 Mpc/h box simulations



First results: Dark matter halo profiles (stacked) 10 Mpc/h box simulations



- Fuzzy dark matter halos form central cores with flat density profiles
- Adaptive resolution better suited to show inner halos

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Summary

Main problems for fuzzy dark matter simulations:

- 1. Time integration $\Delta t \sim \Delta x^2$
- 2. Rapid oscillations even in low-density regions
- 3. Large dynamic range: "large"-scale structure simulations must still resolve de Broglie wavelength, limited to $\lesssim 10~{\rm Mpc}/{\rm h}$ box
- 4. "New" field without decades of experience or refined codes/methods as for CDM

Features of FDM dynamics compared to CDM:

- 1. (Modified initial power spectrum, small scales suppressed)
- 2. Suppression of structure below de Broglie wavelength (\approx kpc) \rightarrow Heisenberg uncertainty principle
- 3. Formation of halo cores
- 4. Fluctuating \approx kpc quantum interference patterns

References

Peccei, R. D. and Helen R. Quinn (June 1977). "CP conservation in the presence of pseudoparticles". In: *Phys. Rev. Lett.* 38, pp. 1440–1443. DOI: 10.1103/PhysRevLett.38.1440.

- Mocz, Philip, Anastasia Fialkov, et al. (Apr. 2020). "Galaxy formation with BECDM II. Cosmic filaments and first galaxies". In: *MNRAS* 494.2, pp. 2027–2044. DOI: 10.1093/mnras/staa738. arXiv: 1911.05746 [astro-ph.CO].
- Mocz, Philip, Lachlan Lancaster, et al. (Apr. 2018). "Schrödinger-Poisson-Vlasov-Poisson correspondence". In: Phys. Rev. D 97.8, 083519, p. 083519. DOI: 10.1103/PhysRevD.97.083519. arXiv: 1801.03507 [astro-ph.C0].
- Schive, Hsi-Yu, Tzihong Chiueh, and Tom Broadhurst (July 2014). "Cosmic structure as the quantum interference of a coherent dark wave". In: *Nature Physics* 10, pp. 496–499. DOI: 10.1038/nphys2996. arXiv: 1406.6586 [astro-ph.GA].