Hunting for Dark Matter in Anisotropies of Gamma-ray Sky: Theory and First Observational Results from Fermi-LAT

Eiichiro Komatsu (Texas Cosmology Center, Univ. of Texas at Austin) MPA Seminar, September 14, 2011

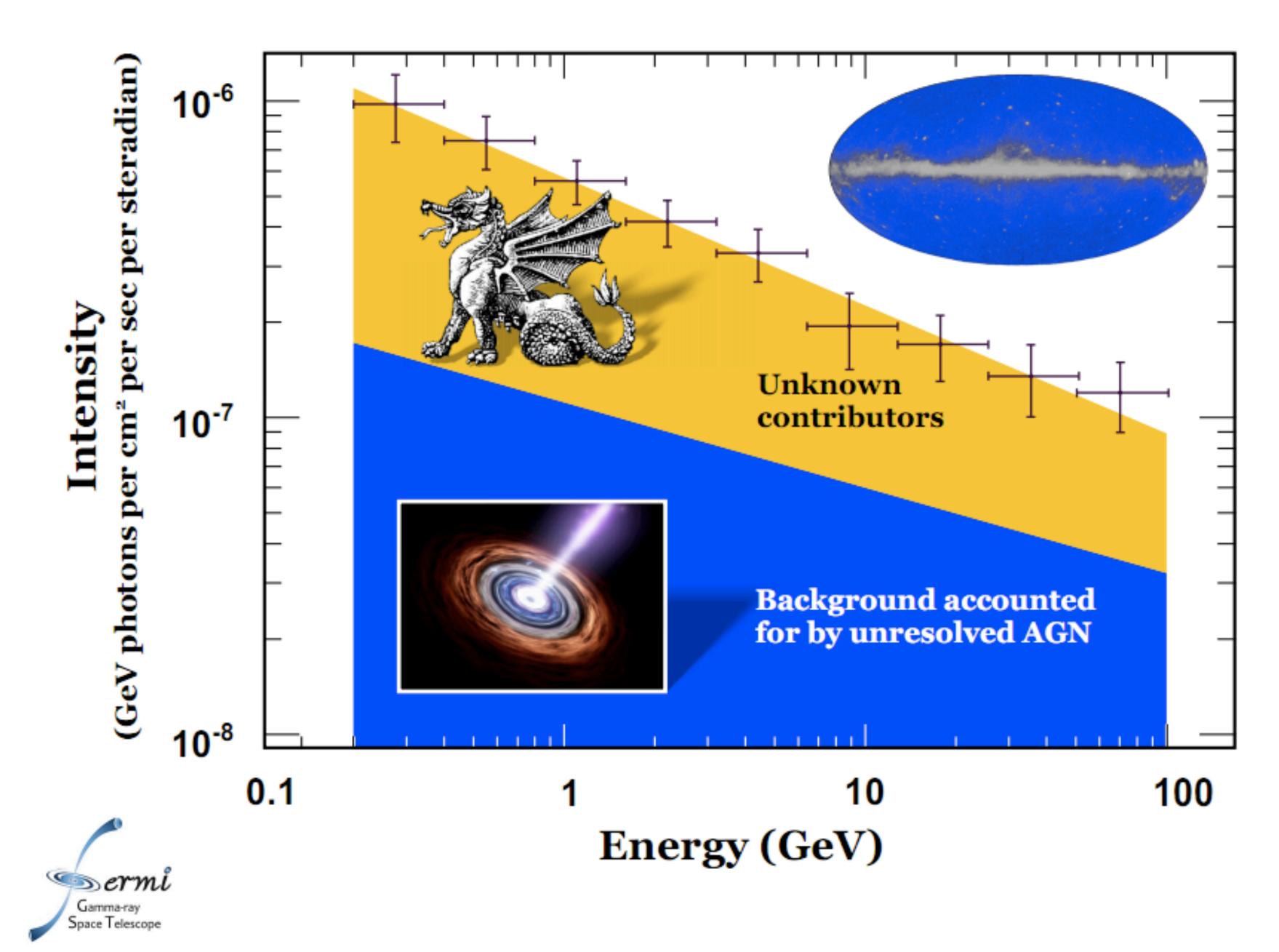
Motivation

• How can we see photons from annihilation/decay of dark matter particles?

Intriguing Observations

- In gamma-ray energies (E>0.1GeV), the origin of 80% of the diffuse emission (after removing the known Galactic emission) is unknown!
 - 20% coming from blazars (Fermi-LAT collaboration)
- In soft gamma-ray energies (E=I-I0MeV), the origin of >90% of the diffuse emission is unknown!
 - <10% coming from supernovae (Ahn, Komatsu and Hoeflich 2005)

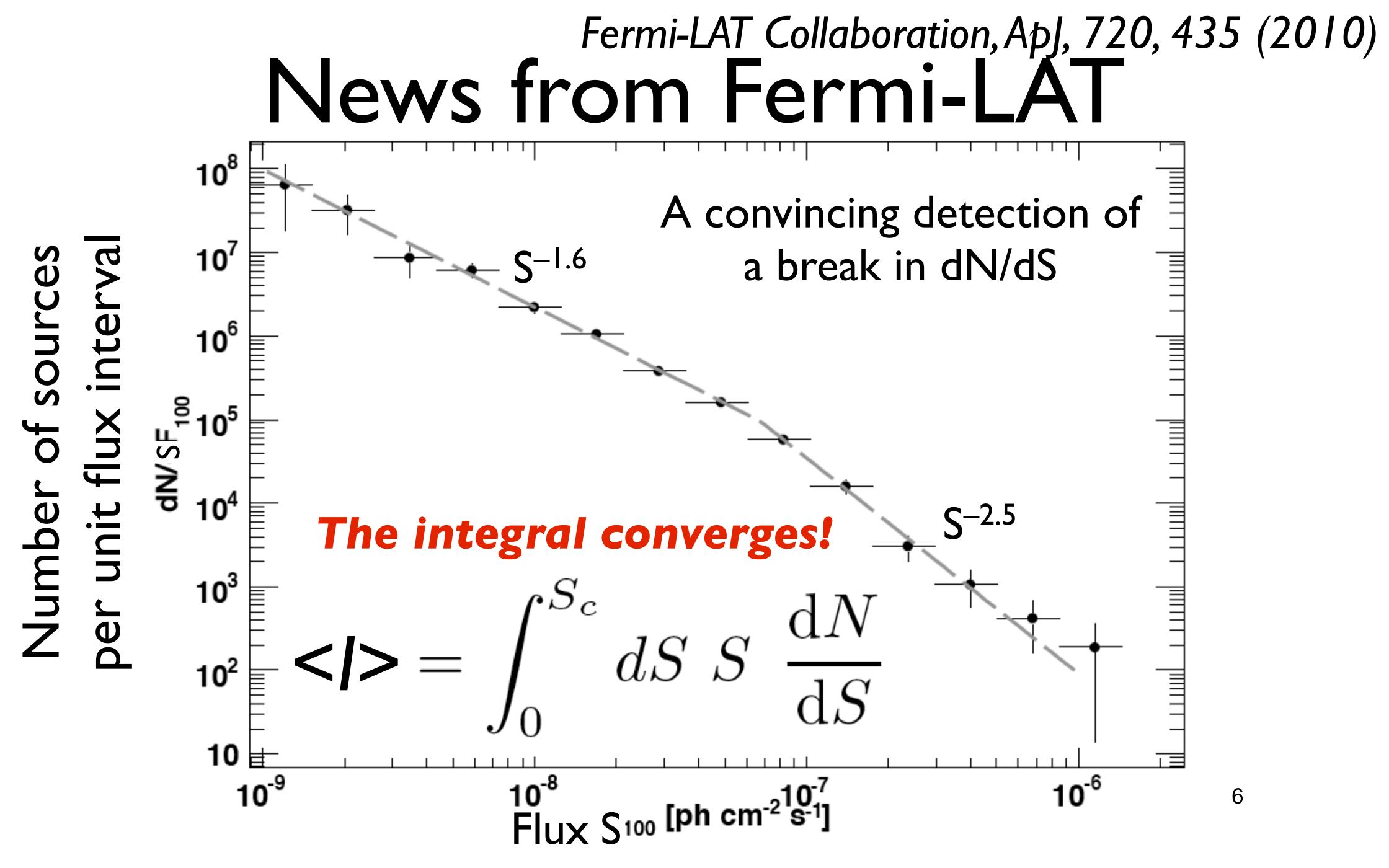
Fermi LAT Extragalactic Gamma-ray Background



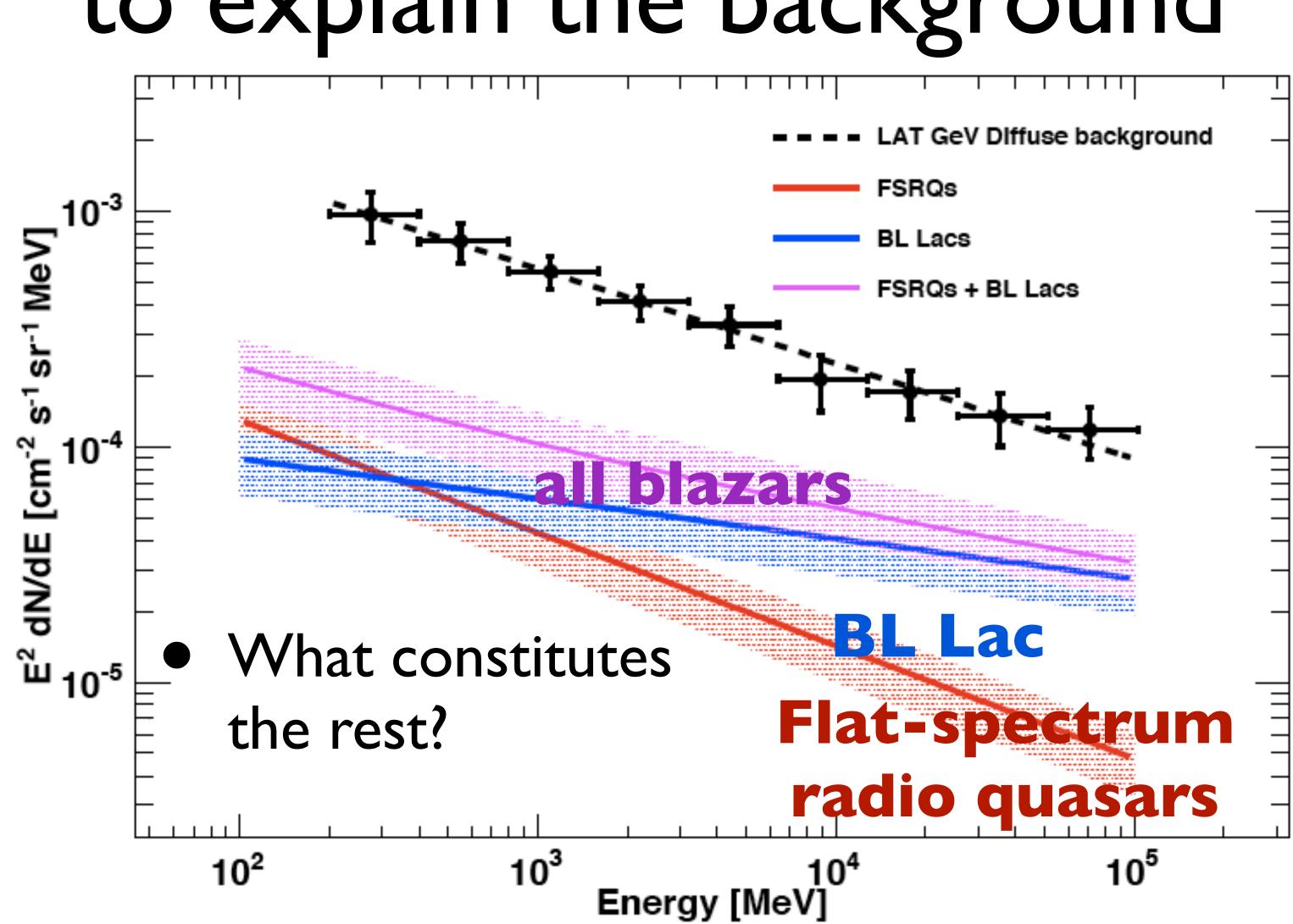


Blazars

- Blazars = A population of AGNs whose relativistic jets are directed towards us.
 Inverse Compton scattering of relativistic particles in jets off photons -> gamma-rays, detected up to TeV
- How many are there? (They are rare.)
 EGRET found ~70 blazars (out of ~100 associated sources) over the full sky
 Fermi-LAT found ~570 blazars (out of ~820 associated sources) over the full sky (LAT 1FGL catalog) 5



Fermi-LAT Collaboration, ApJ, 720, 435 (2010) Unresolved blazars are not enough to explain the background



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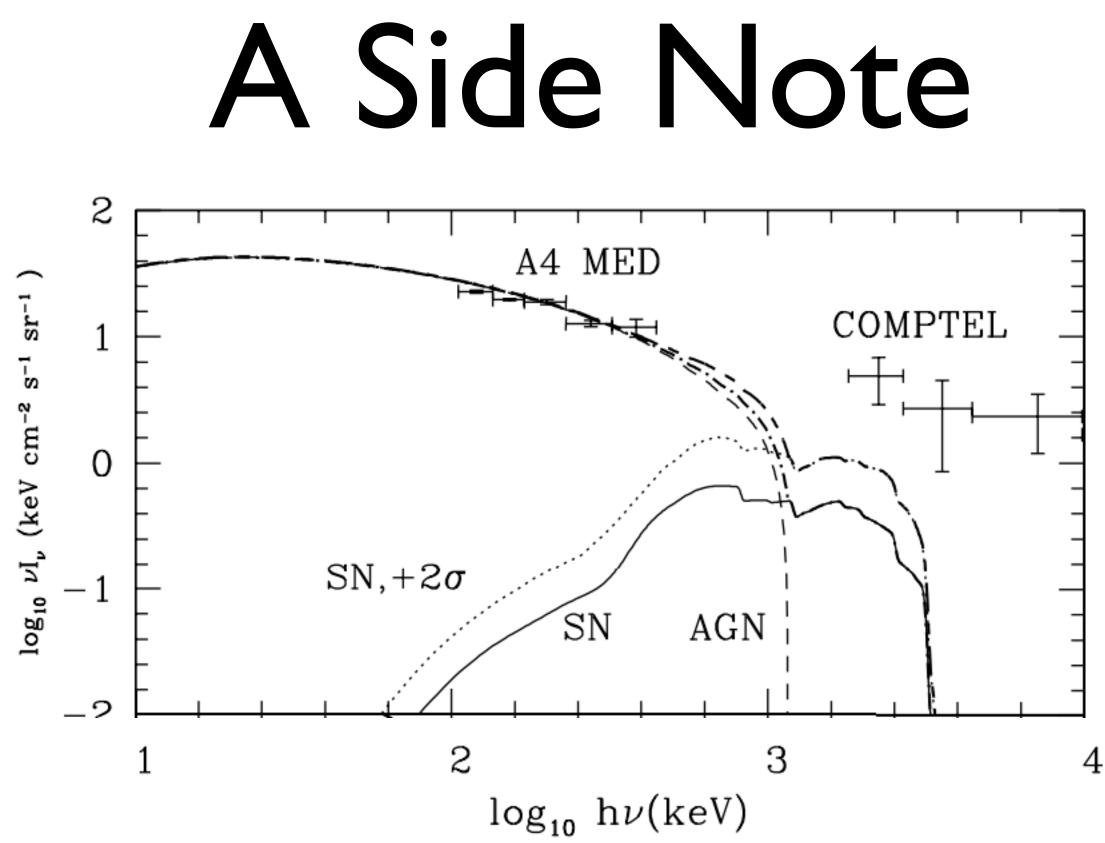
Origin of Diffuse Gamma-ray Background?

- Where do they come from?
 - Star-forming galaxies?
 - Pulsars?
 - Clusters of galaxies?

Origin of Diffuse Gamma-ray Background?

- Where do they come from?
 - Star-forming galaxies?
 - Pulsars?
 - Clusters of galaxies?

or... perhaps... some of them might come from...
Dark matter?



 It was thought that Type Ia supernovae would account that the measured supernova rate is too small for that! The origin of the MeV background is unknown. 10

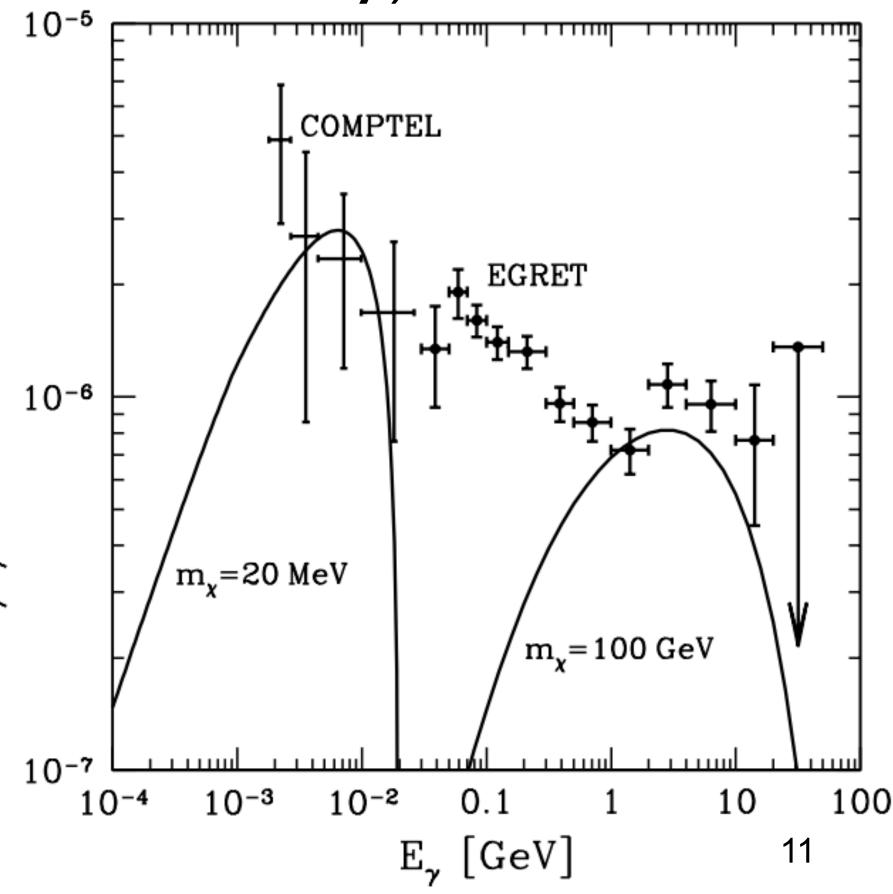
Ahn, Komatsu & Hoeflich (2005)

for most of the MeV gamma-ray background. It turns out

Conventional Method

• Use the energy spectrum of the mean intensity (the number of photons averaged over the sky), and look for spectral features.

However, dark matter is not the only source of gamma-ray photons. How can we distinguish between dark matter signatures and astrophysical sources?



$$A General$$
$$E_{\gamma}I_{\gamma}(\hat{\boldsymbol{n}}, E_{\gamma}) = \frac{c}{4\pi} \int dz \frac{P_{\gamma}([1])}{H(1)}$$

• All we need: $P_{y} =$ "volume emissivity" = energy radiated per unit volume, time, and energy.

E.g., for supernovae:

 $P_{\nu}(\nu, z) = (1 + z)^{3} \text{SNR}_{\text{Ia}}(z) \bar{E}_{\nu}$

Formula $\frac{[+z]E_{\gamma}, z, \hat{n}r}{(z)(1+z)^4}e^{-\tau([1+z]E_{\gamma}, z)}$

A General Formula

$$E_{\gamma}I_{\gamma}(\hat{n}, E_{\gamma}) = \frac{c}{4\pi} \int dz \frac{P_{\gamma}([1+z]E_{\gamma}, z, \hat{n}r)}{H(z)(1+z)^4} e^{-\tau([1+z]E_{\gamma}, z)}$$

• All we need: $P_{y} =$ "volume emissivity" = energy radiated per unit volume, time, and energy.

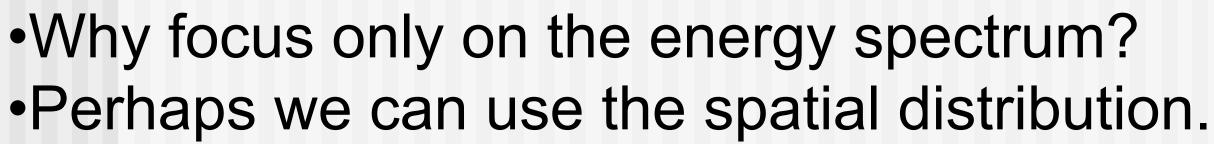
E.g., for dark matter annihilation:

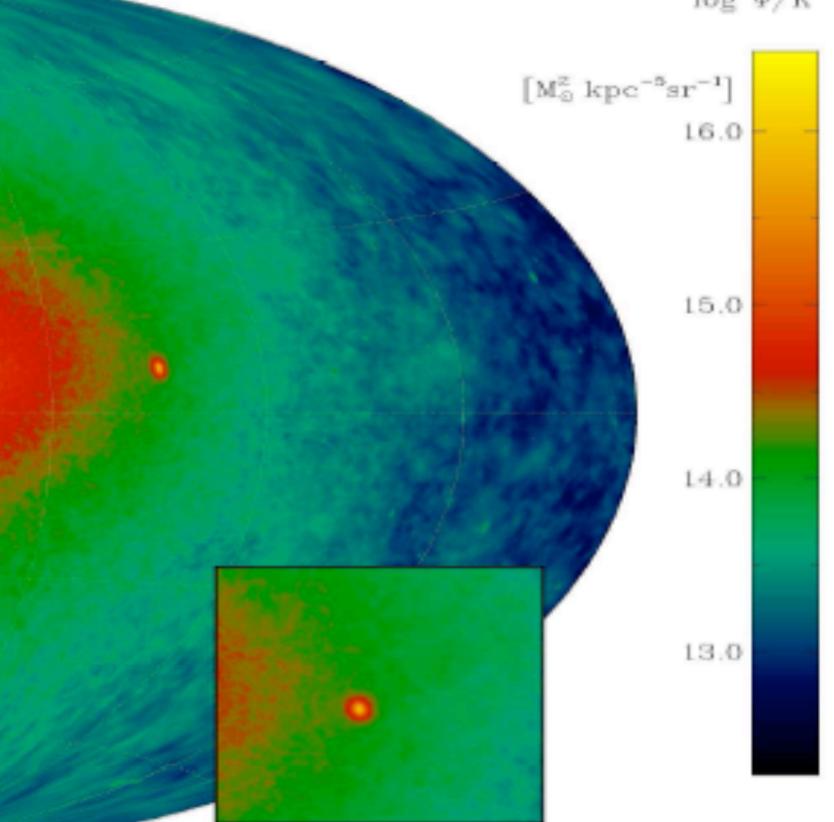
 $P_{\gamma}(E_{\gamma}, z, \hat{\boldsymbol{n}}r) = E$

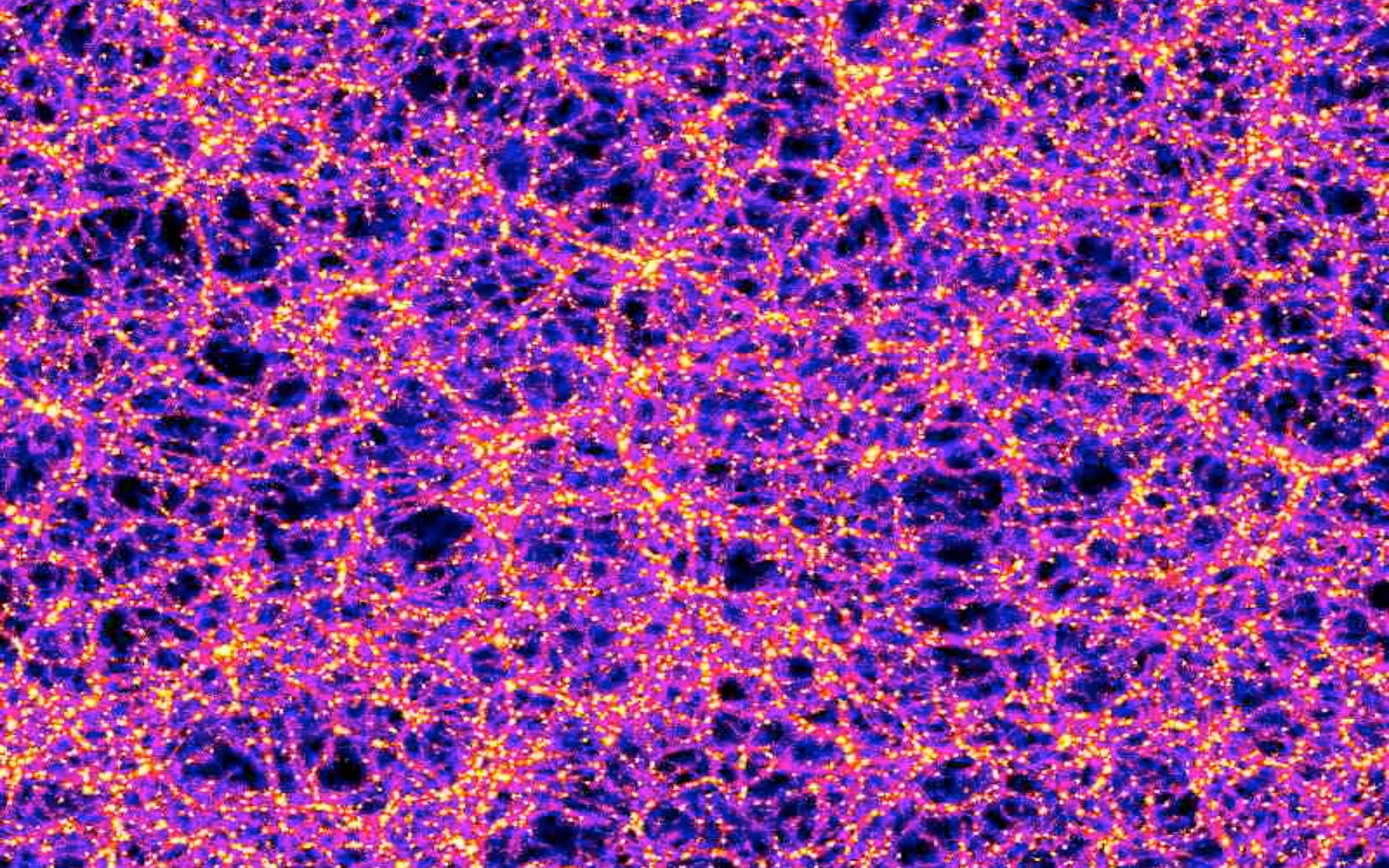
$$\int_{\gamma} \frac{dN_{\gamma}}{dE_{\gamma}} \frac{\langle \sigma v \rangle}{2} \left[\frac{\rho_{\chi}(z, \hat{n} r)}{m_{\chi}} \right]^2$$



Diemand, Khlen & Madau, ApJ, 657, 262 (2007) **Annihilation Signals from Milky Way** log Φ/K







And, not just Milky Way!

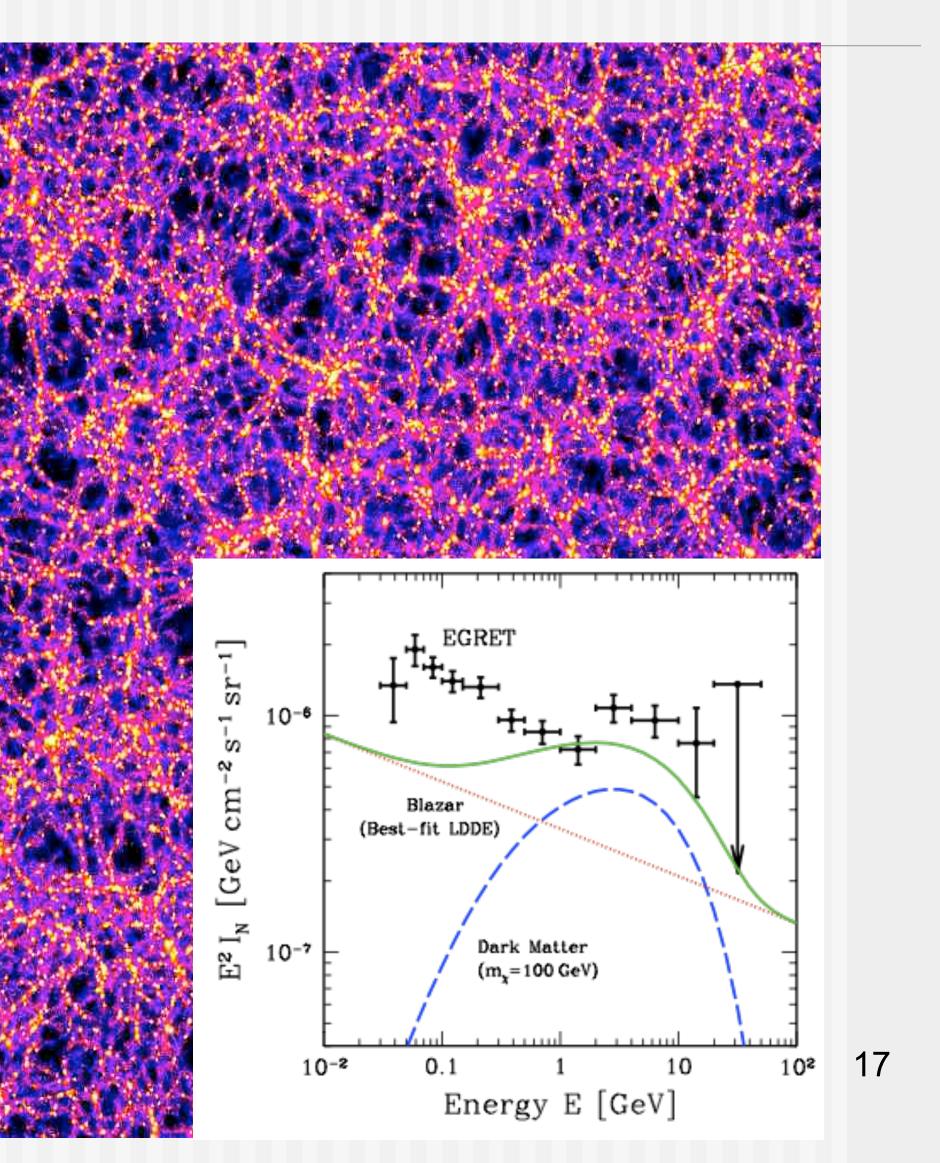
Dark matter particles are annihilating (or decaying) everywhere in the Universe!

• Why just focus on Milky Way?

While we cannot resolve individual dark matter halos, the collective signals can be detected in the diffuse gamma-ray background.

How can we detect such signatures unambiguously?





Ando & EK (2006); Ando, EK, Narumoto & Totani (2007)

Gamma-ray Anisotropy

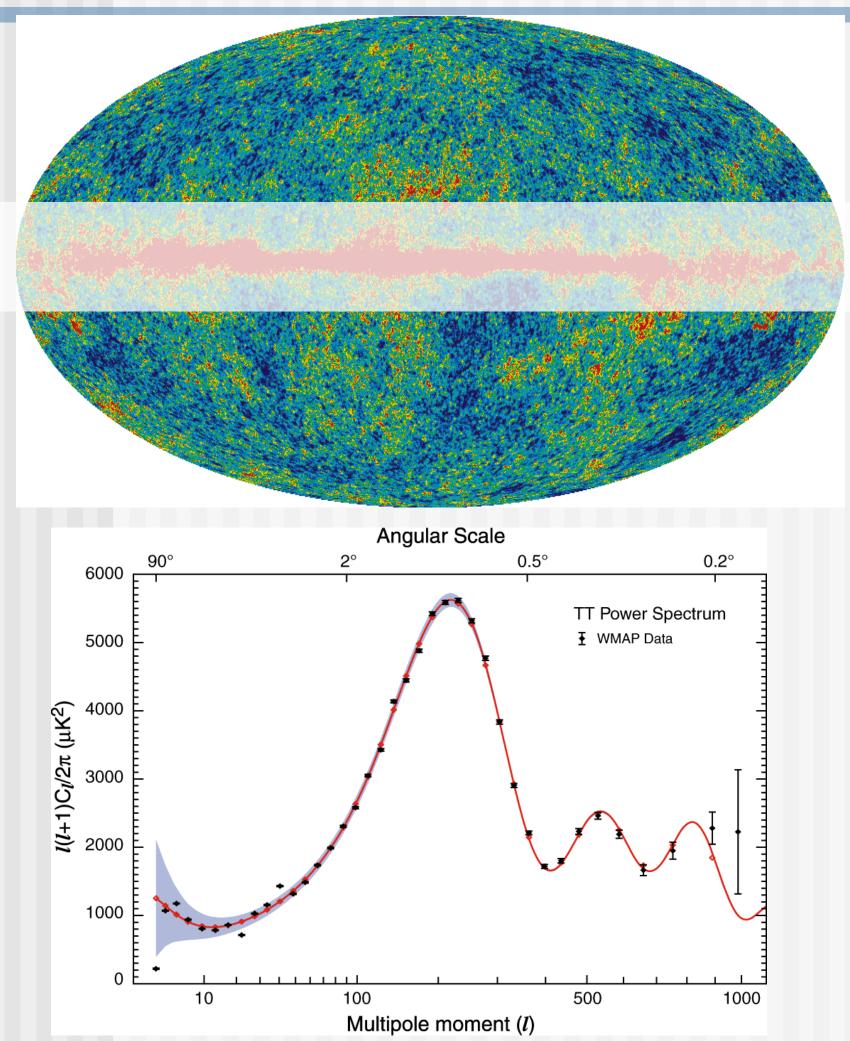
Dark matter halos trace the large-scale structure

Therefore, the gamma-ray background must be anisotropic. If dark matter particles annihilate or decay, anisotropy **must** be there.

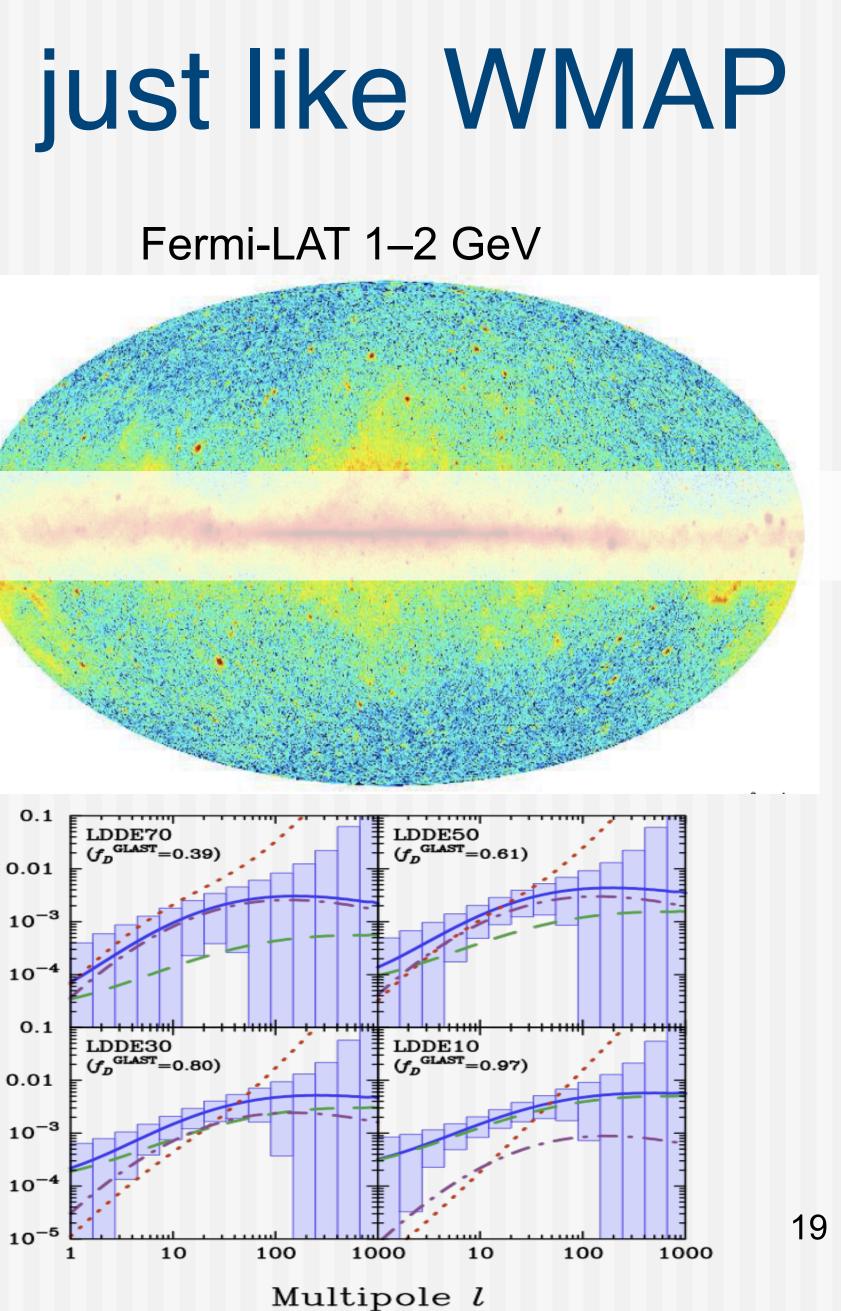
And, their spatial distribution can be calculated within the framework of Lambda-CDM model (using analytical calculations or numerical simulations)

Using Fermi Data, just like WMAP

WMAP 94GHz



 $l(l+1)C_l/2\pi$



Deciphering Gamma-ray Sky Astrophysical: Galactic vs Extra-galactic

- Galactic origin (diffuse)
 - E.g., Decay of neutral pions produced by cosmic-rays interacting with the interstellar medium.

Extra-galactic origin (discrete sources)

- Active Galactic Nuclei (AGNs)
- Blazars (Blazing quasars)
- Gamma-ray bursts

Exotic: Galactic vs Extra-galactic

Galactic Origin

- Dark matter annihilation in the Galactic Center
- Dark matter annihilation in the sub-halos within the Galaxy

Extra-galactic Origin

Dark matter annihilation in the other galaxies

Diffuse Gamma-ray Background

- First, we remove all the resolved (detected) sources from the Fermi-LAT map.
- Then, calculate the mean intensity of the map as a function of energies.
 - The intensity includes contributions from unresolved sources (below the detection threshold) and truly diffuse component (if any).

Why Anisotropy?

 $P_{\gamma}(E_{\gamma}, z, \hat{\boldsymbol{n}}r) = E_{\gamma} \frac{dN_{\gamma}}{dE_{\gamma}} \frac{\langle \sigma \boldsymbol{v} \rangle}{2} \left[\frac{\rho_{\chi}(z, \hat{\boldsymbol{n}}r)}{m_{\gamma}} \right]^{2},$

The shape of the power spectrum is determined by the structure formation, which is well known.

Schematically, we have:

(Anisotropy in Gamma-ray Sky) = (MEAN INTENSITY) $\times \Delta$

- The mean intensity depends on particle physics: annihilation cross-section and dark matter mass.
- The fluctuation power, Δ , depends on structure formation.

Power Spectrum

- Spherical harmonics transform of the intensity map:
 - $I(n) = \sum_{lm} a_{lm} Y_{lm}(n)$
- Squaring the coefficients and summing over m gives the power spectrum:

•
$$C_{I} = (2I+I)^{-1} \sum_{m} |a_{Im}|^{2}$$

 Just like we would do for the analysis of the CMB maps measured by WMAP.

Power Spectrum Formula $C_{l} = \int \frac{dr}{r^{2}} \{ W([1 + z]E_{\gamma}, r) \}^{2} P_{f}\left(k = \frac{l}{r}; r\right)$

• $P_f(k,z)$ is the power spectrum of "density squared," δ^2

 $\langle \tilde{f}_{k} \tilde{f}_{k'} \rangle = (2\pi)^{3} \delta^{(3)}$

where

$$f \equiv \delta^2 - \langle \delta^2$$

 $W(E_{\gamma}, z) = \frac{\langle \sigma v \rangle}{8\pi} \left(\frac{\Omega_{\chi} \rho_c}{m_{\gamma}}\right)^2 (1+z)^3 \frac{dN_{\gamma}}{dE_{\gamma}} e^{-\tau(E_{\gamma}, z)}$

$$^{(3)}(k + k')P_f(k)$$



Power Spectrum Formula $C_{l} = \int \frac{dr}{r^{2}} \{ W([1 + z]E_{\gamma}, r) \}^{2} P_{f}\left(k = \frac{l}{r}; r\right)$

• $P_f(k,z)$ is the power spectrum of "density squared," δ^2

 $\langle \tilde{f}_k \tilde{f}_{k'} \rangle = (2\pi)^3 \delta^{(3)} (\mathbf{k} + \mathbf{k'}) P_f(\mathbf{k})$

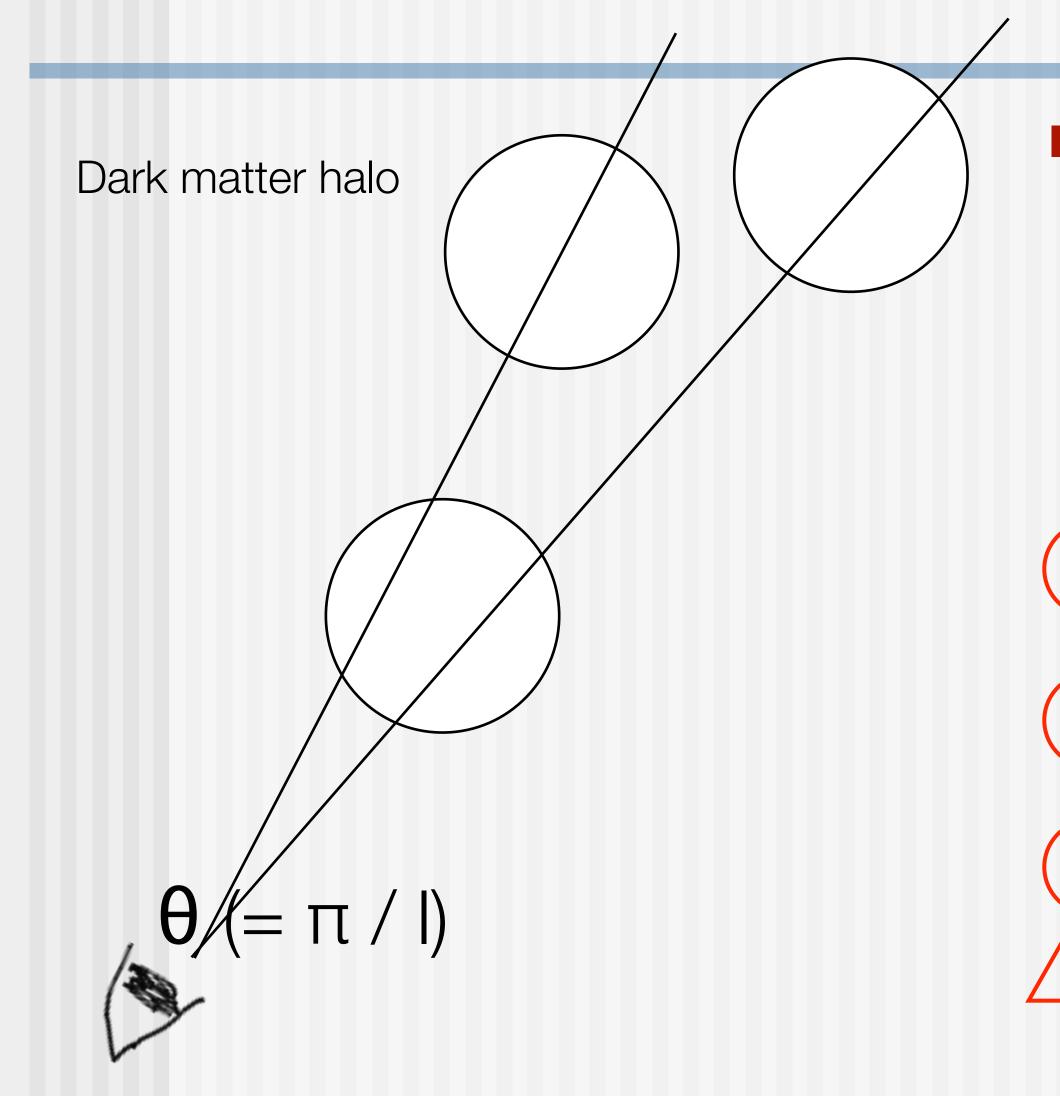
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2-point function of δ^2 \rightarrow = 4-point function 25

A Simple Route to the Angular **Power Spectrum**



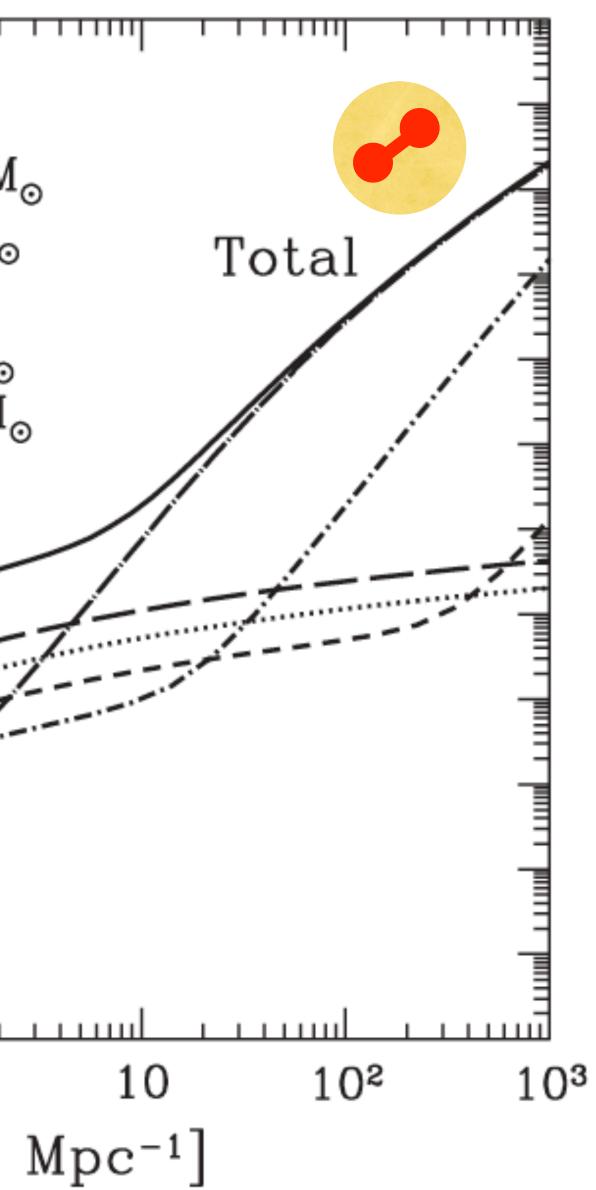
To compute the power spectrum of anisotropy from dark matter annihilation, we need three ingredients:

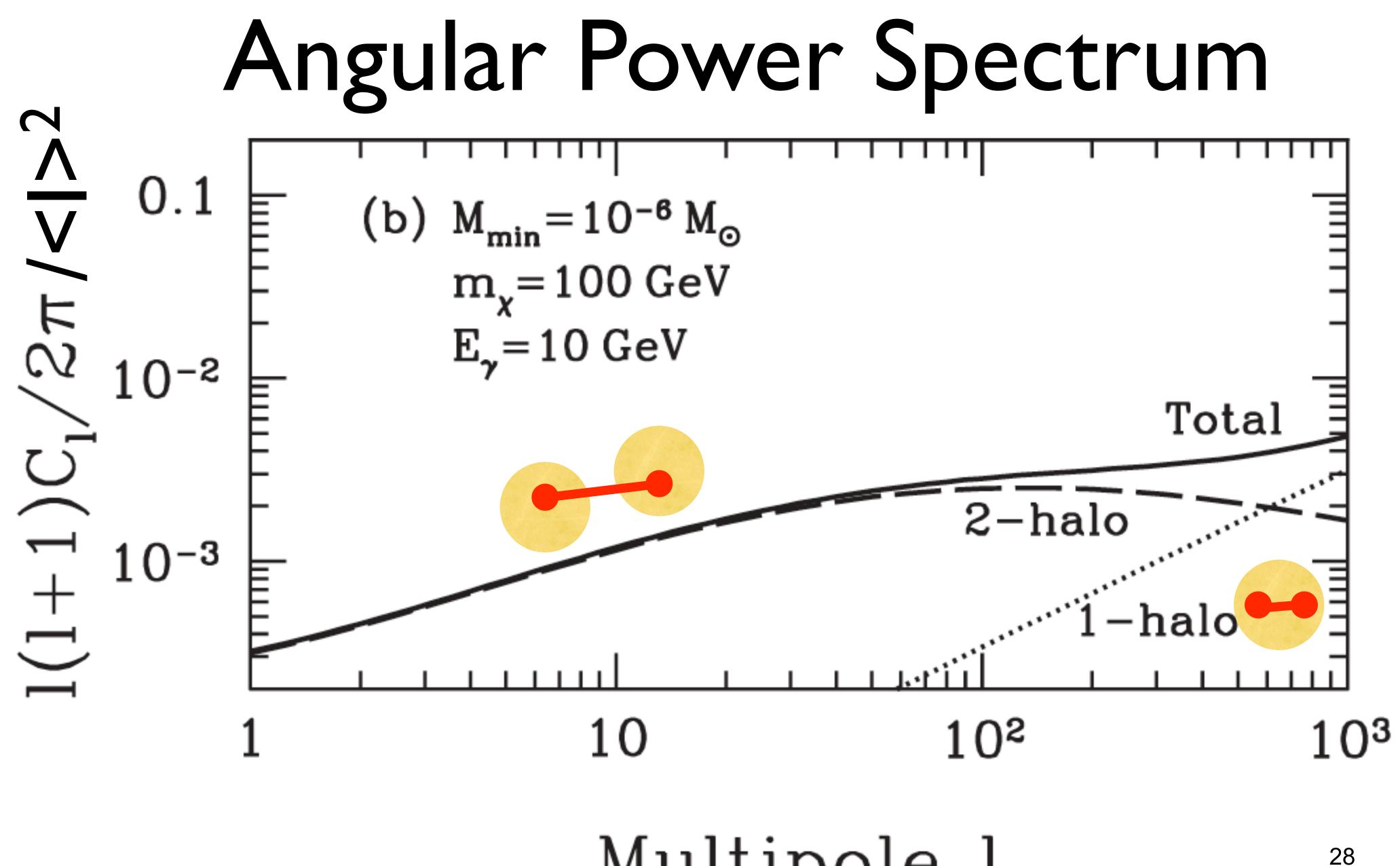
- Number of halos as a function of mass,
- 2.) Clustering of dark matter halos, and

3.) Dark matter density profile (NFW) 4. Substructure inside of each halo.

$$\Delta_{f}^{2}(k) \equiv \frac{k^{3}}{2\pi^{2}} \frac{P_{f}(k)}{\langle \delta^{2} \rangle^{2}} \stackrel{10^{2}}{}_{10}$$

trum of δ^2





Multipole l

Anisotropies in the Diffuse Gamma-ray Background Measured by the Fermi-LAT

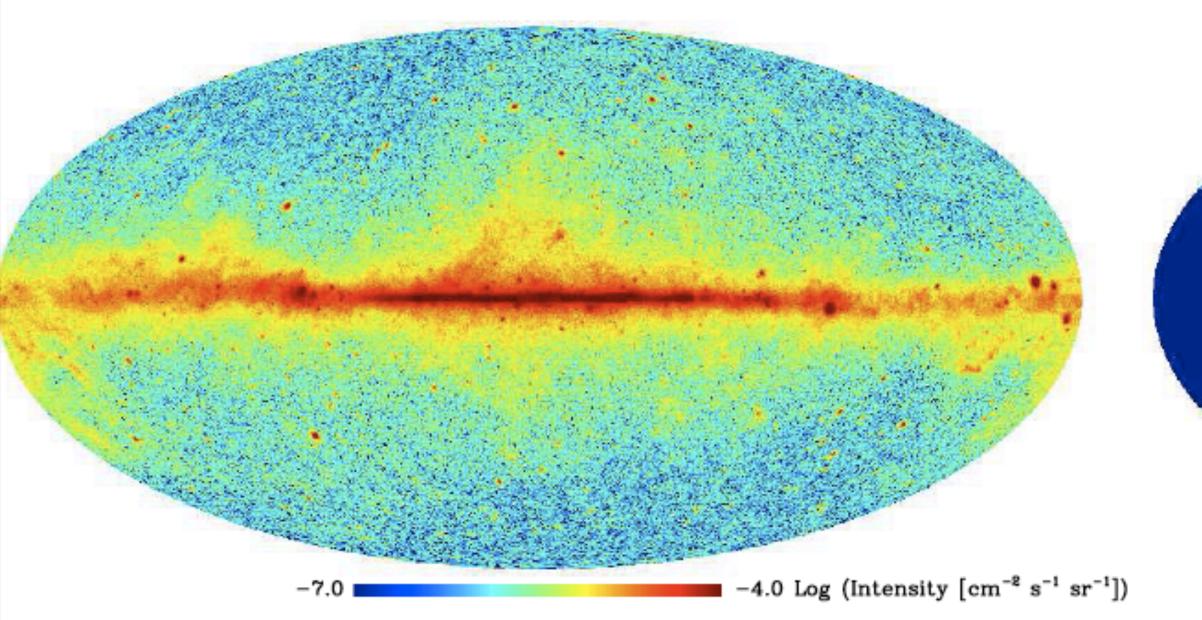
in collaboration with J. Siegal-Gaskins, A. Cuoco, T. Linden, M.N.Mazziotta, and V.Vitale (on behalf of Fermi-LAT Team)

Data Analysis

- Use the same Fermi-LAT map (~22mo, diffuse-class events)
- Apply the usual spherical harmonics transform, and measure the power spectrum!
 - $I(n) = \sum_{lm} a_{lm} Y_{lm}(n)$
 - $C_{l} = (2l+1)^{-1} \sum_{m} |a_{lm}|^{2}$
- Just like we did for the analysis of the CMB maps measured by WMAP.

1.0-2.0 GeV

DATA (P6_V3 diffuse), 1.0-2.0 GeV

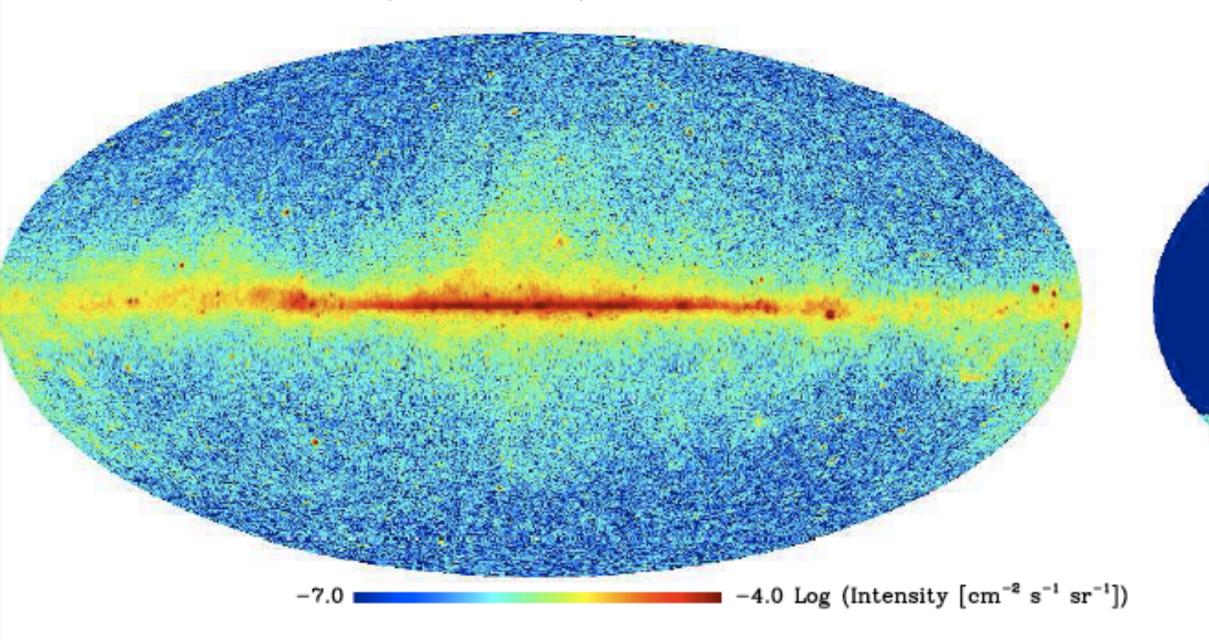


DATA (P6_V3 diffuse), 1.0-2.0 GeV

Mask b <30 degrees -4.0 Log (Intensity [cm⁻² s⁻¹ s -7.0

2.0-5.0 GeV

DATA (P6_V3 diffuse), 2.0-5.0 GeV



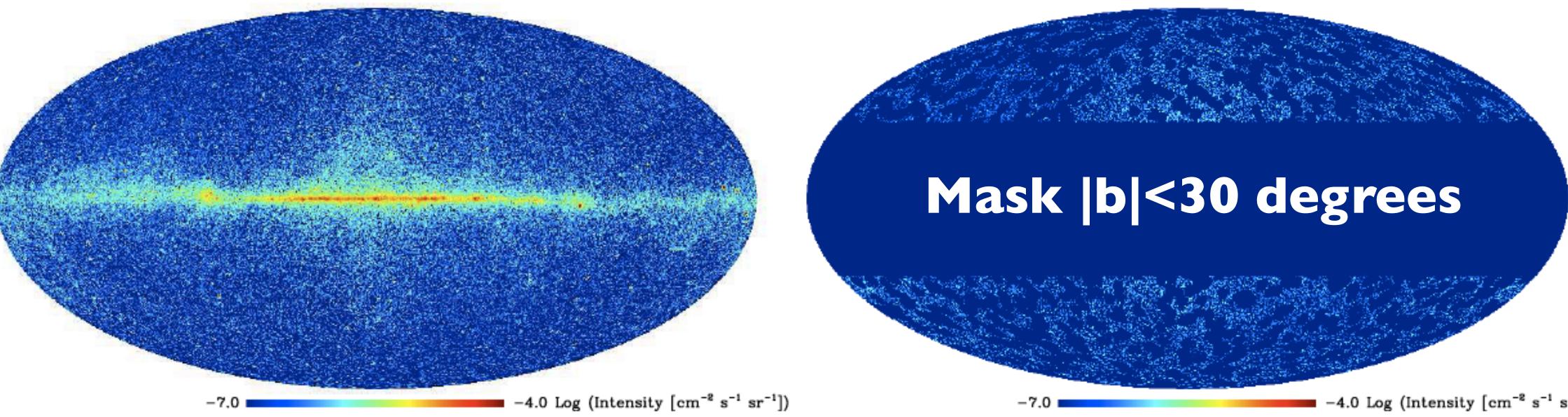
DATA (P6_V3 diffuse), 2.0-5.0 GeV

Mask b <30 degrees

-4.0 Log (Intensity [cm⁻² s⁻¹ s -7.0

5.0–10.4 GeV

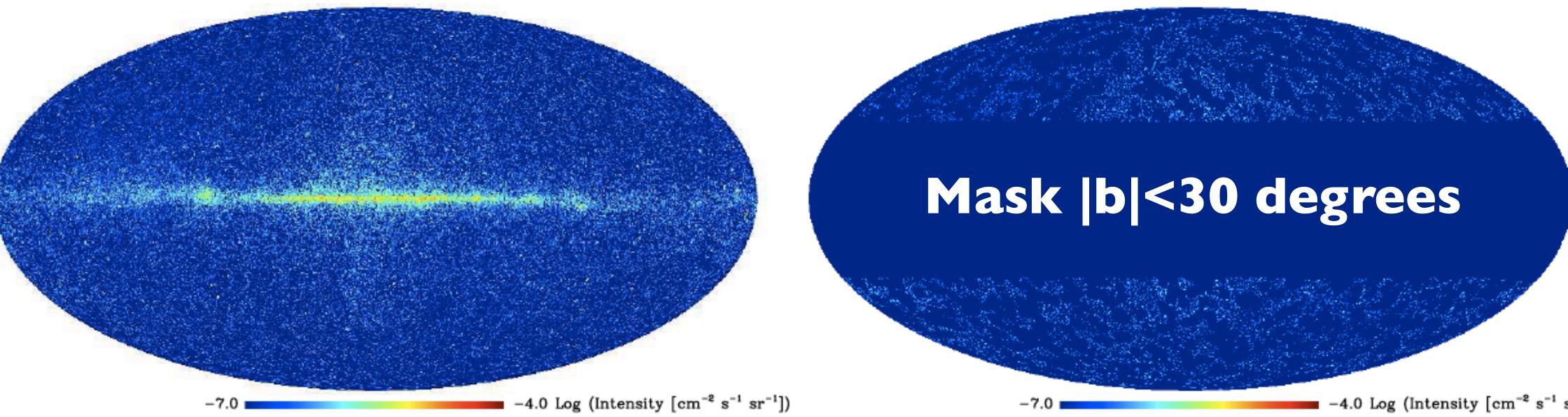
DATA (P6_V3 diffuse), 5.0-10.4 GeV



DATA (P6_V3 diffuse), 5.0-10.4 GeV

10.4–50.0 GeV

DATA (P6_V3 diffuse), 10.4-50.0 GeV



DATA (P6_V3 diffuse), 10.4-50.0 GeV

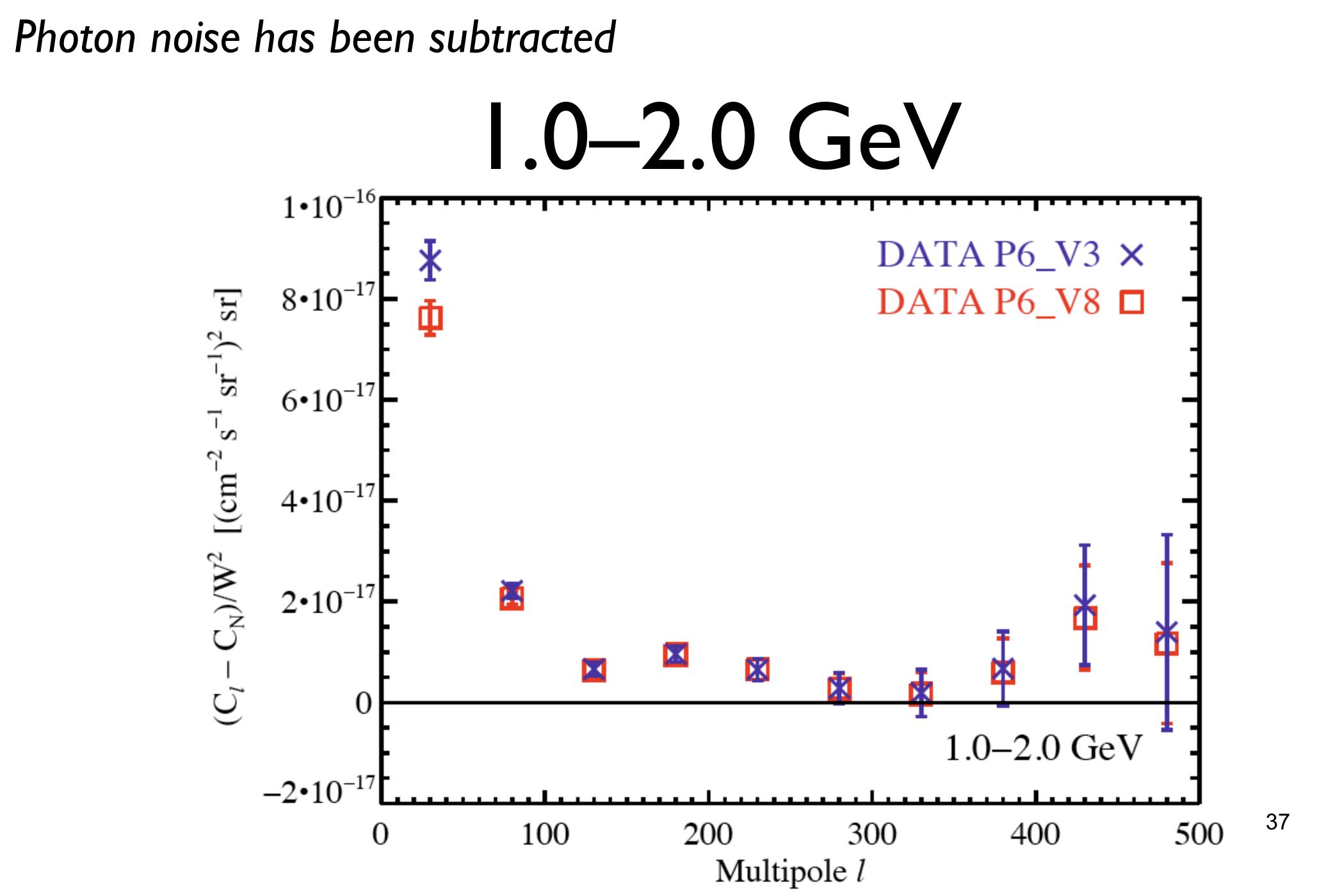
Fermi vs WMAP

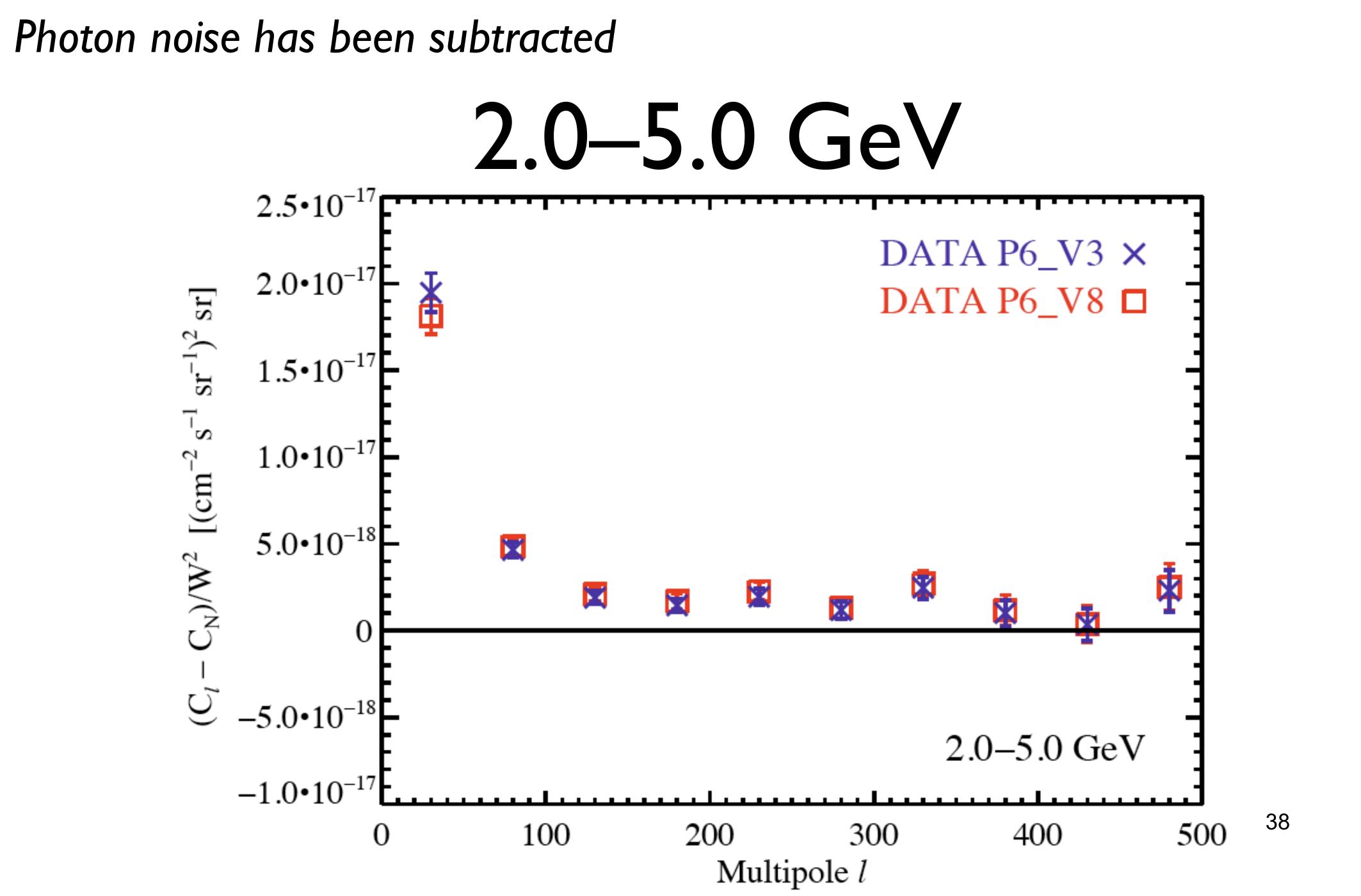
- There is an important difference between Fermi and WMAP maps
 - We count photons to produce Fermi maps; thus, there is the "photon noise" (Poisson statistics) in the power spectrum, which we must subtract.
 - Photon noise, C_N , is independent of multipoles, and is given by the mean number density of photons over the sky (which is precisely calculable).

Point Spread Function

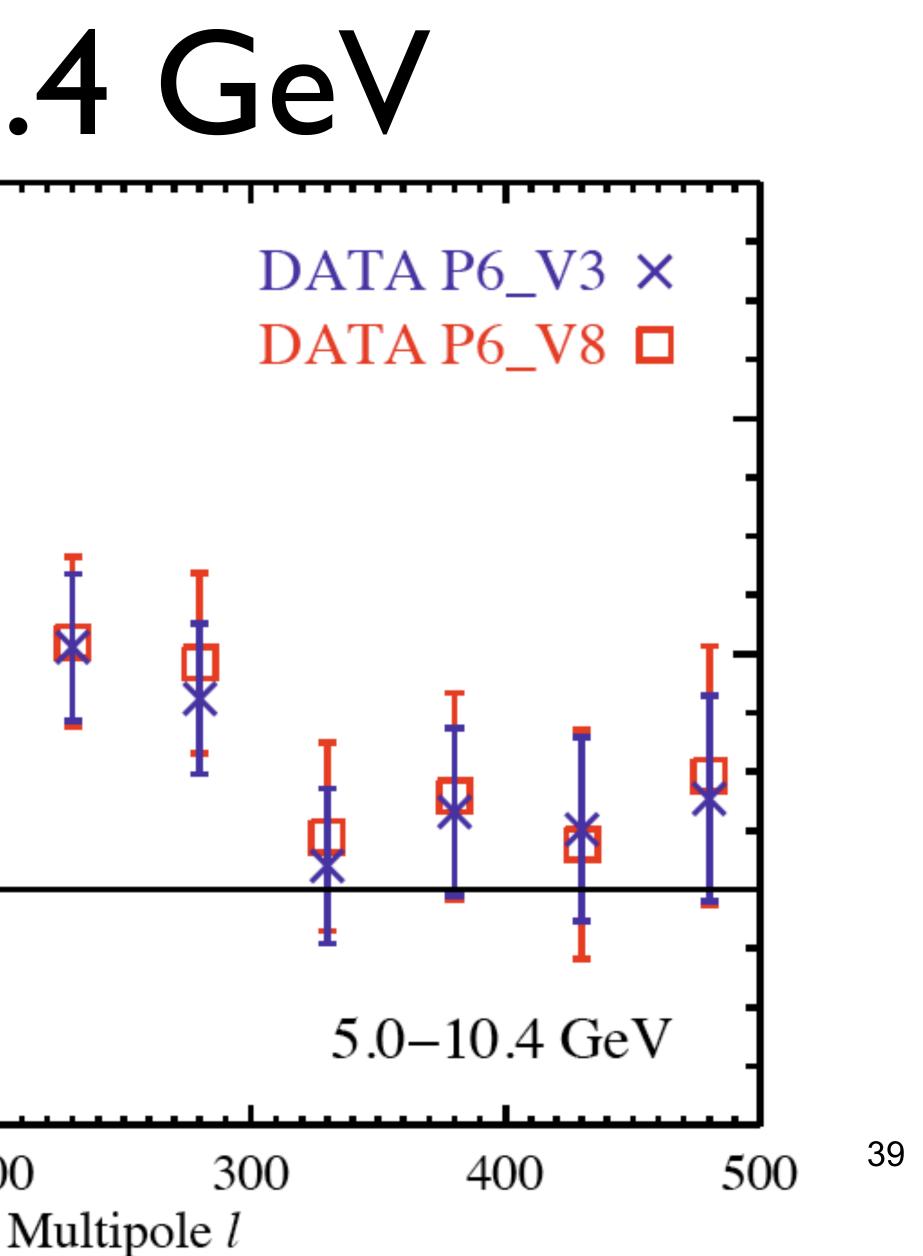
- The measured power spectrum is the true power spectrum multiplied by the harmonic transform of the "point spread function" (PSF). (It is called the "beam transfer function" in the WMAP analysis.)
- PSF is by no means a Gaussian we use two different versions of Fermi-LAT instrument response functions and compute PSF.
- We then compute $W_{\ell}^{\text{beam}}(E) =$
- The attenuation by PSF is corrected as $(C_I-C_N)/W_I^2$.
 - Two versions of PSF gave consistent answers.

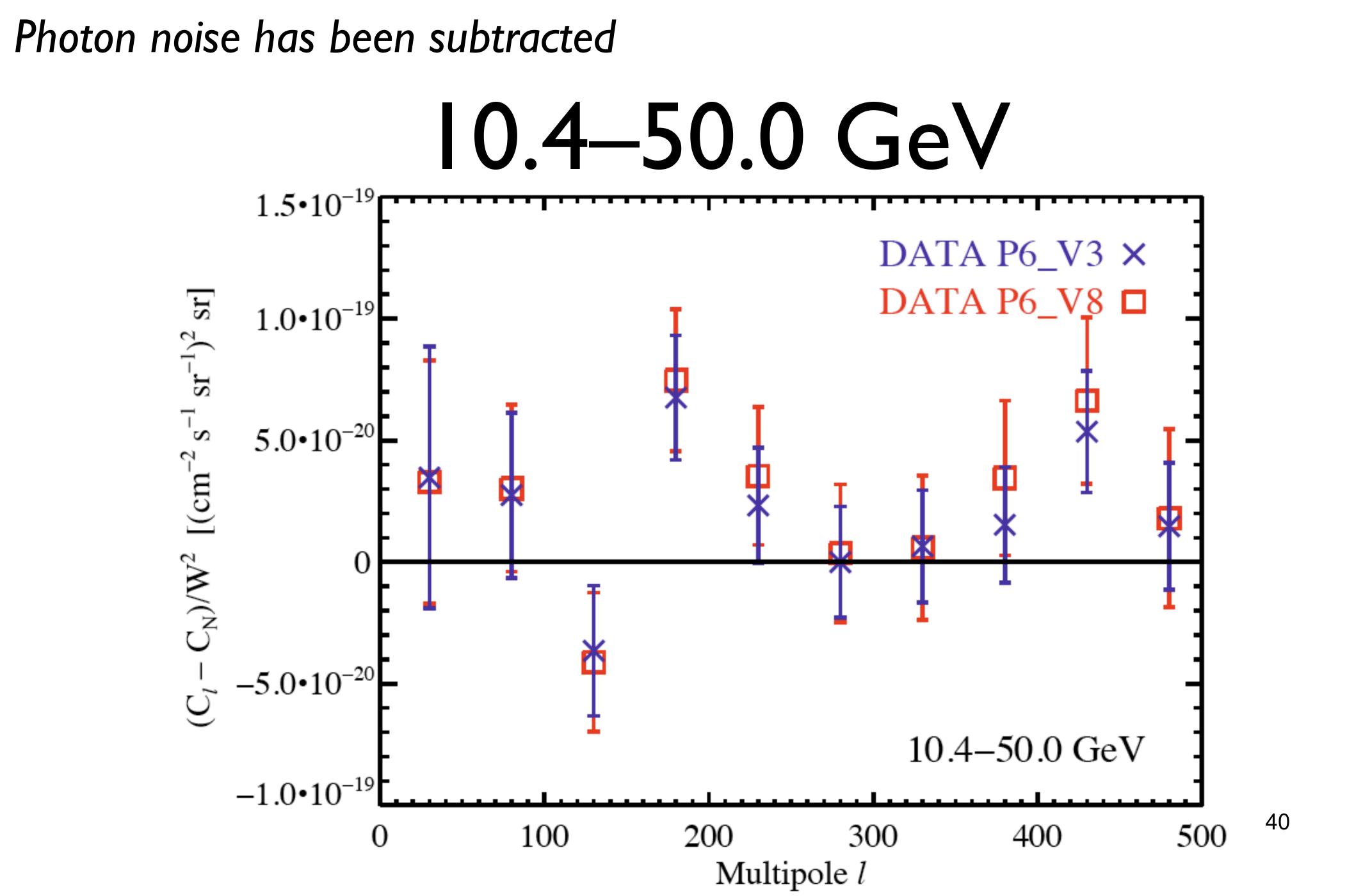
$$2\pi \int_{-1}^{1} d\cos\theta P_{\ell}(\cos(\theta)) \text{PSF}(\theta; \text{E})$$





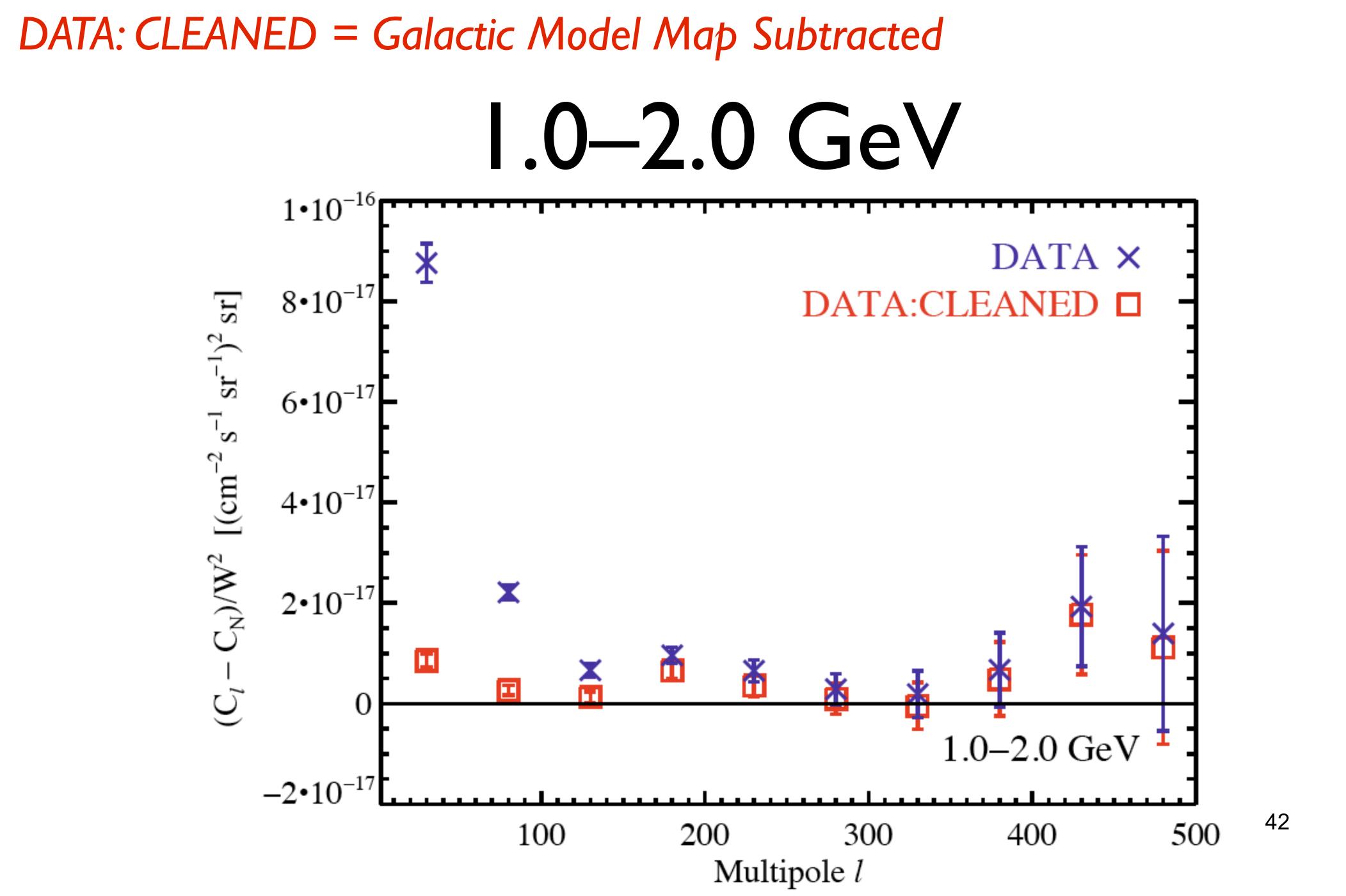
Photon noise has been subtracted 5.0–10.4 GeV 6•10⁻¹⁹ $[(\text{cm}^{-2} \text{ s}^{-1} \text{ sr}^{-1})^2 \text{ sr}]$ $4 \cdot 10^{-19}$ $2 \cdot 10^{-19}$ 閗 $(C_l - C_N)/W^2$ 0 $-2 \cdot 10^{-19}$ 100 200 0

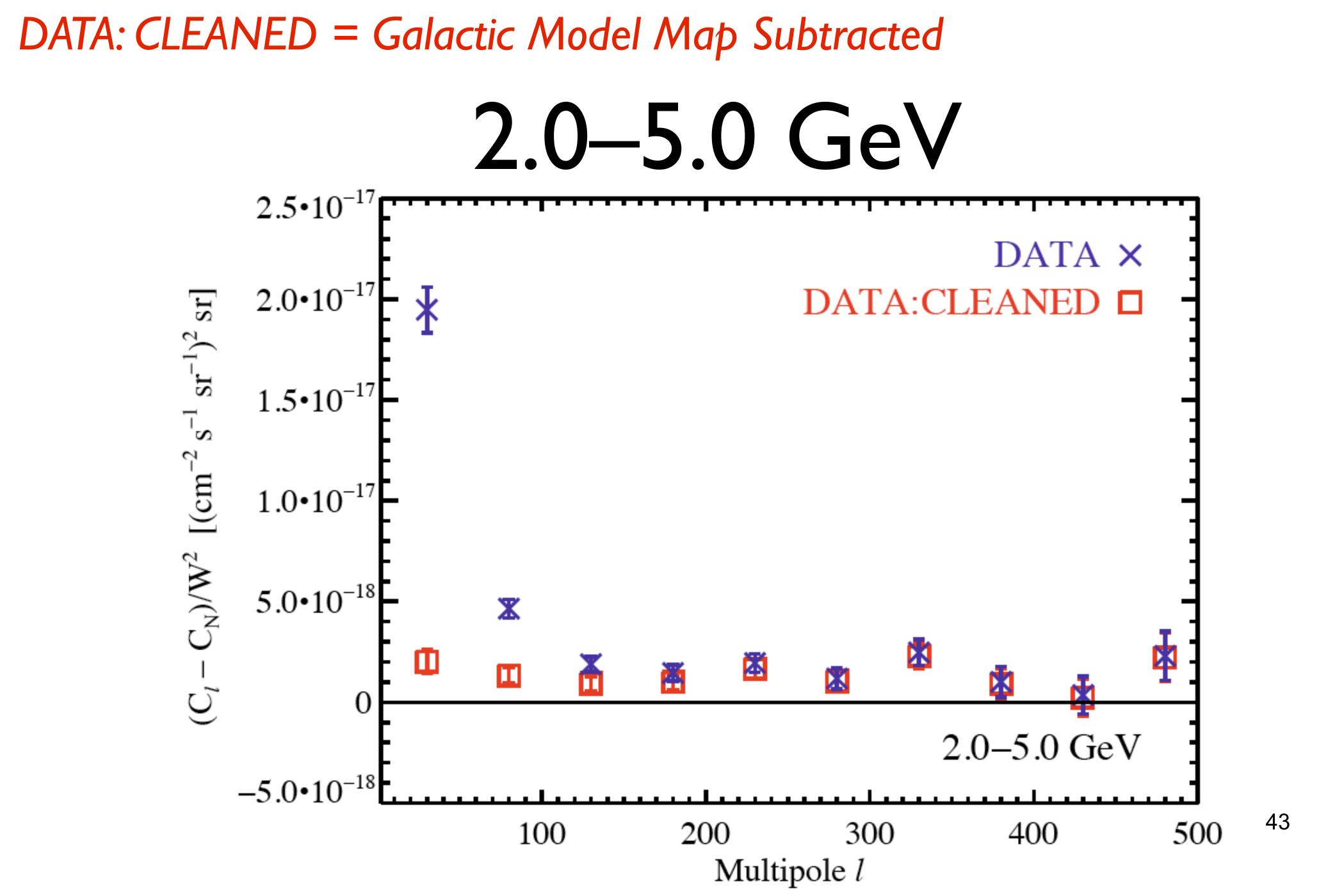


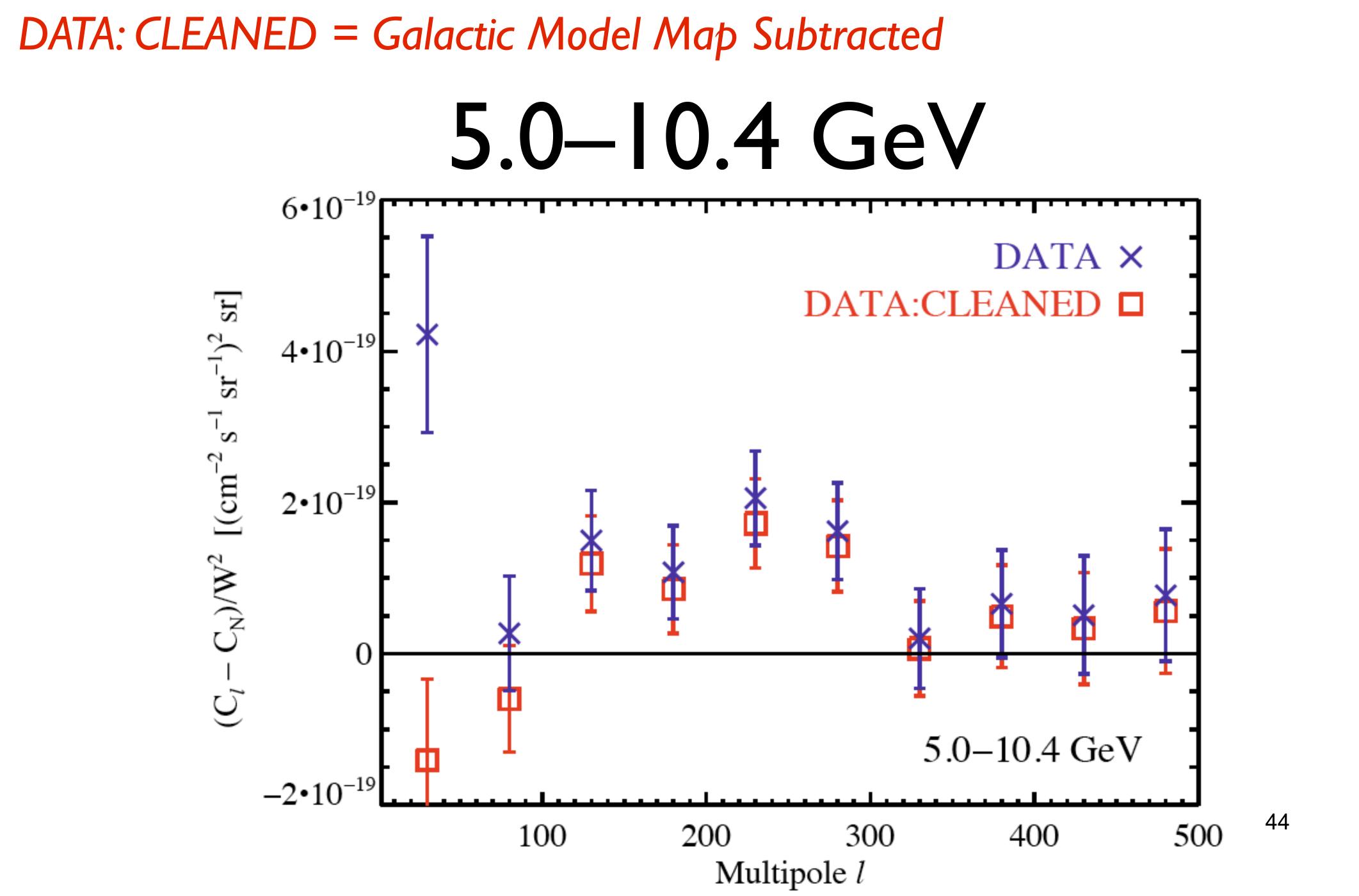


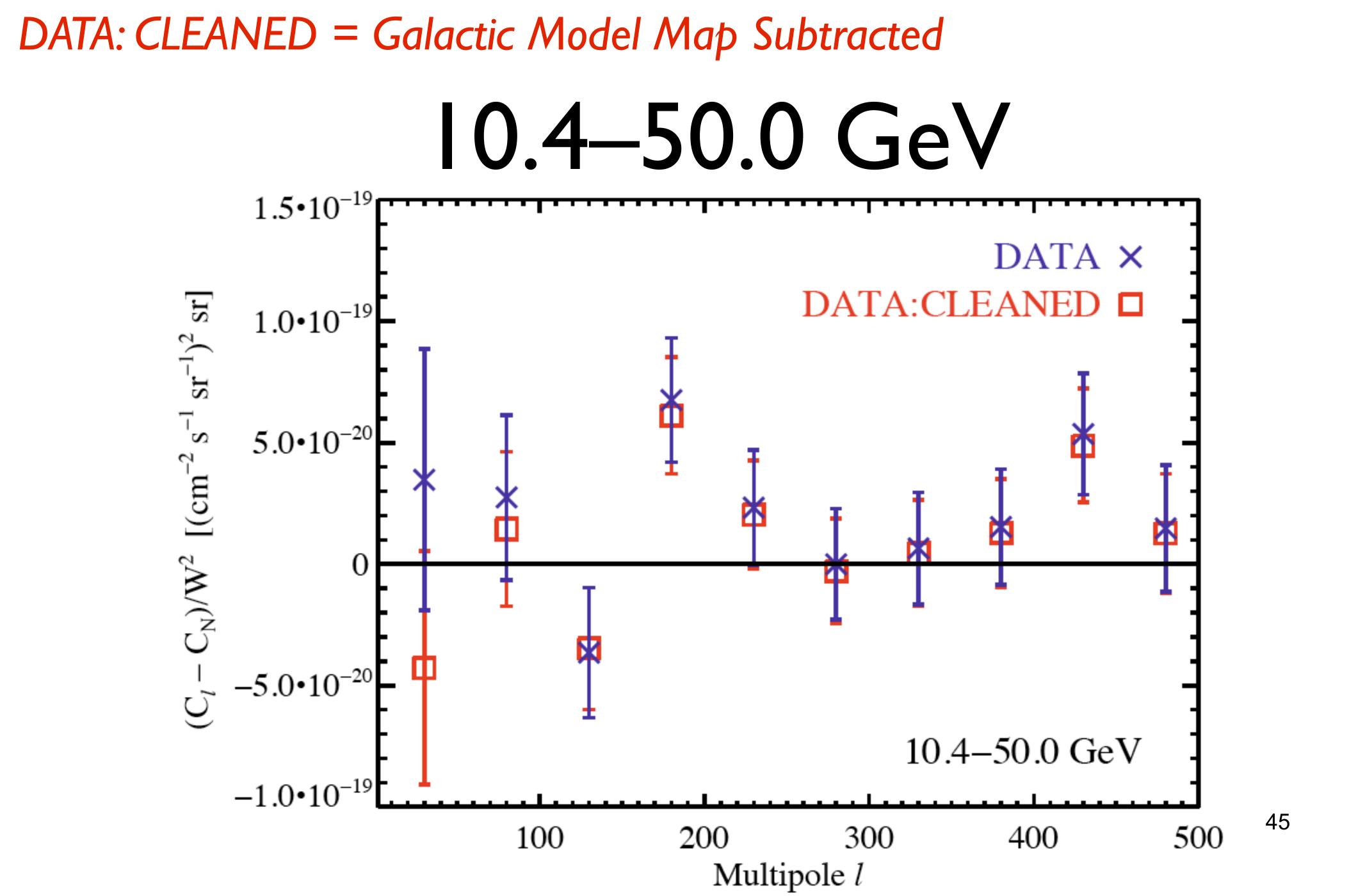
Observations

- At I<150, the power spectrum rises towards lower multipoles (larger angular scales).
 - The Galactic foreground contribution (more later)
- At I>150, we detect the excess power over the photon noise.
 - The excess power appears to be constant over multipoles, indicating the contribution from unclustered point sources (more later)









Focus on |>150

- The Galactic model maps indicate that the power we see at I<150 is largely coming from the Galactic foreground.
- The small-scale power at I>150 is not very much affected by the foreground, and thus is usable for investigating the extra-galactic gamma-ray background.

No Scale Dependence

• Fitting the measured power spectrum at I>150 to a single power-law: $C_{l} \sim l^{n}$

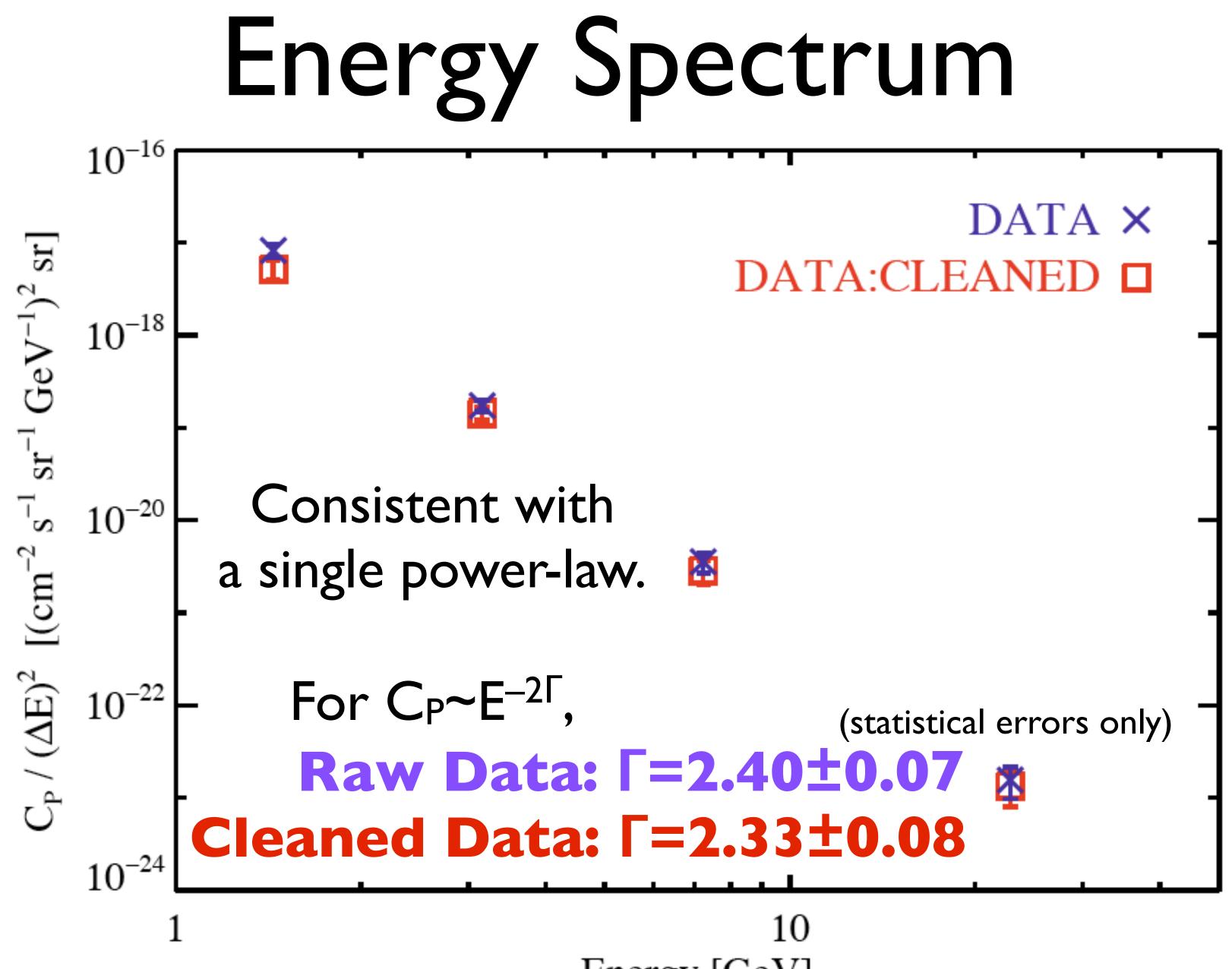
| E_{\min} | E_{\max} | n | $\chi^2/d.o.f.$ |
|------------|------------|------------------|-----------------|
| 1.04 | 1.99 | -1.33 ± 0.78 | 0.38 |
| 1.99 | 5.00 | -0.07 ± 0.45 | 0.43 |
| 5.00 | 10.4 | -0.79 ± 0.76 | 0.37 |
| 10.4 | 50.0 | -1.54 ± 1.15 | 0.39 |

Therefore, we will find the best-fitting constant power, C_{P} . ("P" stands for "Poisson contribution") 47

First detection of the extragalactic Y-ray anisotropy

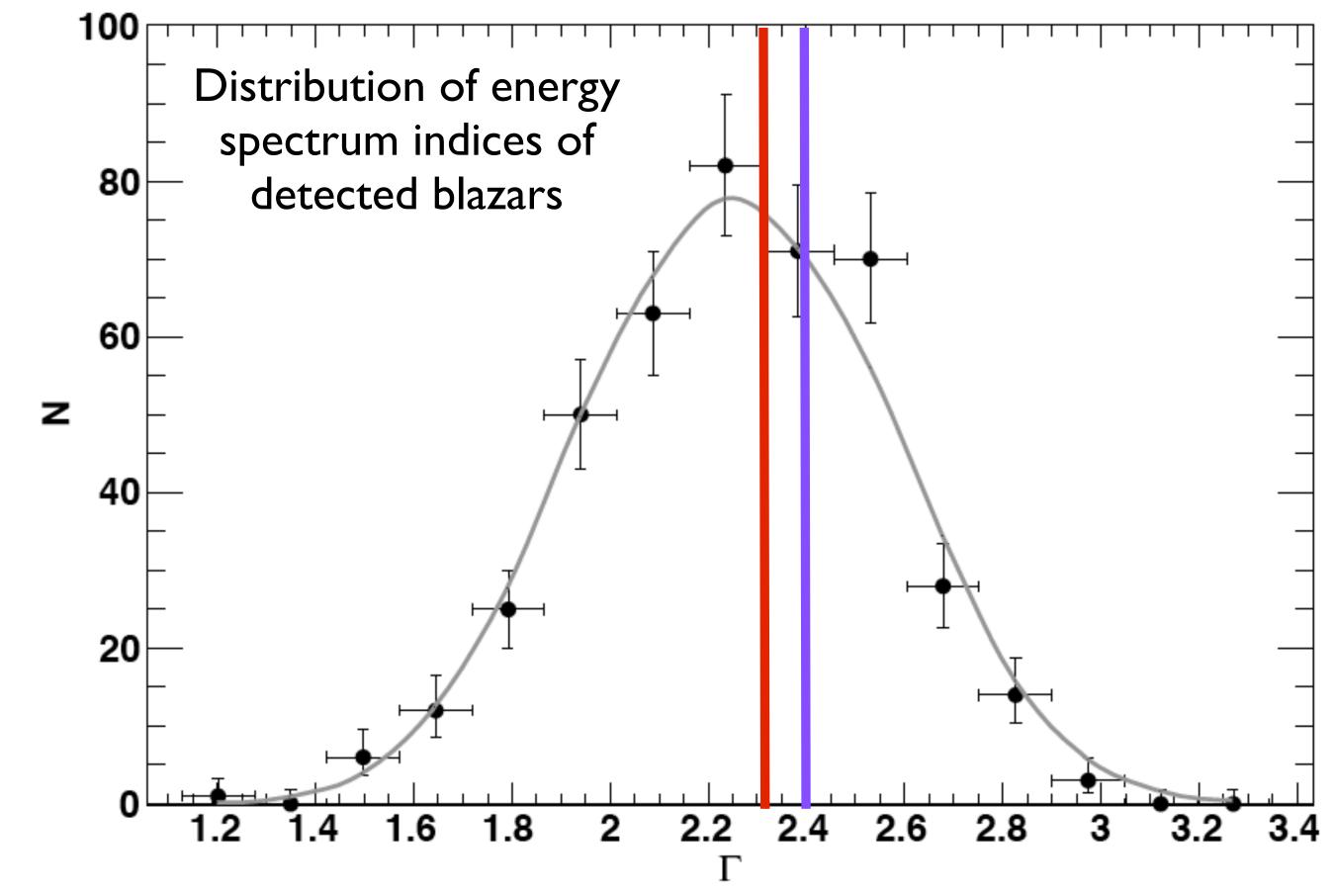
| | E_{\min} | E_{\max} | $C_{\mathbf{P}}$ | Significance |
|--------------|------------|------------|---|--------------|
| | [GeV] | [GeV] | $[(\mathrm{cm}^{-2} \mathrm{s}^{-1} \mathrm{sr}^{-1})^2 \mathrm{sr}]$ | |
| DATA | 1.04 | 1.99 | $7.39 \pm 1.14 \times 10^{-18}$ | 6.5σ |
| | 1.99 | 5.00 | $1.57 \pm 0.22 \times 10^{-18}$ | 7.2σ |
| | 5.00 | 10.4 | $1.06 \pm 0.26 \times 10^{-19}$ | 4.1σ |
| | 10.4 | 50.0 | $2.44 \pm 0.92 \times 10^{-20}$ | 2.7σ |
| DATA:CLEANED | 1.04 | 1.99 | $4.62 \pm 1.11 \times 10^{-18}$ | 4.2σ |
| | 1.99 | 5.00 | $1.30 \pm 0.22 \times 10^{-18}$ | 6.0σ |
| | 5.00 | 10.4 | $0.845 \pm 0.246 \times 10^{-19}$ | 3.4σ |
| | 10.4 | 50.0 | $2.11 \pm 0.86 \times 10^{-20}$ | 2.4σ |

• Many-sigma detections up to 10 GeV!



Energy [GeV]

Fermi-LAT Collaboration, ApJ, 720, 435 (2010) Are we seeing blazars?

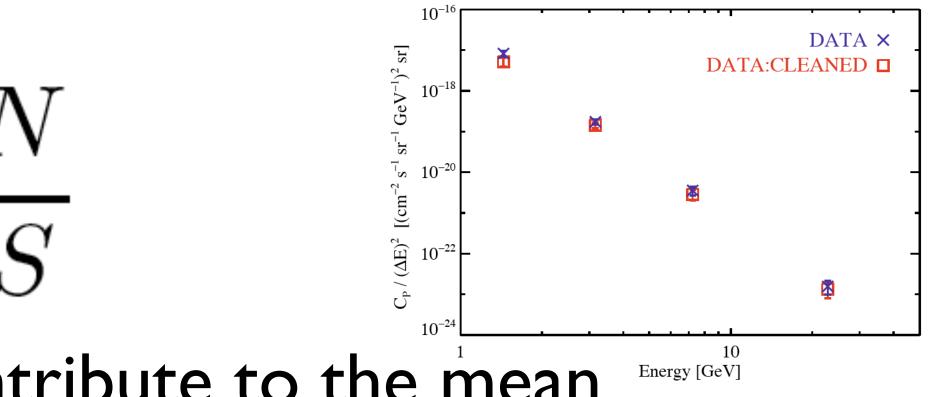


 The energy spectrum of anisotropy (from unresolved sources) agrees with that of **detected** blazars.

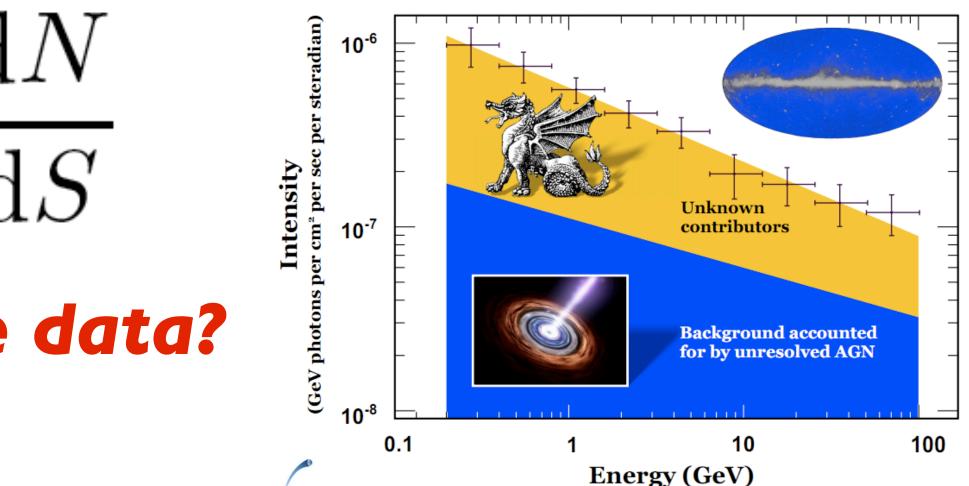
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Interpreting the Results

- Unresolved, unclustered point sources contribute to C_P as $C_{\rm P} = \int_0^{S_c} dS \ S^2 \frac{\mathrm{d}N}{\mathrm{d}S}$
- Unresolved, point sources contribute to the mean intensity as $\int_{0}^{S_{c}} dS S \frac{\mathrm{d}N}{\mathrm{d}S}$
- Are they consistent with the data?

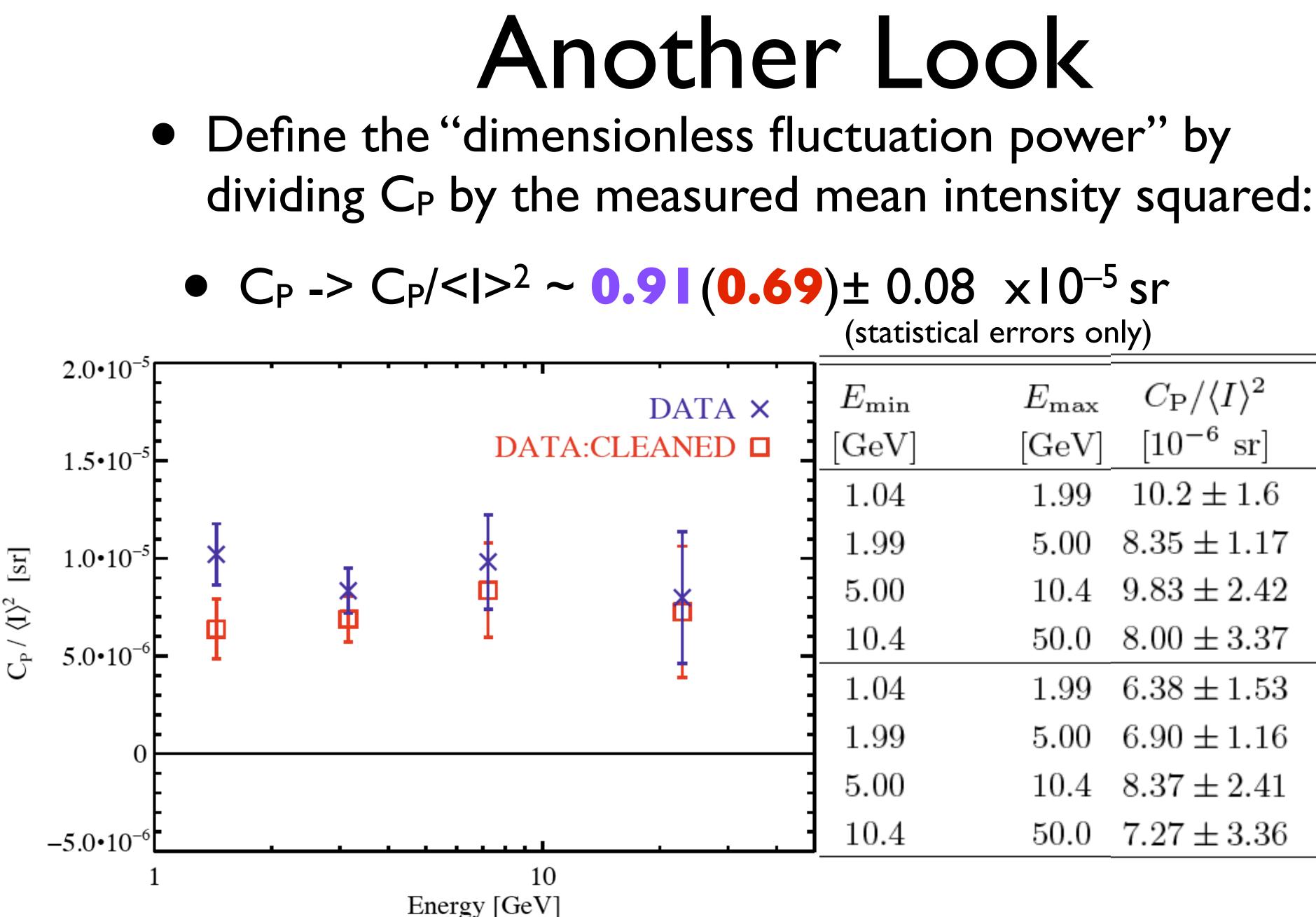


Fermi LAT Extragalactic Gamma-ray Background



The answer seems YES

- Our results are consistent with the following interpretation:
 - The detected anisotropy is largely due to unresolved blazars.
 - The amplitude of anisotropy is consistent with the fact that the same unresolved blazars contribute only to a fraction of the mean gamma-ray background.
- These two, independent measurements give us a consistent picture of the gamma-ray sky.



(statistical errors only)

| E_{\min} | $E_{\rm max}$ $C_{\rm P}/\langle I \rangle^2$ | Significance |
|------------|---|--------------|
| [GeV] | [GeV] [10 ⁻⁶ sr] | |
| 1.04 | $1.99 10.2 \pm 1.6$ | 6.5σ |
| 1.99 | 5.00 8.35 ± 1.17 | 7.1σ |
| 5.00 | $10.4 9.83 \pm 2.42$ | 4.1σ |
| 10.4 | 50.0 8.00 ± 3.37 | 2.4σ |
| 1.04 | $1.99 6.38 \pm 1.53$ | 4.2σ |
| 1.99 | 5.00 6.90 ± 1.16 | 5.9σ |
| 5.00 | 10.4 8.37 ± 2.41 | 3.5σ |
| 10.4 | 50.0 7.27 ± 3.36 | 2.2σ |

What about Dark Matter?

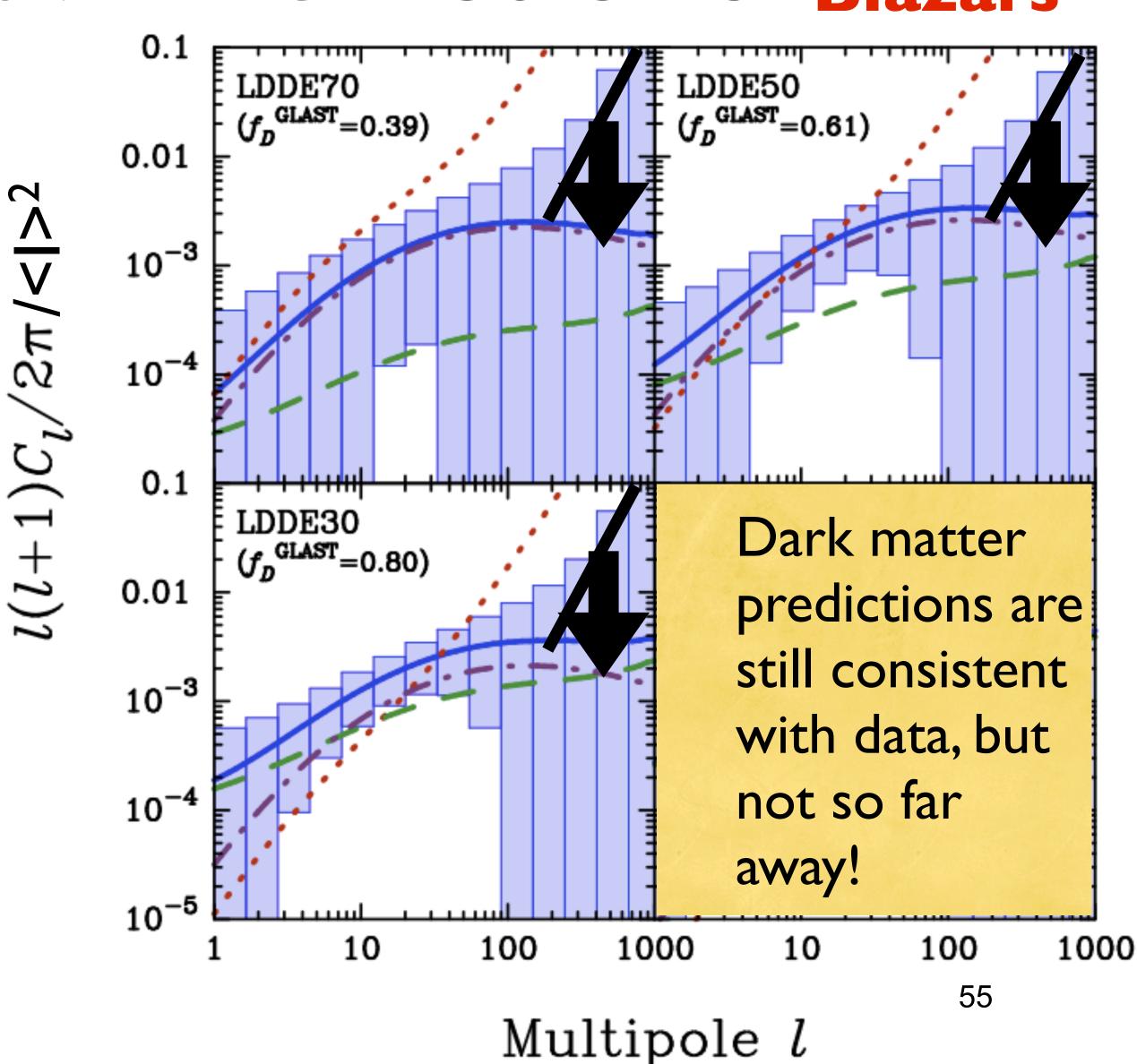
- Our results can be used to place limits on the dark matter properties.
- Subtracting the blazar contribution, the upper limit on the constant power at |>150 is

• $C_P / <|>^2 < |0^{-6} sr$

• What would this mean?

Ando & EK (2006); Ando, EK, Narumoto & Totani (2007) 2006/2007 Predictions Blazars

- Watch out for the factor of I(I+I).
 - Poisson spectrum gives ~l²
- We constrain C_I only at I>150



Bottom-line Message

- We have the new observable: **power spectrum of** the gamma-ray background.
- And, it has been detected from the data.

Conclusions

- We have detected anisotropy in the extra-galactic diffuse gamma-ray background from Fermi-LAT 22mo maps.
- The detected anisotropy is consistent with the contribution from unresolved blazars
 - Also consistent with the mean intensity data
 - The origin of the bulk of diffuse background remains a mystery
- Dark matter annihilation contributions may not be so far away from the current limit. Wait for results from the future Fermi analysis (3 to 7 more years to go!) 57