

# ***Fundamental physics without spacetime: ideas, results and challenges from quantum gravity and beyond***

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Bayes Forum

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# Plan

1. spacetime as we know it, in GR - and the relational perspective
2. the problem of Quantum Gravity: conceptual, physical, mathematical
3. emergent space, emergent time in QG
4. an example of a fundamental QG formalism and of the emergence of spacetime from it
5. a possible quantum statistical foundation of the formalism based on Jaynes' principle
6. remarks about foundational/philosophical issues (and the role of agency) in light of QG

different kind of talk

this talk

other talks at  
the Bayes Forum



goal

outline general issues in QG, and some research directions, more than specific results, as well broader implications

general survey, with lots of material (mostly for later discussion)

mostly focusing on conceptual aspects



**Nature of spacetime:**

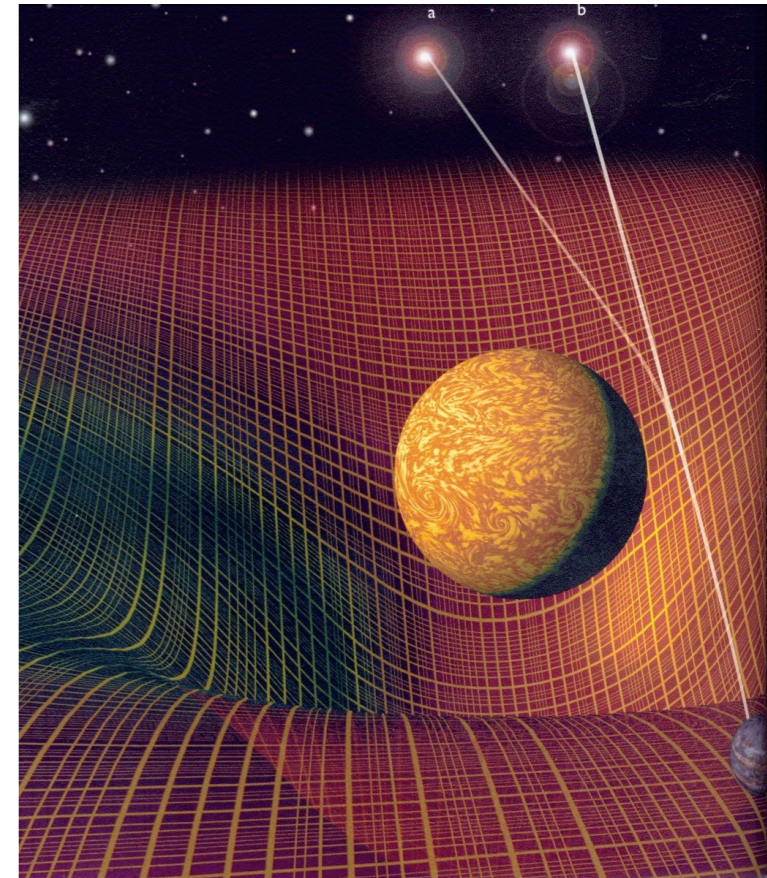
**lessons from General Relativity**

# Nature of spacetime: lessons from General Relativity

main lesson: spacetime is a physical system itself (own dofs, see gravitational waves)

$$g_{\mu\nu}(t, x) \quad ds^2 = g_{tt}dt^2 + g_{tx_i}dtdx_i + g_{x_ix_j}dx_idx_j$$

- gravity = spacetime geometry (spatial distances, time intervals, curvature of space, volumes, ..... ) = field
- spacetime geometry is generically non-flat and dynamical
- matter mass-energy “deforms” spacetime, deformation affects motion of matter
- no preferred space or time direction
- physics is the same in all (idealized) frames
- causal structure non-trivial and dynamical; spatial regions can be causally inaccessible (horizons)



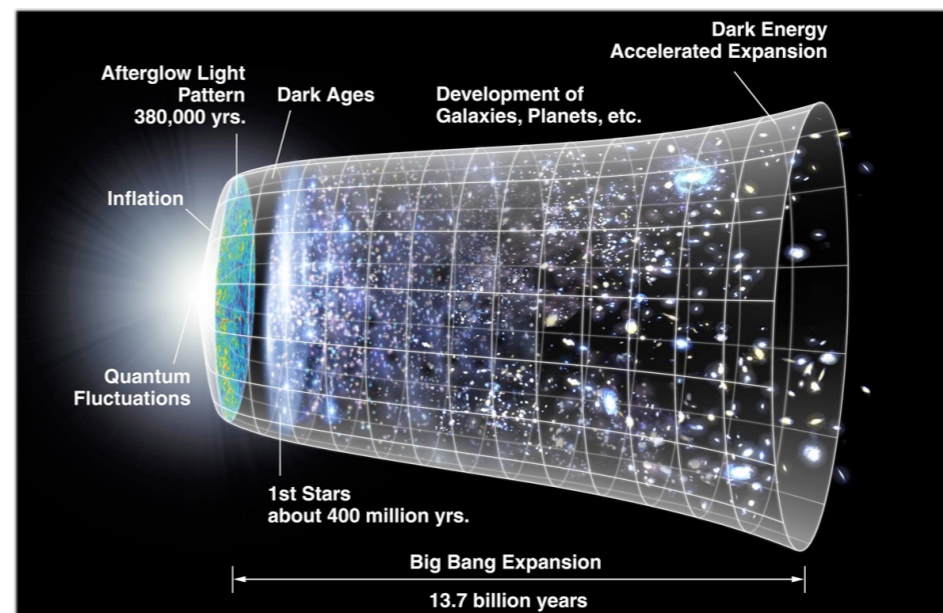
Einstein's equations

$$R_{\mu\nu}[g(x)] - \frac{1}{2}R[g(x)] + \Lambda g_{\mu\nu}(x) = 8\pi G_N T_{\mu\nu}[\phi(x), \dots]$$

give eqns for evolution of universe:

e.g. Friedmann eqn (homogeneity + isotropy):

$$H^2 = \left(\frac{\dot{a}}{a}\right)^2 = \frac{8\pi G\rho}{3} - \frac{kc^2}{a^2}$$





**Nature of spacetime:**

**lessons from General Relativity**

**in fact, there is more to it .....**

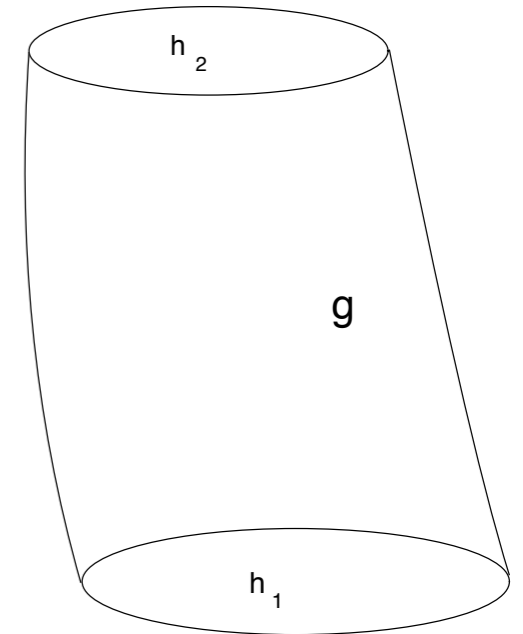


# Diffeo-invariance, spacetime (local) observables and relational strategy

diffeomorphism invariance + background independence

D. Giulini, '06

- no absolute notion of temporal or spatial direction/location/distance
- manifold has only global role (topological restriction)
- local manifold structures (points, directions, paths, coordinate frames, ...) have no physical significance
- what is physical is values of (continuum) dynamical fields, among which the metric field, and their relations

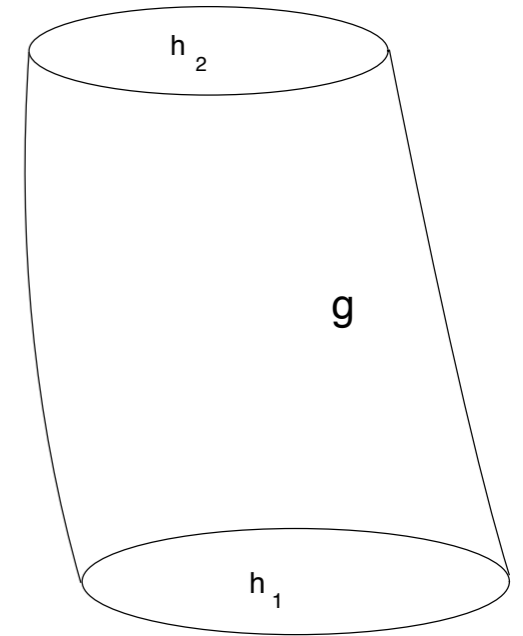


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so what are spacetime observables?

only "global" ones?

$$\mathcal{O} = \int_{\mathcal{M}} d^4x \sqrt{-g} O(g(x), \varphi(x))$$

yes, wrt manifold, because manifold is not spacetime

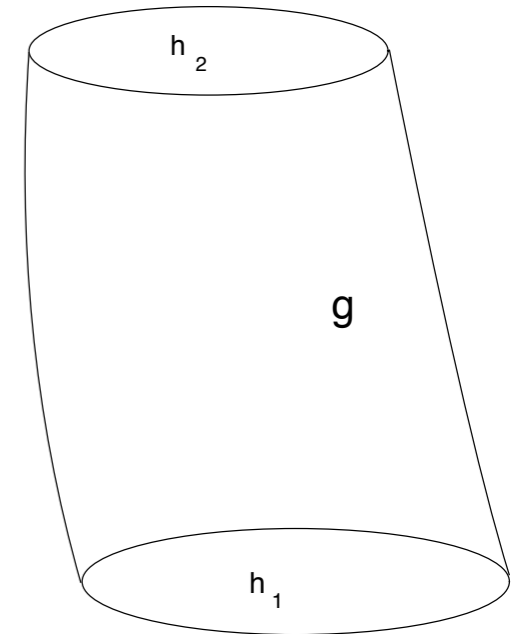
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**relational strategy:**

Rovelli '90s+ [related ideas DeWitt '60s; Bargmann & Komar 90's]

Hoehn, '00s

$R(t) \quad \Phi(t)$

$\implies$

$R(t) \quad t(\Phi)$

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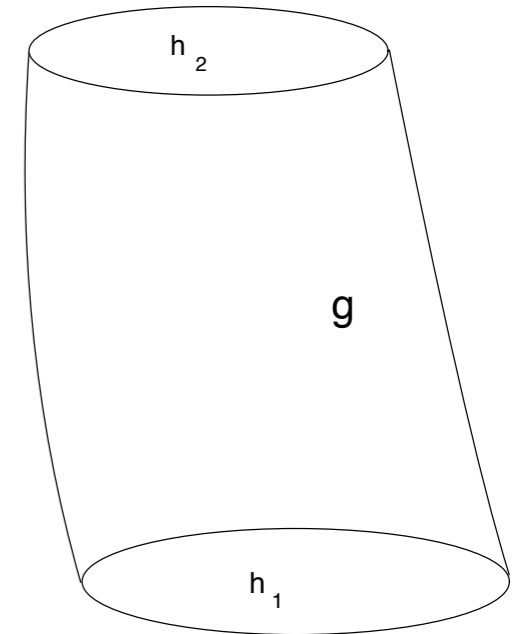
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ideally, spacetime physics should only be expressed in terms of such relational quantities

points, coordinates, trajectories on manifold are "useful fictions" representing physical frames (clocks and rods) in the limit in which their physical properties (energy, interactions, ...) are negligible

relational perspective: physics is in the relations between dynamical fields  $g_{\mu\nu}(x)$   $A_\mu(x)$   $\varphi(x)$   
 (complete, Dirac) observables = correlations on superspace (space of fields)

simplest example: parametrized pendulum

### classical single 1d pendulum

physical quantities: pendulum position as function of physical time  $Q = Q(T)$  (value of some clock)

dynamics:  $\frac{d^2 Q}{dT^2} = -\omega^2 Q \longrightarrow$  general solution:  $Q(T) = A \sin(\omega T + \phi)$

true physical system is pendulum + clock

physics is in the relation  $Q(T)$

$Q$  and  $T$  can be measured (partial observables); what can be predicted is only  $Q(T)$  (complete observable)

### parametrized classical single 1d pendulum

turn dynamical variables into functions of new "time parameter" (i.e. scalar fields in  $d=1$ ):  $Q(\tau)$   $T(\tau)$

$$\frac{dQ}{d\tau} = P_Q \quad \frac{dT}{d\tau} = P_T \quad H(Q, P_Q, T, P_T) = P_T(\tau) + \frac{1}{2} P_Q^2(\tau) + \frac{1}{2} \omega^2 Q^2(\tau)$$

$$\frac{dQ}{d\tau} = P_Q = \frac{dH}{dP_Q} \quad \frac{dT}{d\tau} = P_T = \frac{dH}{dP_T} = 1 \quad \frac{dP_Q}{d\tau} = P_Q = -\frac{dH}{dQ} = -\omega^2 Q \quad \frac{dP_T}{d\tau} = -\frac{dH}{dT} = 0$$

+ invariance (covariance of equations) under 1d diffeos:  $\tau \rightarrow f(\tau)$  **1d manifold not physical**

only diffeo-invariant observable, evaluated on solutions on the dynamics, is:  $Q(T) = A \sin(\omega T + \phi)$

$Q(\tau)$   $T(\tau)$  are neither measurable nor predictable (as functions of affine parameter)

only  $Q(T)$  (complete observable) can be predicted -  $Q$  and  $T$  are only "physical" in relational sense

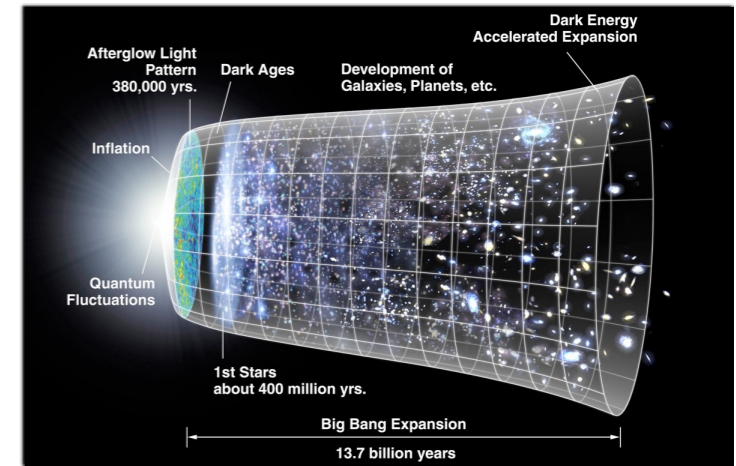
diffeomorphism invariance indicates what is physical and what is not

general point: physics is on superspace (space of field configurations), not manifold (only auxiliary structure)

difficult to express/extract it in general QG case

things much simpler in cosmological context

restriction to global features of universe: (approximately) homogeneous fields



example: flat Friedmann universe (homogeneous, isotropic)  $ds^2 = -N^2(t)dt^2 + a^2(t)\delta_{ab}dx^a dx^b$

dynamical variables = scale factor (universe volume) and massless scalar field

GR action reduces to:  $S = \frac{3}{8\pi G} \int dt N \left( -\frac{aV_0\dot{a}^2}{N^2} + \frac{V}{N} \frac{\dot{\chi}^2}{2N} \right)$   $V \equiv V_0 a^3$  invariant under 1d diffeos

configuration space is 2d flat manifold  $\{a, \chi\}$  only relational observable  $V(\chi)$

can be fully deparametrized to give relational evolution:  $\left( \frac{1}{3V} \frac{dV}{d\chi} \right)^2 \equiv \left( \frac{V'}{3V} \right)^2 = \frac{4\pi G}{3}$

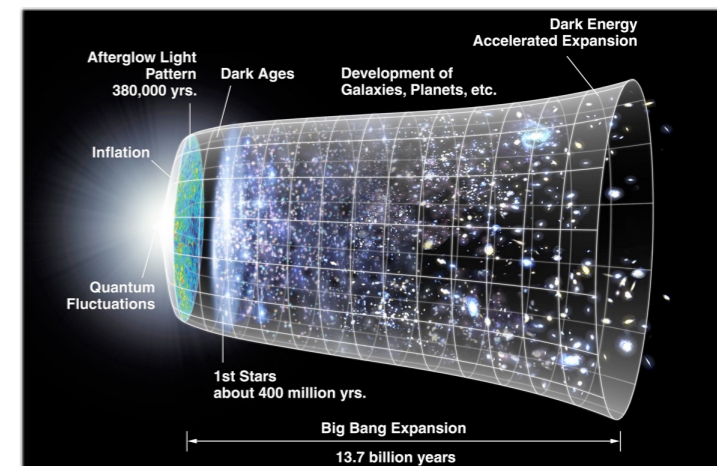
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summary to identify "spacetime = manifold" or "spacetime physics = physics on manifold" is approximation at best

(corresponds to case in which set of four scalar fields behave like test fields covering manifold, and can be used as coordinates for manifold points)

do not expect to find manifold etc neither at fundamental QG level, nor in its effective description

Note: physics may be different with respect to different PHYSICAL reference frames!

# **Lessons:**

**spacetime physics is "fields (values) in relation to fields (values)"**

**physical frames and physical covariance: observer matters**



# **The Quantum Gravity problem**

## starting point: conceptual, physical, mathematical clash

framework and ingredients of GR are incompatible with what we learned from Quantum Mechanics

### GR

- spacetime (geometry) is a dynamical entity itself
- there are no preferred temporal (or spatial) directions
- physical systems are local and locally interacting
- everything (incl. spacetime) evolves deterministically
- all dynamical fields are continuous entities
- every property of physical systems (incl. spacetime) and of their interactions can be precisely determined, in principle

### QFT

- spacetime is fixed background for fields' dynamics
- evolution is unitary (conserved probabilities) with respect to a given (preferred) temporal direction
- nothing can be perfectly localised
- everything evolves probabilistically
- interaction and matter fields are made of "quanta"
- every property of physical systems and their interactions is intrinsically uncertain, in general

- in fact, no proper understanding of interaction of geometry with quantum matter, if gravity is not quantized

$$R_{\mu\nu} - \frac{1}{2}g_{\mu\nu}R + \Lambda g_{\mu\nu} = \frac{8\pi G}{c^4} \langle \Psi | \hat{T}_{\mu\nu} | \Psi \rangle$$

not a consistent fundamental theory

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two frameworks come with different associated mathematical language and tools

conceptual + mathematical clash is clear

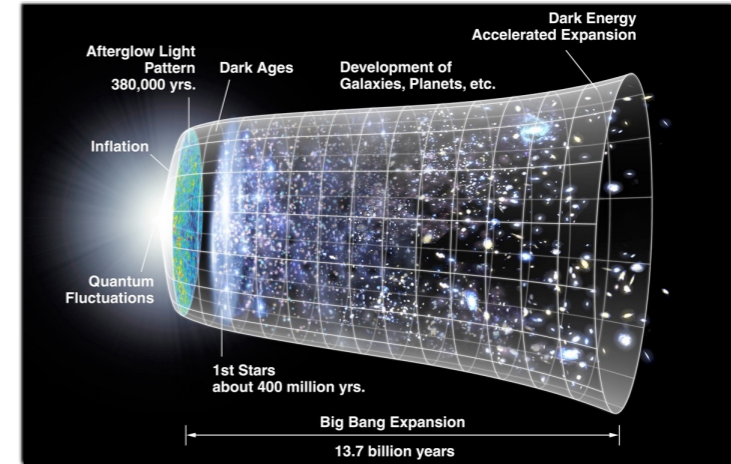
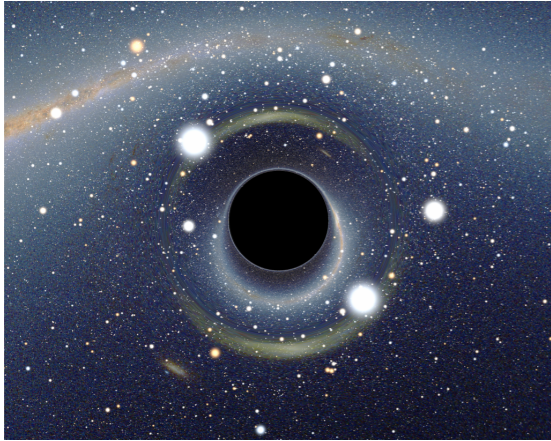
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## summary of physical issues

- spacetime singularities: breakdown of GR for strong gravitational fields/large energy densities - inevitable in classical GR  
center of black holes, big bang - quantum effects expected to be important



- cosmological scenarios for the early universe need QG completion

R. Brandenberger, '10, '11, '14

why a close to homogeneous and isotropic universe?  
why an approximately scale invariant power spectrum?  
what happens "at" the Big Bang?

### Inflation

- what produces inflation?
- physics of trans-Planckian modes (for long inflation)?
- inflation too close to Planck regime?
- inflationary spacetime still contains singularity

### Bouncing cosmology

- new physics needed to describe/justify cosmological bounce

### Emergent universe (pre-big bang static phase)

- static phase and phase transition require new physics

## summary of physical issues

new QG dofs? primordial  
(quantum) black holes?

new type of matter?

new QG dof?

cosmological constant?

why doesn't it gravitate?

modified gravity?

why holographic entropy?

spacetime microstructure?

violation of unitarity? locality? .....

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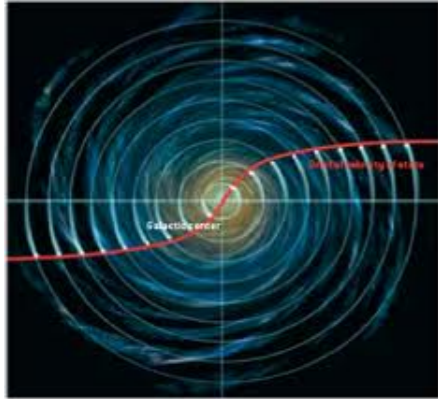
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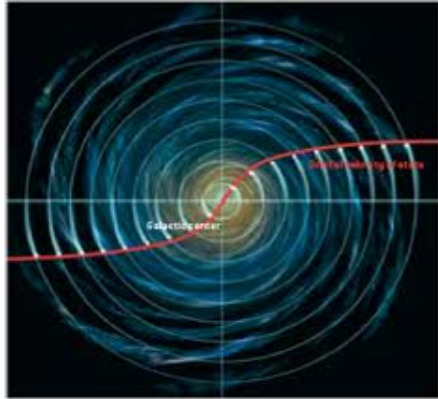
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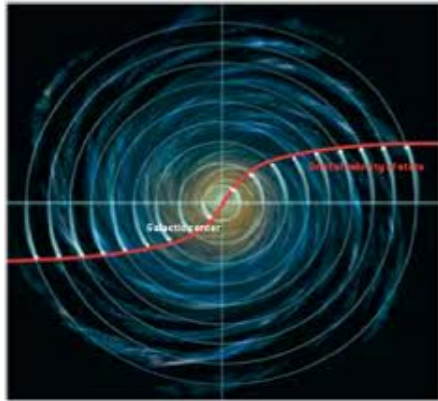
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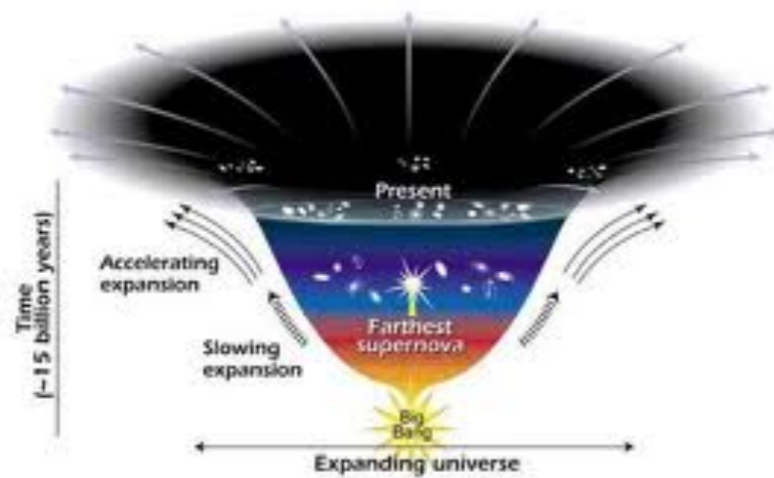
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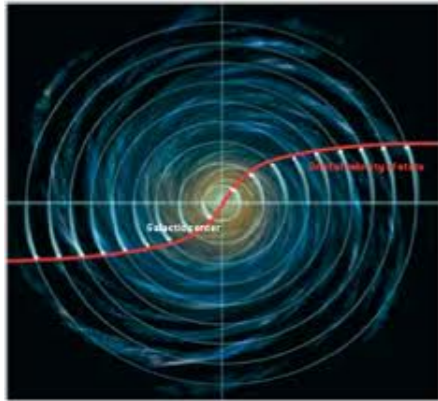
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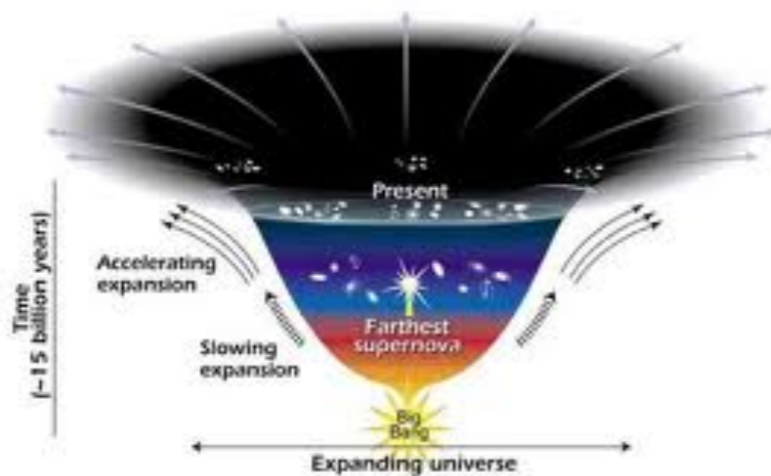
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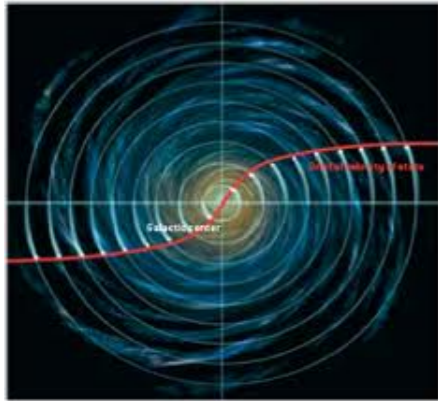
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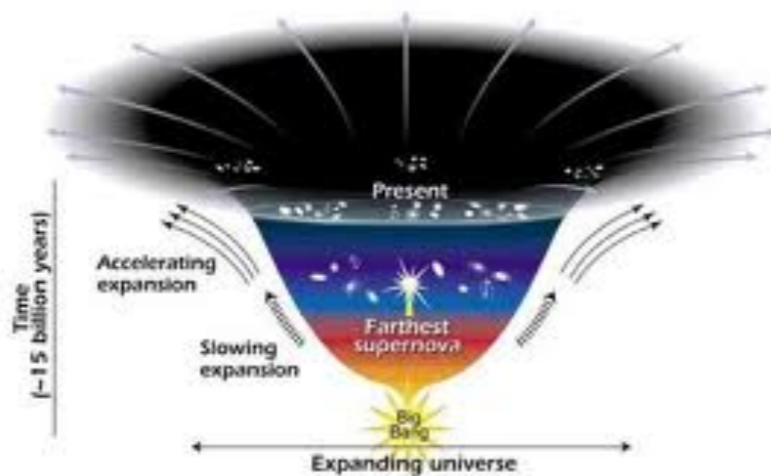
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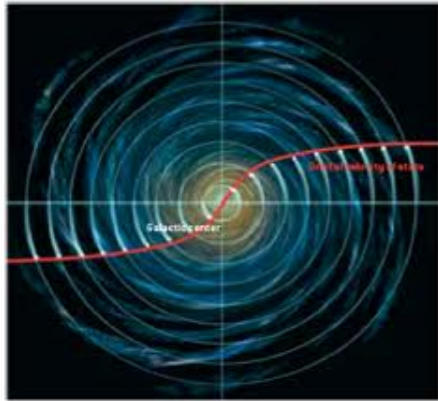
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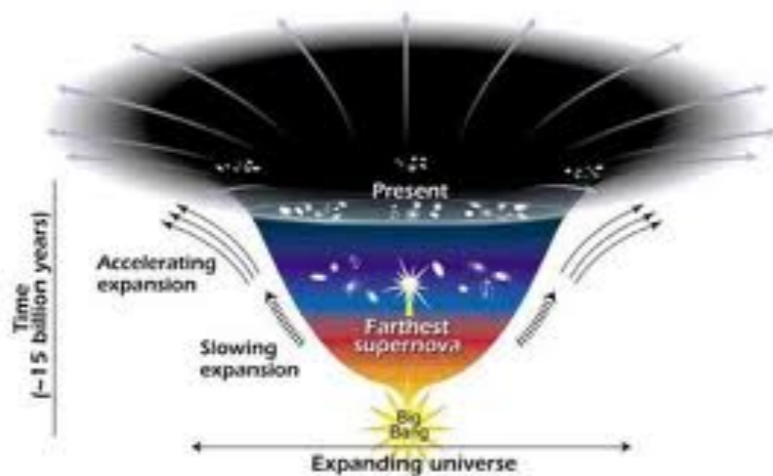
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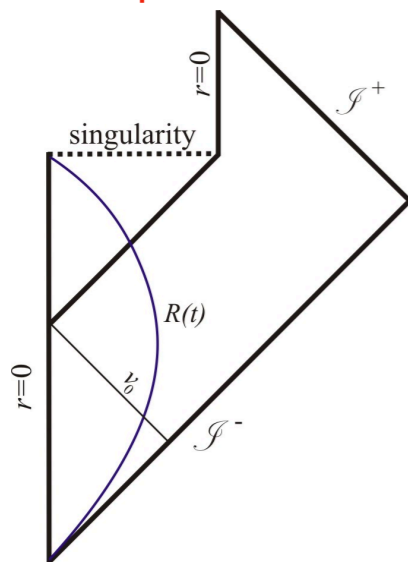
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why difficult?

spacetime and geometry (and matter) should become "quantum" physical systems themselves

full non-perturbative quantum theory of gravitational field (not just perturbations around spacetime background)



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simple perturbative methods fail, non-perturbative methods nightmare

canonical:

Hilbert space of physical (diffeo-invariant)  
quantum geometries (incl. scalar product)

$\mathcal{H} \ni |h_{ij}\rangle = | \text{spatial geometry} \rangle =$   
 $= | \text{spatial distances, curvature, volumes, ...} \rangle$

algebra of observables (distances, curvature, volumes, ...)

semiclassical approximation

covariant:

non-perturbative QG path integral  
(sum-over-geometries), incl. measure

$$\langle h_2 | h_1 \rangle = \sum_{g_{\mu\nu} | h_1, h_2} \mathcal{A}(g)$$

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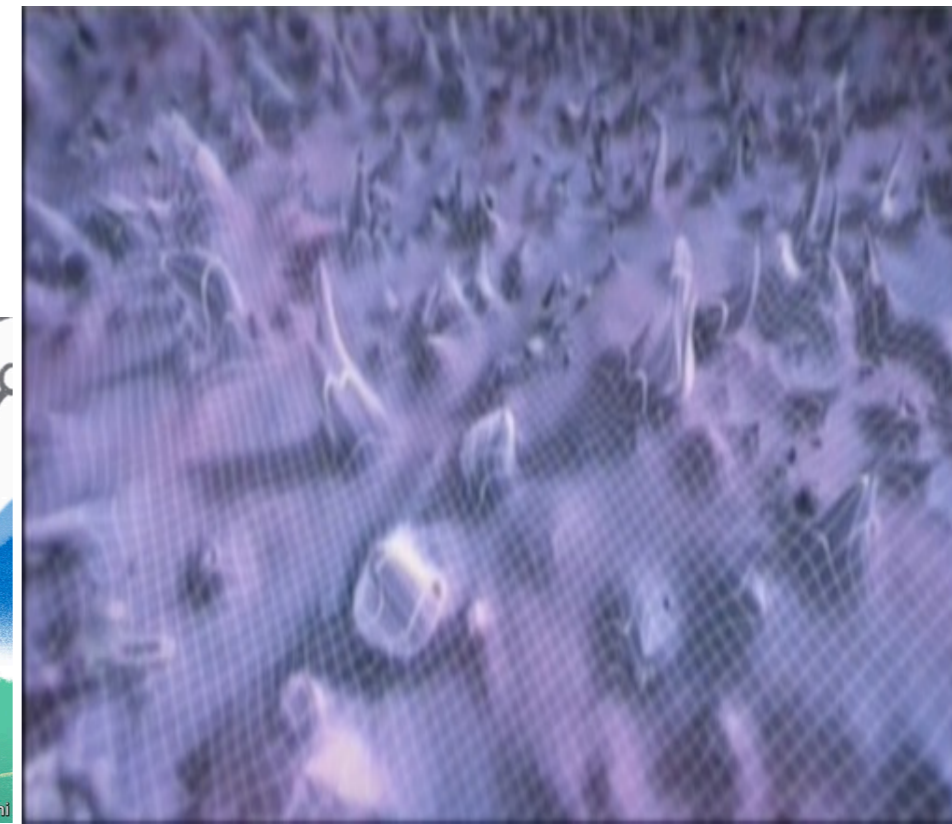
## conceptual challenges

fluctuating geometry/causal structure, entanglement, ....

thinking without fixed background spacetime/geometry

diffeo-invariance, spacetime observables, relational strategy, but fully quantum!

quantum clocks and rods

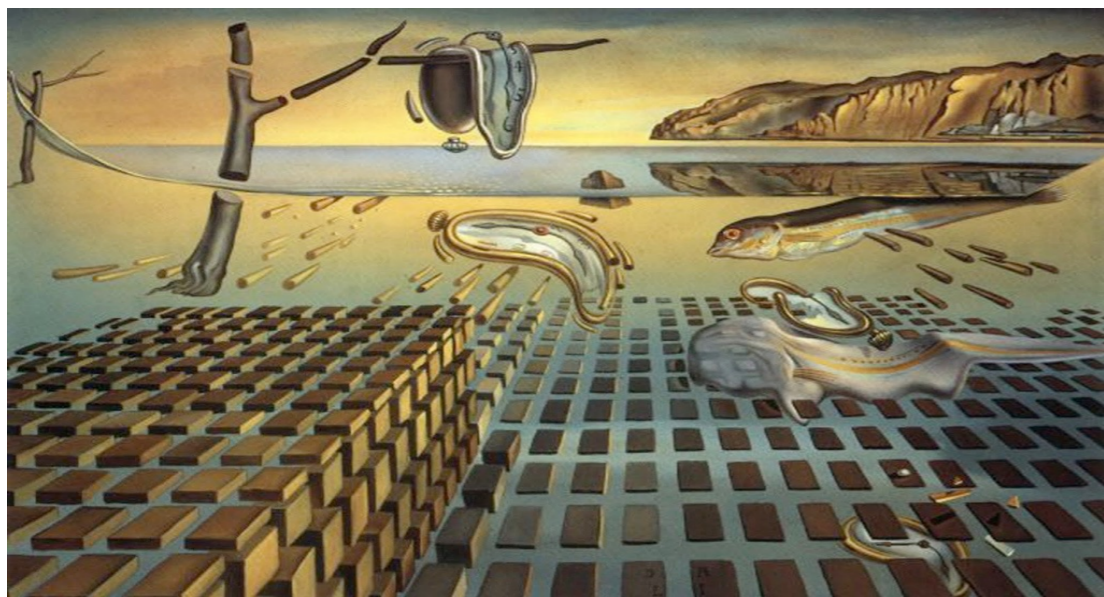


# **The emergent spacetime scenario**

# Is spacetime emergent?

suggestions that spacetime and geometry are not fundamental but emergent, collective entities

- challenges to “localization” in semi-classical GR minimal length scenarios
- spacetime singularities in GR breakdown of continuum itself?
- black hole thermodynamics space itself is a thermodynamic system
- black hole information paradox some fundamental principle has to go: locality?
- Einstein’s equations as equation of state  
GR dynamics is effective equation of state for any microscopic dofs collectively described by a spacetime, a metric and some matter fields
- entanglement ~ geometry  
geometric quantities defined by quantum (information) notions (examples from AdS/CFT, and various quantum many-body systems)
- many suggestions and results from several QG approaches (string theory, LQG, causal sets, ...)



new (quantum) dofs?

discrete structures?

which "dynamics"?

# Quantum gravity problem reloaded

quantum theory of "new" non-spatiotemporal entities

continuum spacetime and geometric quantum observables  
reconstructed from collective quantum dynamics of  
"atoms of space"



quantum spacetime as a (background-independent) quantum many-body system

extraction of spacetime and cosmology similar to typical problem in condensed matter theory  
(from atoms to macroscopic effective continuum physics)

- all GR structures and dynamics are to be approximately obtained (in relational language) at effective level
- not just emergent gravity; flat spacetime itself would be emergent, highly excited, collective state of "QG atoms"

further issues and possibilities open up in "emergent spacetime" scenarios

besides quantum effects of spacetime, we will have collective effects of "spacetime constituents"

which may manifest in new (or newly explained) spacetime features

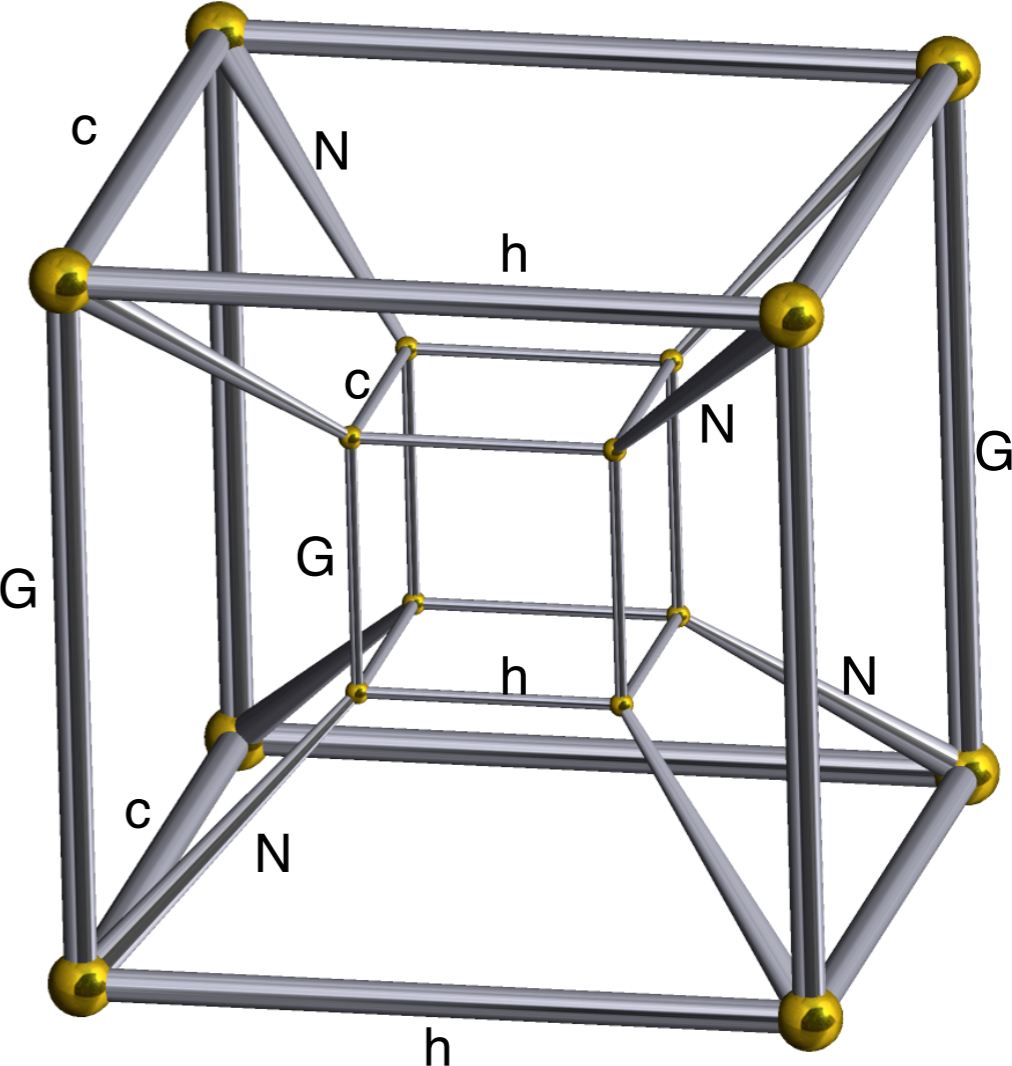
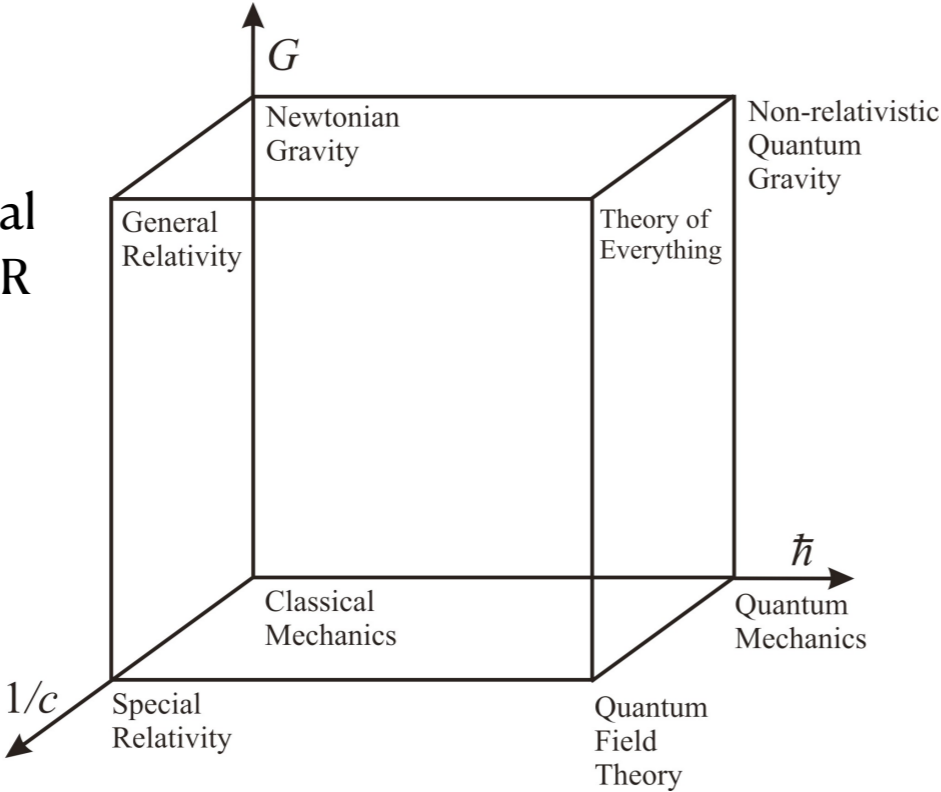
main conceptual point:

but if fundamental d.o.f.s are not smooth spacetimes (geometries) .....

the Bronstein cube .....

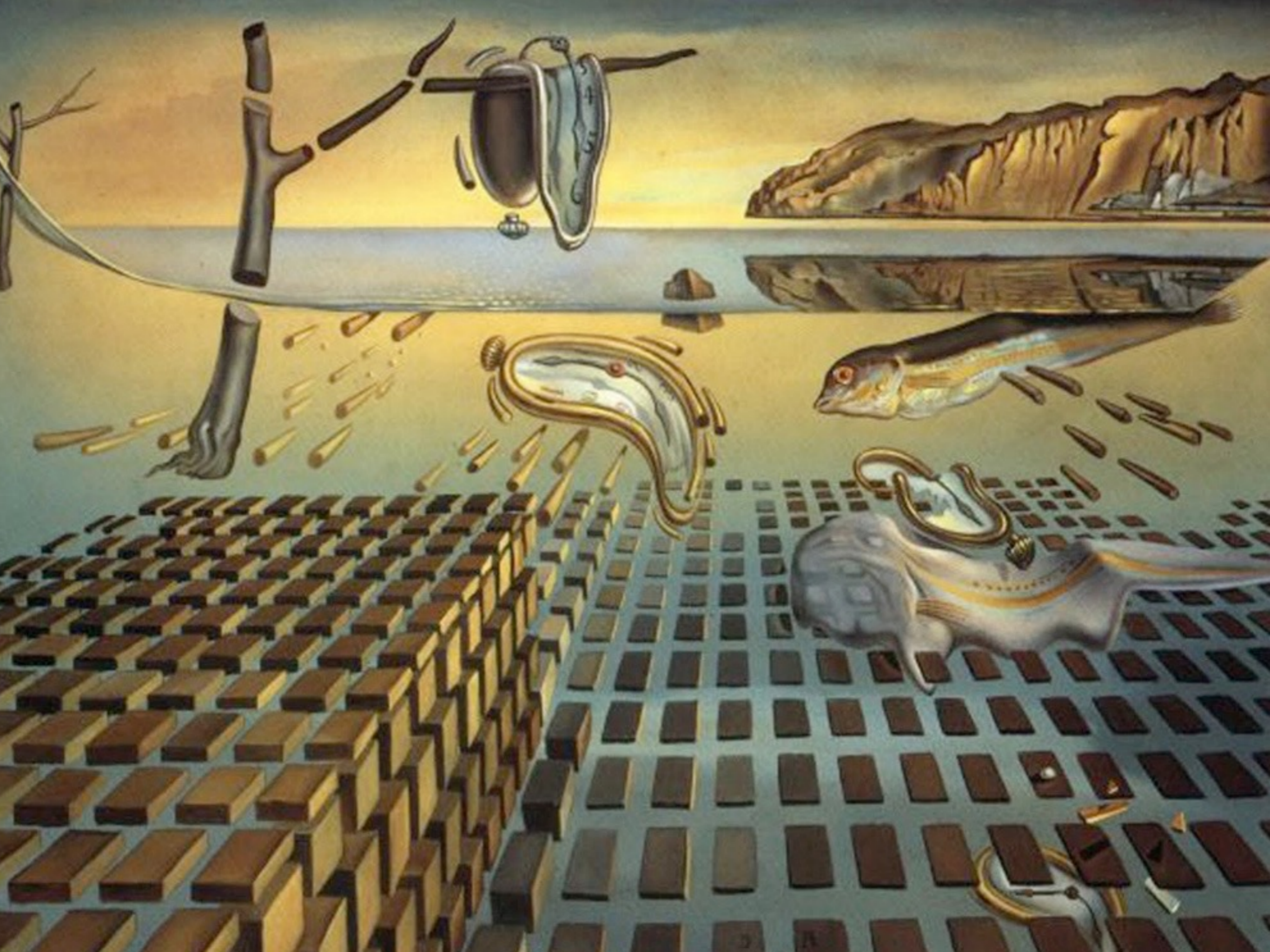
the Bronstein hypercube of Quantum Gravity

corresponds to traditional view of QG = quantum GR



adding a new direction to our understanding of the world....  
.... understanding the collective physics of many QG d.o.f.s

N-direction is where emergent behaviour takes place:  
"More is different"



# **Key point:**

**space and time may not be fundamental**

**physics may not be "fields (values) in relation to fields (values)"**

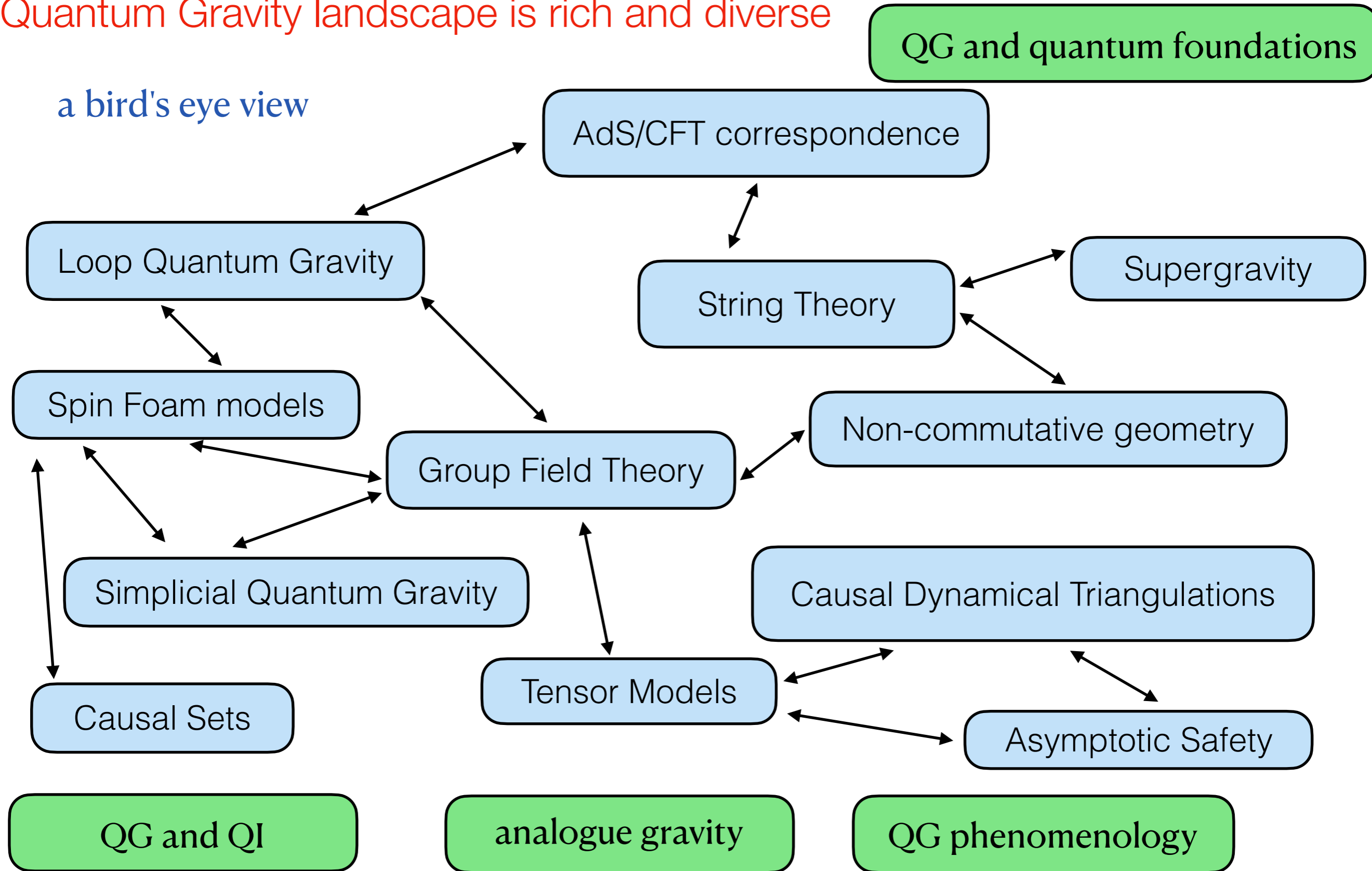
**some other structures/entities may replace continuum fields  
at more fundamental level**



**A proviso:**

**Quantum Gravity landscape  
is rich and diverse**

# Quantum Gravity landscape is rich and diverse



great variety; many mutual relations; many shared issues; mostly same goals

all approaches incomplete, missing parts (and achievements) depend on chosen strategy

# Quantum Gravity landscape is rich and diverse



several sub-communities

with sometimes difficult relationships



very differently-sized communities - strings  $\sim O(1000)$  , LQG  $\sim O(100)$  , others  $\sim O(10)$

but counting is very ambiguous, because boundaries are not sharp, and actual research directions very diverse

different historical roots of different communities:

some in particle physics tradition, others more in GR tradition; some more mathematical, others more physics-oriented

communication difficult because of different languages, and different definitions of and perspectives on QG problem

scarce resources do not help

**A proviso:**

**Quantum Gravity landscape  
is rich and diverse**

**here, just one example.....**

**Example:**

**Tensorial Group Field Theories  
for Quantum Gravity**

**(here, quantum geometric models)**

# GFTs: basics

4d case - specific class of models

atoms of space ~ quantum 3-simplices with extra scalar dofs

- geometric variables: triangle vectors ~  $\mathfrak{su}(2)$  Lie algebra elements

$$A_i n_i^I = b_i^I \in \mathbb{R}^{3,1}$$

$$b_i \cdot N = 0$$

$$\sum_i b_i = 0$$

4 triangle vectors (with modulus equal to area)

normal vector to 3d hypersurface

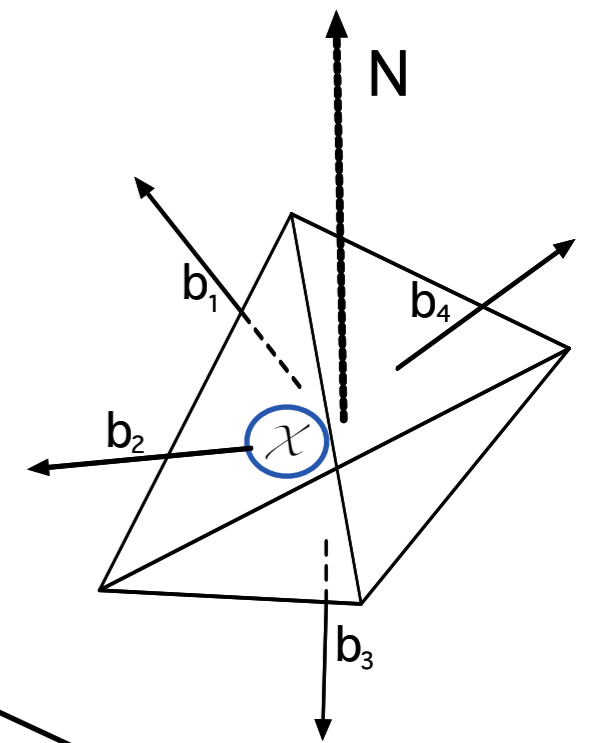
thus vectors are effectively  $b_i \in \mathbb{R}^3$

all vectors lie in same hypersurface  
(spacelike if normal is timeline)

triangle vectors close  
(triangles form closed 3-cell)

identified with Lie algebra elements  $\mathfrak{su}(2) \simeq \mathbb{R}^3$

- observables: e.g. triangle areas, volume  $A_i = |b_i|$   $V = \frac{1}{6} \sqrt{b_1 \cdot b_2 \times b_3}$  become operators:  $\vec{b}_i \rightarrow \hat{J}_i$



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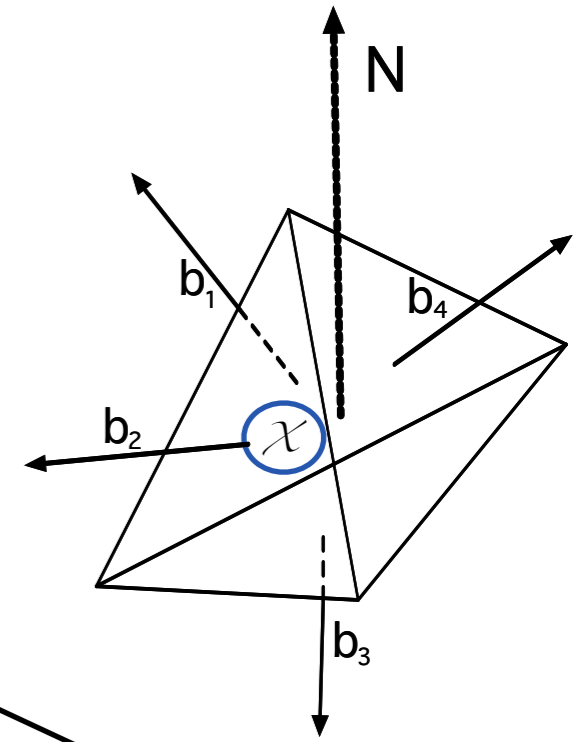
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## Hilbert space of quantum tetrahedron

(in terms of SU(2) irreps)

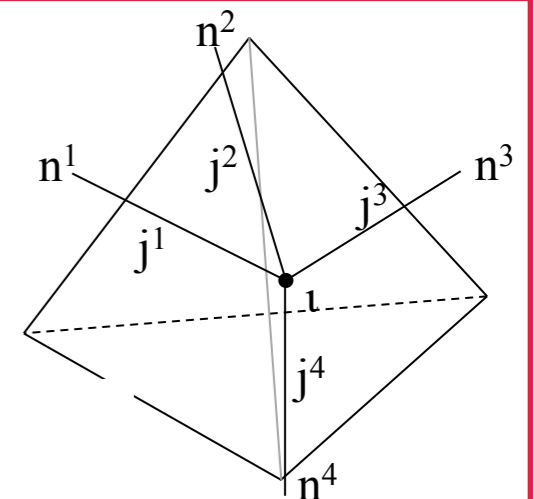
spin network vertex ~ quantum tetrahedron

$$\mathcal{H}_v = \bigoplus_{\vec{j}_v} \left( \bigotimes_{i=1}^d \underbrace{V^{j^i}_v}_{\text{repr. space}} \otimes \underbrace{\mathcal{I}^{\vec{j}_v}}_{\text{intertwiner space}} \right)$$

quantum geometric operators act on this Hilbert space:

$|j^i n^i\rangle \in V^{j^i}$  diagonalises area operator

$|\vec{j}\rangle \in \mathcal{I}^{\vec{j}} = \text{Inv}_G [V^{j^1} \otimes \dots \otimes V^{j^d}]$  diagonalises volume operator



+ scalar dofs  $\otimes L^2(\mathbb{R} \times \dots \times \mathbb{R})$

## GFTs: basics

4d case - specific class of models

- equivalent representation:  $\Psi(g_1, \dots, g_4) = \Psi(g_1 h, \dots, g_4 h) = \sum_{\{j_i, m_i; I\}} \Psi_{m_1 \dots m_4}^{j_1 \dots j_4; I} D_{m_1 n_1}^{j_1}(g_1) \dots D_{m_4 n_4}^{j_4}(g_4) C_{n_1 \dots n_4}^{j_1 \dots j_4; I}$   
thus  $L^2(SU(2)^4 / SU(2))$  (quantum geometry dofs)



# GFTs: basics

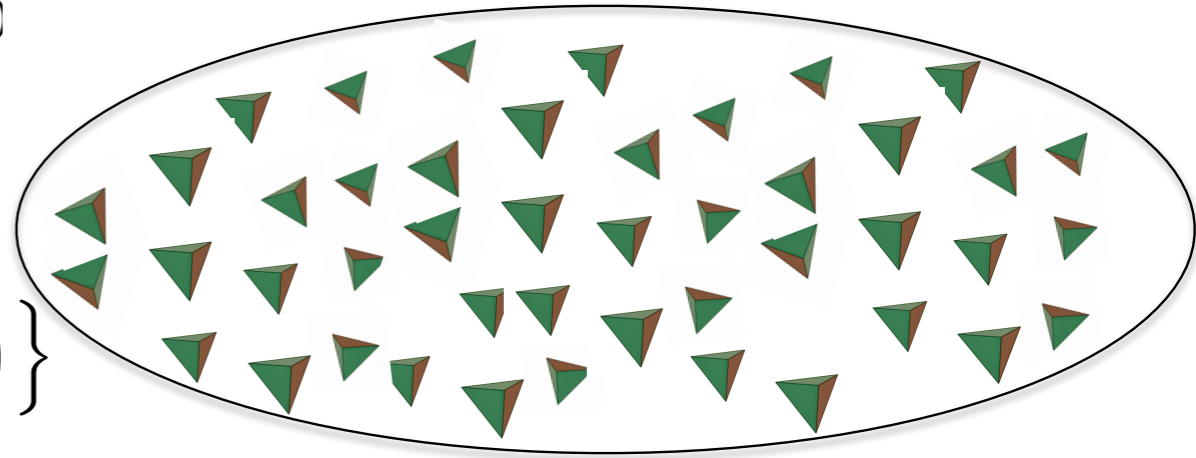
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- Fock space

$$\mathcal{F}(\mathcal{H}_v) = \bigoplus_{V=0}^{\infty} \text{sym} \left\{ \left( \mathcal{H}_v^{(1)} \otimes \mathcal{H}_v^{(2)} \otimes \dots \otimes \mathcal{H}_v^{(V)} \right) \right\}$$



- GFT field operators (creating/annihilating tetrahedra):

$$\hat{\varphi}(g_I, \chi^a) \equiv \hat{\varphi}(g_I, \chi^1, \dots, \chi^n) \quad \left[ \hat{\varphi}(\vec{g}), \hat{\varphi}^\dagger(\vec{g}') \right] = \mathbb{I}_G(\vec{g}, \vec{g}') \quad \left[ \hat{\varphi}(\vec{g}), \hat{\varphi}(\vec{g}') \right] = \left[ \hat{\varphi}^\dagger(\vec{g}), \hat{\varphi}^\dagger(\vec{g}') \right] = 0$$

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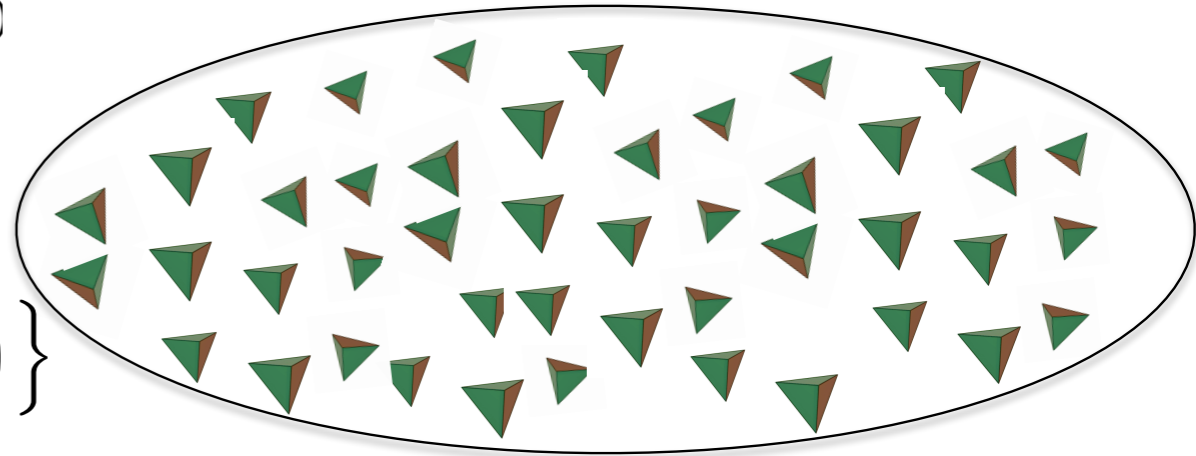
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- discrete (collective) quantum geometric observables

e.g. volume  $\hat{V}_{tot} = \int [dg_i][dg'_j] \hat{\varphi}^\dagger(g_i) V(g_i, g'_j) \hat{\varphi}(g'_j) = \sum_{J_i} \hat{\varphi}^\dagger(J_i) V(J_i) \hat{\varphi}(J_j)$

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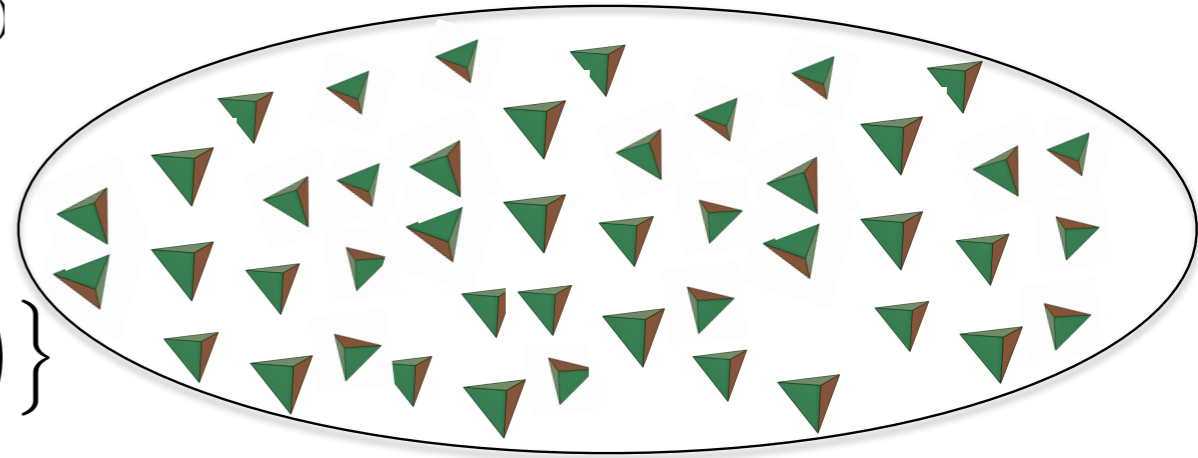
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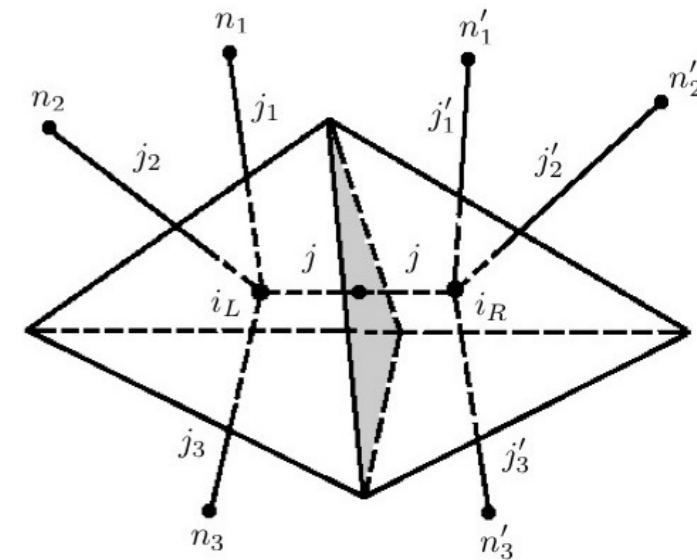
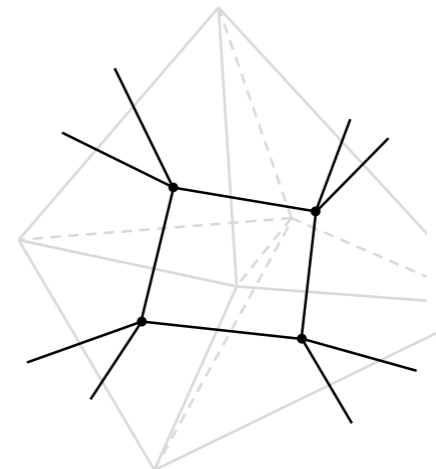
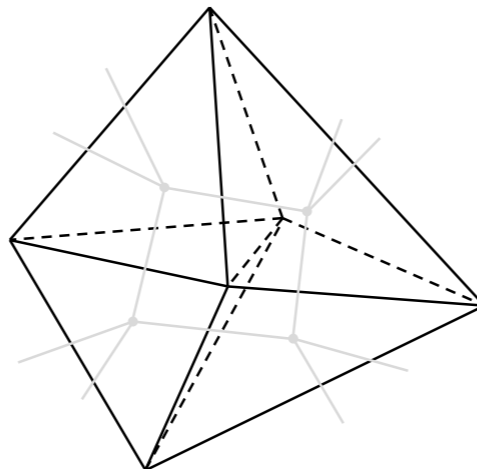
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- maximal entanglement of "triangle dofs" ~ gluing of tetrahedra across triangle

entangled states ~ extended simplicial complexes

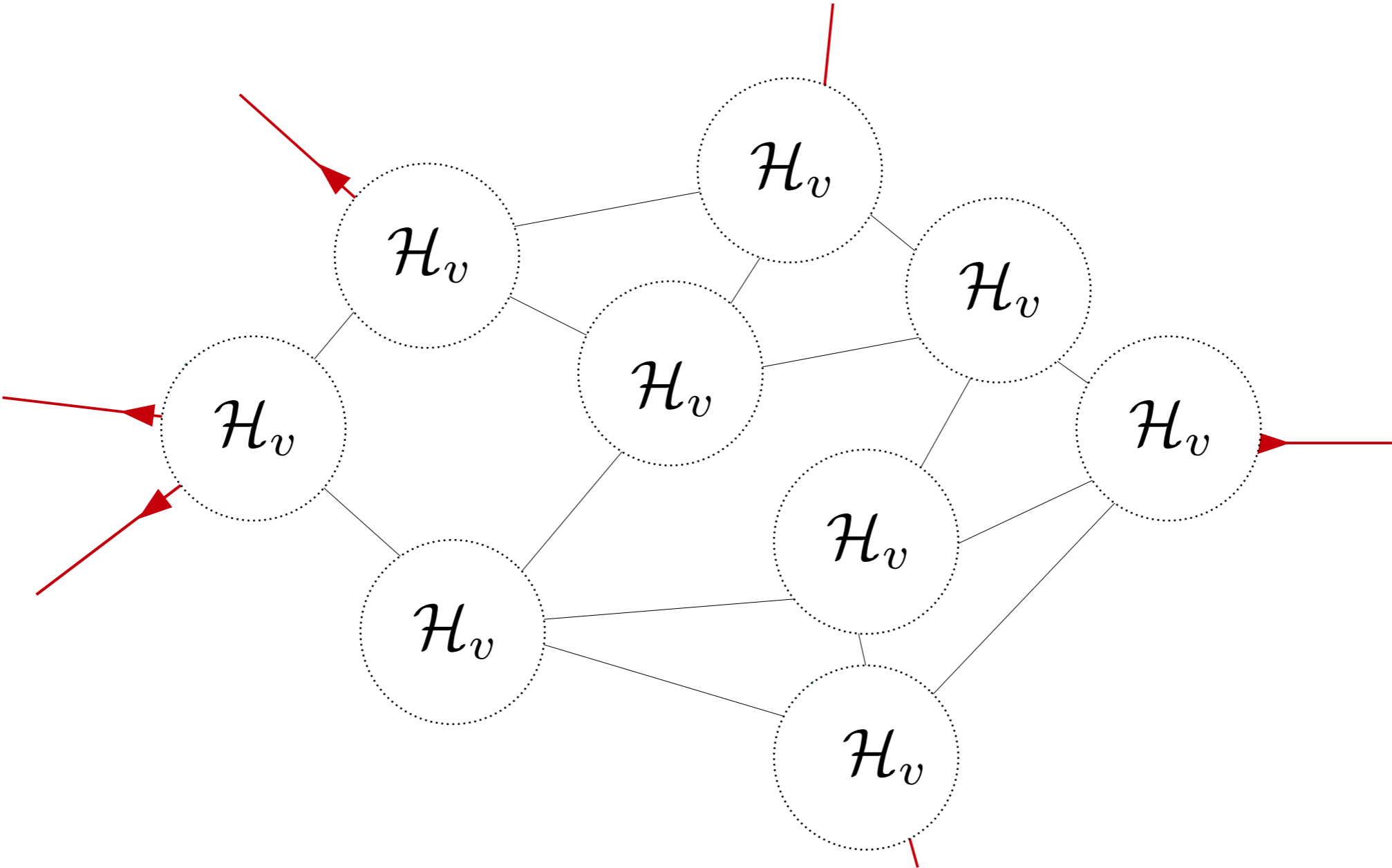


# QG states = entanglement networks of quantum geometric blocks

algebraic data on graph

elementary quantum systems on nodes

graph ~ pattern of entanglement across nodes



structure shared by several QG formalisms (LQG, spin foams, lattice QG, TGFT)

# GFTs: basics

4d case - specific class of models

dynamics of quantum atomic geometry

GFT action = prescription for weights associated to building blocks of 4d lattice in sum over discrete geometries

$$S(\varphi, \bar{\varphi}) = \frac{1}{2} \int [dg_i] \overline{\varphi(g_i)} \mathcal{K}(g_i) \varphi(g_i) + \frac{\lambda}{D!} \int [dg_{ia}] \varphi(g_{i1}) \dots \varphi(g_{iD}) \mathcal{V}(g_{ia}, \bar{g}_{iD}) + c.c.$$
$$\mathcal{Z} = \int \mathcal{D}\varphi \mathcal{D}\bar{\varphi} e^{i S_\lambda(\varphi, \bar{\varphi})} = \sum_{\Gamma} \frac{\lambda^{N_\Gamma}}{\text{sym}(\Gamma)} \mathcal{A}_\Gamma$$

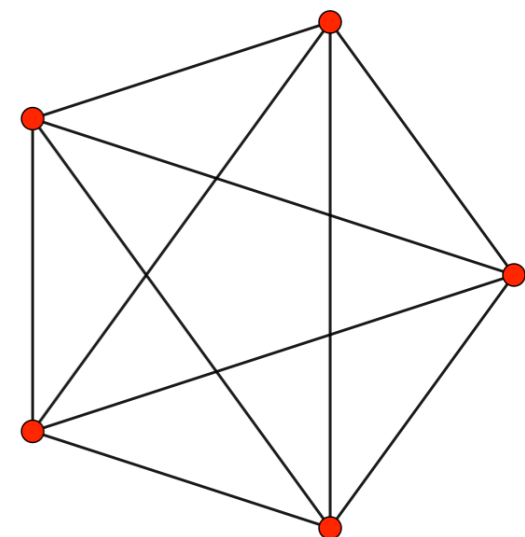
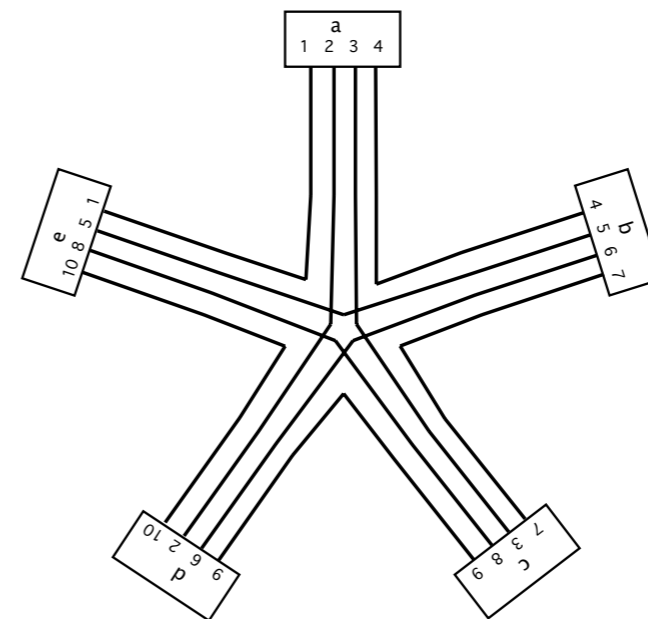
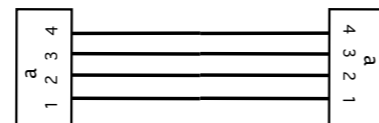
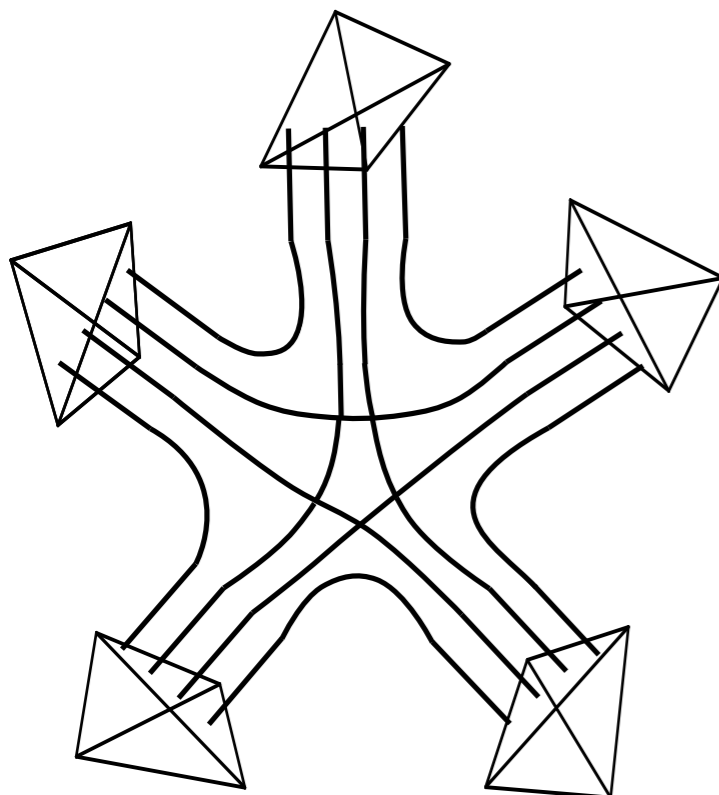
Feynman diagrams = stranded diagrams dual to cellular complexes of arbitrary topology

De Pietri, Petronio, '00; R. Gurau, '10; ...

labelled by group-theoretic data (group elements, group irreps, ...)

example: 3-simplices/4-tensors (4d)

generalises to any dimension (rank of tensor)



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Feynman amplitudes (model-dependent) = convolution of propagation kernels with interaction kernels = sum over group-theoretic data (group elements, Lie algebra elements, group irreps) associated to complex dual to Feynman diagram

Reisenberger, Rovelli, '00

- GFT Feynman amplitudes = lattice gravity path integrals (in group/algebra variables) on lattice dual to GFT Feynman diagram = spin foam models (in irreps variables)

A. Baratin, DO, '11

M. Finocchiaro, DO, '18

basic guideline for model-building (choosing GFT action):

GFT Feynman amplitudes = simplicial path integrals for gravity coupled to scalar fields

# Quantum "transition amplitudes" for QG processes

- quantum dynamics: assignment of quantum amplitude to each possible process (directed graph ~ cellular complex)

- full amplitude obtained from elementary operators (kernels):

$$\mathcal{V}_a : \bigotimes_{p \in \partial a} \mathcal{H}_p \longrightarrow \mathbb{C} \quad \text{node operator}$$

$$\mathcal{K}_e : \mathcal{H}_{p_1} \otimes \mathcal{H}_{p_2} \longrightarrow \mathbb{C} \quad \text{gluing operator}$$

with appropriate dualization reflecting orientation

- amplitude associated to whole complex):

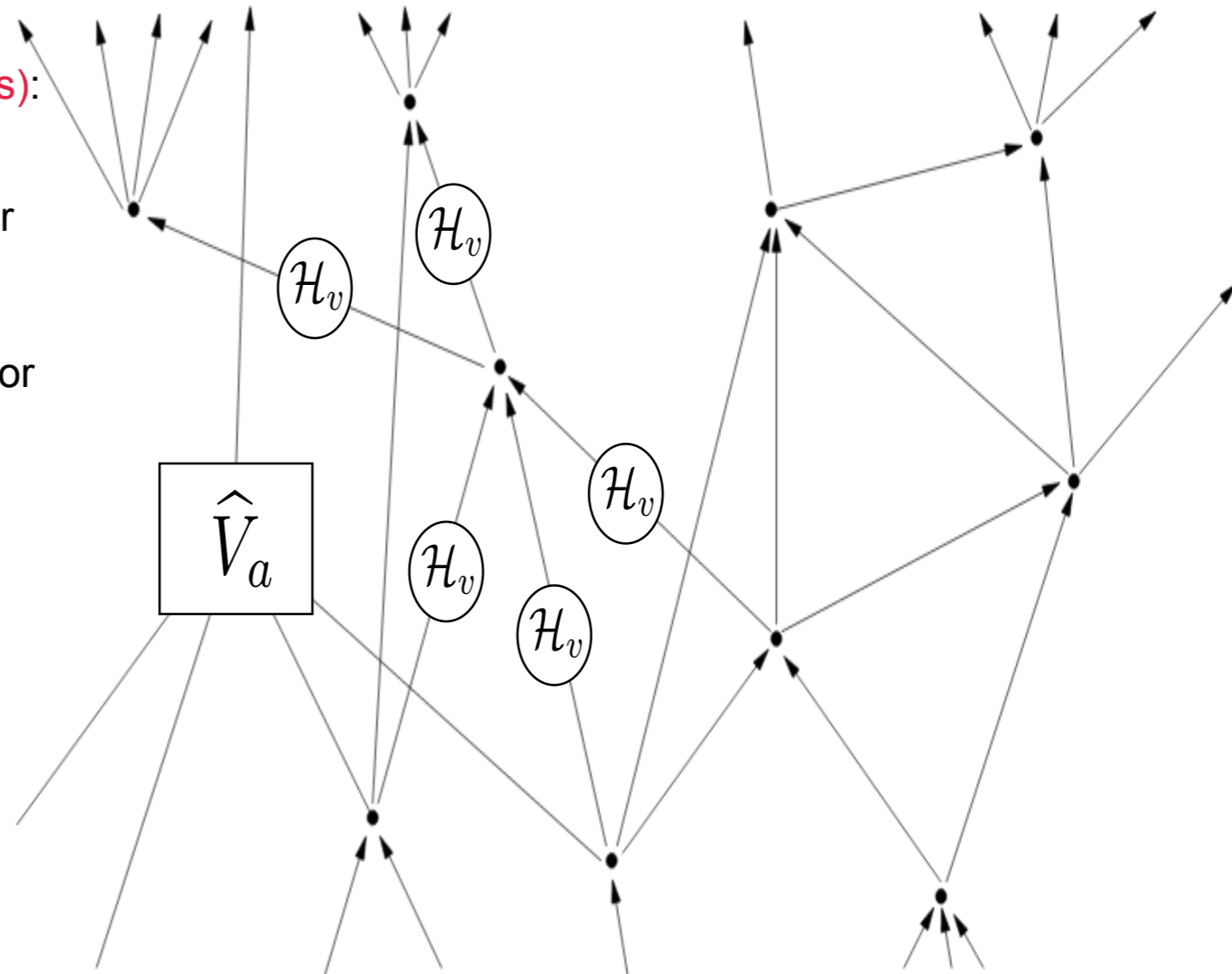
$$\mathcal{A}(\mathfrak{m}) = \text{Tr}_{p \in \mathfrak{m}} \left( \prod_{e|m} \mathcal{K}_e \prod_{a \in \mathfrak{m}} \mathcal{V}_a \right)$$

(trace defined over complete basis in link Hilbert spaces)

- different quantum gravity models (spin foam, LQG, lattice QG, TGFT) = different choices of elementary operators (and Hilbert spaces)

note: classical approx on given lattice ---> sum over saddles of discrete gravity path integral = good simplicial geometries on given lattice

$$\mathcal{A}(\mathfrak{m}) \sim \sum_{g_\Delta = \text{classical}} e^{iS(g_\Delta)}$$



## GFTs: example

Boulatov model - topological 3d euclidean QG (no matter)

$$\varphi : SU(2)^{\times 3} \rightarrow \mathbb{C} \quad \sim \text{quantum triangles}$$

$$S(\varphi) = \frac{1}{2} \int [dg] \varphi^2(g_1, g_2, g_3) + \frac{\lambda}{4!} \int [dg] \varphi(g_1, g_2, g_3) \varphi(g_3, g_4, g_5) \varphi(g_5, g_2, g_6) \varphi(g_6, g_4, g_1) + \text{cc}$$

for fields satisfying:

$$\varphi(g_1, g_2, g_3) = \varphi(hg_1, hg_2, hg_3) \quad \forall h \in SU(2)$$

partition function & perturbative expansion

$$\mathcal{Z} = \int \mathcal{D}\varphi \mathcal{D}\bar{\varphi} e^{i S_\lambda(\varphi, \bar{\varphi})} = \sum_{\Gamma} \frac{\lambda^{N_\Gamma}}{\text{sym}(\Gamma)} \mathcal{A}_\Gamma$$

Feynman diagrams dual to 3d simplicial lattices



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Feynman amplitudes in different representations:

$$\begin{aligned} \mathcal{A}_\Gamma &= \int \prod_l dh_l \prod_f \delta(H_f(h_l)) = \int \prod_l dh_l \prod_f \delta\left(\overrightarrow{\prod_{l \in \partial f}} h_l\right) = \\ &= \sum_{\{j_e\}} \prod_e d_{j_e} \prod_\tau \left\{ \begin{array}{ccc} j_1^\tau & j_2^\tau & j_3^\tau \\ j_4^\tau & j_5^\tau & j_6^\tau \end{array} \right\} = \int \prod_l [dh_l] \prod_e [d^3 x_e] e^{i \sum_e \text{Tr } x_e H_e} \end{aligned}$$

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lattice gauge theory formulation of 3d gravity/BF theory

spin foam formulation of 3d gravity

i.e. quantum covariant dynamics of spin networks (LQG)

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Feynman diagrams dual to 3d simplicial lattices

Feynman amplitudes in different representations:

$$\begin{aligned} \mathcal{A}_\Gamma &= \int \prod_l dh_l \prod_f \delta(H_f(h_l)) = \int \prod_l dh_l \prod_f \delta\left(\prod_{l \in \partial f} h_l\right) = \\ &= \sum_{\{j_e\}} \prod_e d_{j_e} \prod_\tau \left\{ \begin{matrix} j_1^\tau & j_2^\tau & j_3^\tau \\ j_4^\tau & j_5^\tau & j_6^\tau \end{matrix} \right\} = \int \prod_l [dh_l] \prod_e [d^3 x_e] e^{i \sum_e \text{Tr } x_e H_e} \end{aligned}$$

lattice gauge theory formulation of 3d gravity/BF theory

spin foam formulation of 3d gravity

discrete 1st order path integral for 3d gravity on simplicial complex dual to GFT Feynman diagram

i.e. quantum covariant dynamics of spin networks (LQG)

# GFTs: example

## Boulatov model - topological 3d euclidean QG (no matter)

$$\varphi : SU(2)^{\times 3} \rightarrow \mathbb{C} \quad \sim \text{quantum triangles}$$

$$S(\varphi) = \frac{1}{2} \int [dg] \varphi^2(g_1, g_2, g_3) + \frac{\lambda}{4!} \int [dg] \varphi(g_1, g_2, g_3) \varphi(g_3, g_4, g_5) \varphi(g_5, g_2, g_6) \varphi(g_6, g_4, g_1) + \text{cc}$$

for fields satisfying:

$$\varphi(g_1, g_2, g_3) = \varphi(hg_1, hg_2, hg_3) \quad \forall h \in SU(2)$$

partition function & perturbative expansion

$$\mathcal{Z} = \int \mathcal{D}\varphi \mathcal{D}\bar{\varphi} e^{i S_\lambda(\varphi, \bar{\varphi})} = \sum_{\Gamma} \frac{\lambda^{N_\Gamma}}{\text{sym}(\Gamma)} \mathcal{A}_\Gamma$$

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discretization of Palatini gravity:

$$S(e, \omega) = \int \text{Tr}(e \wedge F(\omega))$$

# GFT example - 4d Lorentzian QG

(motivated by: LQG + simplicial quantum geometry)

**EPRL model**  $S = \sum_{\substack{j_{v_a i} \\ m_{v_a i}, \ell_a}} \bar{\varphi}_{m_{v_1}}^{j_{v_1} \ell_1} \varphi_{m_2}^{j_{v_2} \ell_2} (\mathcal{K}_2)_{m_{v_1}, m_{v_2}}^{j_{v_1} j_{v_2} \ell_1 \ell_2} + V$

$$V = \sum_{j_i, m_i, \ell_i} \left[ \varphi_{m_1 m_2 m_3 m_4}^{j_1 j_2 j_3 j_4 \ell_1} \varphi_{m_4 m_5 m_6 m_7}^{j_4 j_5 j_6 j_7 \ell_2} \varphi_{m_7 m_3 m_8 m_9}^{j_7 j_3 j_8 j_9 \ell_3} \varphi_{m_9 m_6 m_2 m_{10}}^{j_9 j_6 j_2 j_{10} \ell_4} \varphi_{m_{10} m_8 m_5 m_1}^{j_{10} j_8 j_5 j_1 \ell_5} \times \tilde{\mathcal{V}}_5(j_1, \dots, j_{10}; \ell_1, \dots, \ell_5) \right]$$

$$\tilde{\mathcal{V}}_5(j_{ab}, i_a) = \sum_{n_a} \int d\rho_a (n_a^2 + \rho_a^2) \left( \bigotimes_a f_{n_a \rho_a}^{i_a}(j_{ab}) \right) 15j_{SL(2, \mathbb{C})}((2j_{ab}, 2j_{ab}\gamma); (n_a, \rho_a))$$

$$f_{n\rho}^i := i^{m_1 \dots m_4} \bar{C}_{(j_1, m_1) \dots (j_4, m_4)}^{n\rho} \quad \rho = \gamma n \quad n = 2j$$

based on SU(2) group and irreps - relation between SL(2,C) and SU(2) data; (almost) SU(2) spin network states

Feynman amplitudes:

no need here

specific form of action implements:

conditions for well-defined simplicial quantum geometry of GFT quanta (3-simplices)

conditions for producing 4d lattices with proper simplicial quantum geometry in perturbative expansion

**Example:**

**TGFT cosmology**

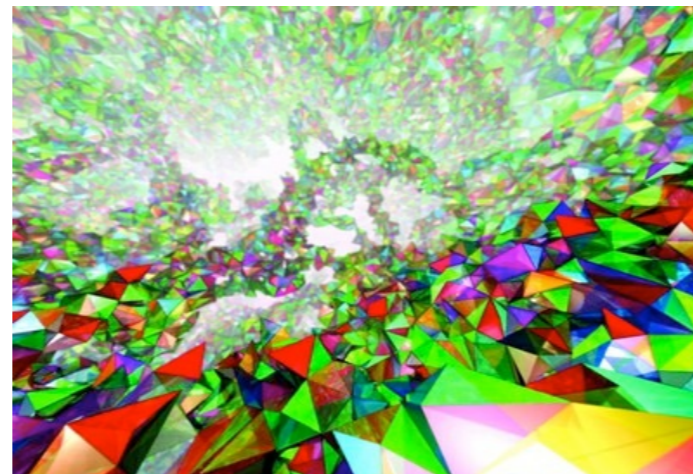
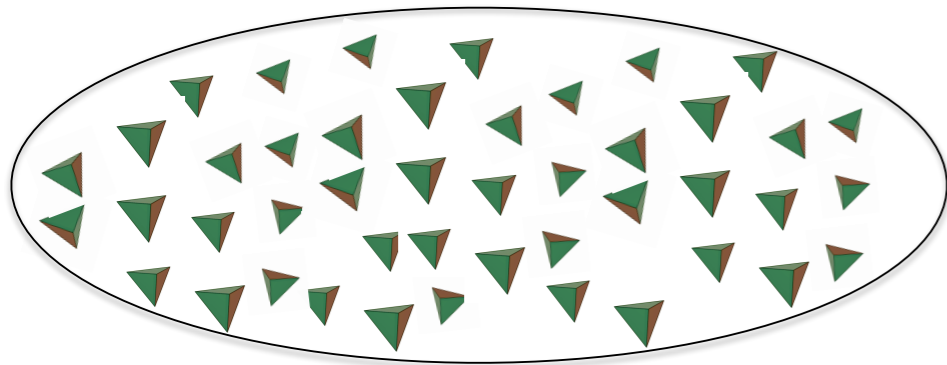
**emergent spacetime physics from QG**

# spacetime and geometry are emergent in GFT

from perspective of fundamental QG atoms of space:

continuum geometry = coarse-grained description of discrete geometry of many (infinite) QG atoms

GR dynamics = approximate description of collective quantum dynamics of many (infinite) QG atoms





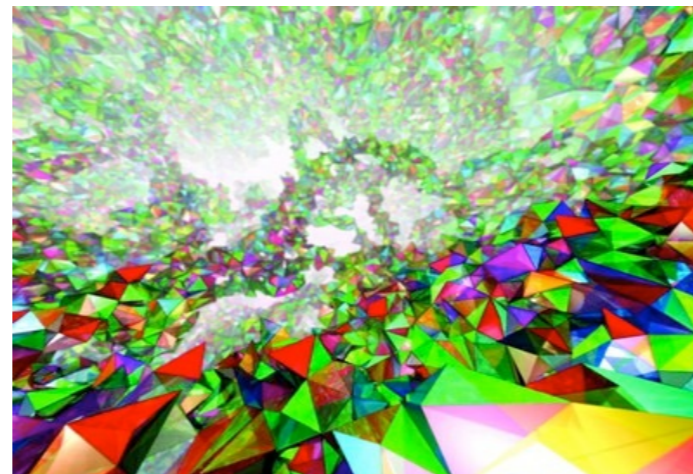
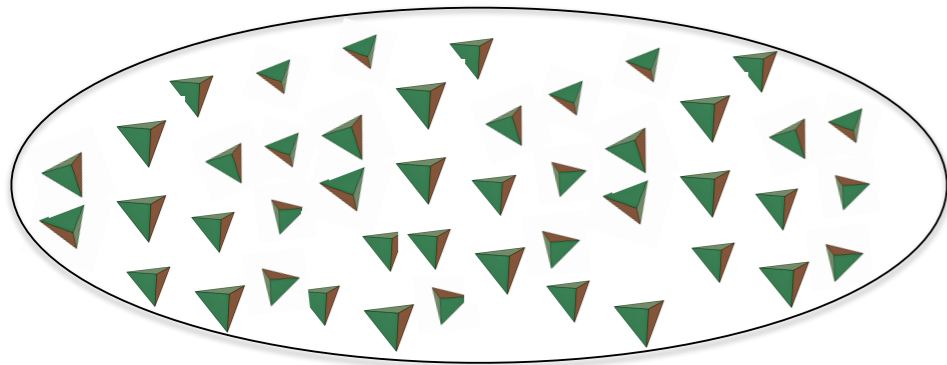
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cosmology expected to correspond to "most coarse-grained" dynamics

→ in other words: effective dynamics of special (global) observables of full theory

→ QG hydrodynamics

## GFT cosmology

- **general strategy:**
  - hypothesis: universe as QG quantum fluid (condensate)
  - extract approximate hydrodynamic eqns for QG fluid (density and phase)
  - compute relational cosmological observables in hydrodynamic approximation, as functions of density & phase
  - translate hydrodynamic eqns into eqns for cosmological observables

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$$S(\varphi, \bar{\varphi}) = \frac{1}{2} \int [dg_i] \overline{\varphi(g_i)} \mathcal{K}(g_i) \varphi(g_i) + \frac{\lambda}{D!} \int [dg_{ia}] \varphi(g_{i1}) \dots \varphi(\bar{g}_{iD}) \mathcal{V}(g_{ia}, \bar{g}_{iD}) + c.c.$$

$$\mathcal{Z} = \int \mathcal{D}\varphi \mathcal{D}\bar{\varphi} e^{i S_\lambda(\varphi, \bar{\varphi})} = \sum_{\Gamma} \frac{\lambda^{N_\Gamma}}{\text{sym}(\Gamma)} \mathcal{A}_\Gamma$$

$$F_\lambda(J) = \ln Z_\lambda[J] \quad \Gamma[\phi] = \text{sup}_J (J \cdot \phi - F(J)) \quad \langle \varphi \rangle = \phi$$

\* simplest approximation:  
mean field hydrodynamics

$$\Gamma[\phi] \approx S_\lambda(\phi)$$

mean field ~ condensate wavefunction

- **corresponding quantum states:**

(simplest): GFT condensate, GFT field coherent state

$$|\sigma\rangle := \exp(\hat{\sigma}) |0\rangle$$

$$\hat{\sigma} := \int d^4g \sigma(g_I) \hat{\varphi}^\dagger(g_I) \quad \sigma(g_I k) = \sigma(g_I)$$

# GFT cosmology

## general facts

- cosmological interpretation natural and clear:

isomorphism between domain of TGFT condensate wavefunction and minisuperspace

$$\begin{aligned} \sigma(\mathcal{D}) \quad \mathcal{D} &\simeq \{ \text{geometries of tetrahedron} \} \simeq \\ &\simeq \{ \text{continuum spatial geometries at a point} \} \simeq \\ &\simeq \text{minisuperspace of homogeneous geometries} \end{aligned}$$

S. Gielen, DO, L. Sindoni, '13

- general form of resulting (Gross-Pitaevskii) equations of motion for condensate wavefunction (mean field):

$$\int [dg'] d\chi' \mathcal{K}(g, \chi; g', \chi') \sigma(g', \chi') + \lambda \frac{\delta}{\delta \varphi} \mathcal{V}(\varphi) |_{\varphi \equiv \sigma} = 0$$

Gielen, DO, Sindoni, '13; DO, Sindoni, Wilson-Ewing, '16

← polynomial functional of condensate wavefunction

cosmology as QG hydrodynamics ~ non-linear extension of (loop) quantum cosmology

that is, in isotropic restriction and with just one matter field:

$$\sigma(a, \phi)$$

"wavefunction" on minisuperspace

$$\mathcal{K}(a, \partial_a, \phi, \partial_\phi) \sigma(a, \phi) + \mathcal{V}[\sigma(a, \phi)] = 0 \quad \text{hydrodynamic (non-linear, possibly non-local) eqn on minisuperspace}$$

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polynomial functional of condensate wavefunction

cosmology as QG hydrodynamics ~ non-linear extension of (loop) quantum cosmology

- cosmological observables are fluid averages = mean values of fundamental QG operators in Fock space

relationally localized in time/space as functions of values of physical (e.g. scalar matter) dofs, specified by the GFT state (for GFT models including such dofs) e.g. volume operator

$$\hat{V} = \int d^n \chi \int dg_I dg'_I \hat{\varphi}^\dagger(g_I, \chi^a) V(g_I, g'_I) \hat{\varphi}(g'_I, \chi^a) \longrightarrow V(x^0, x^i) = \langle \sigma_{(x^0, x^i)} | \hat{V} | \sigma_{(x^0, x^i)} \rangle | \sigma_{(x^0, x^i)}$$

- eqn for condensate wavefunction -----> eqn for geometric/cosmological observables

# GFT cosmology

concrete example of cosmology from "quantum geometric" GFT models

valid for EPRL & BC models, possibly more

general mean field eqns for quantum geometry coupled to 5 scalar fields in peaked states

general form of dynamics - work with parametrized ambiguities

$$S_{\text{GFT}} = K + U + U^*$$

$$K = \int dg_I dh_I \int d^d \chi d^d \chi' d\phi d\phi' \bar{\varphi}(g_I, \chi) \mathcal{K}(g_I, h_I; (\chi - \chi')_\lambda^2, (\phi - \phi')^2) \varphi(h_I, (\chi')^\mu, \phi')$$

$$U = \int d^d \chi d\phi \int \left( \prod_{a=1}^5 dg_I^a \right) \mathcal{U}(g_I^1, \dots, g_I^5) \prod_{\ell=1}^5 \varphi(g_I^\ell, \chi^\mu, \phi)$$

simple mean field approx. - classical GFT eqns

S. Gielen, DO, L. Sindoni, '13

$$\left\langle \frac{\delta S_{\text{GFT}}[\hat{\varphi}, \hat{\varphi}^\dagger]}{\delta \hat{\varphi}^\dagger(g_I, \chi_0)} \right\rangle_{\sigma_{\epsilon\mu}; x^\mu, \pi_\mu} \equiv \left\langle \sigma_{\epsilon\mu}; x^\mu, \pi_\mu \left| \frac{\delta S_{\text{GFT}}[\hat{\varphi}, \hat{\varphi}^\dagger]}{\delta \hat{\varphi}^\dagger(g_I, \chi_0)} \right| \sigma_{\epsilon\mu}; x^\mu, \pi_\mu \right\rangle = 0$$

restriction to "good clock+rods" condensate states - peakedness properties on clock/rod values

$$\sigma_{\epsilon, \delta, \pi_0, \pi_x; x^\mu}(g_I, \chi^\mu, \phi) = \eta_\epsilon(\chi^0 - x^0; \pi_0) \eta_\delta(|\boldsymbol{\chi} - \mathbf{x}|; \pi_x) \tilde{\sigma}(g_I, \chi^\mu, \phi)$$

L. Marchetti, DO, '20, '21

$$|\boldsymbol{\chi} - \mathbf{x}|^2 = \sum_{i=1}^d (\chi^i - x^i)^2 \quad \mathbb{C} \ni \delta = \delta_r + i\delta_i \quad \delta_r > 0 \quad \epsilon, |\delta| \ll 1 \quad z_0 \equiv \epsilon\pi_0^2/2 \quad z \equiv \delta\pi_x^2/2$$

# GFT cosmology

## Observables and their relational (mean) values

- number operator

$$\hat{N} = \int d^n \chi \int dg_I \hat{\varphi}^\dagger(g_I, \chi^a) \hat{\varphi}(g_I, \chi^a)$$

- universe volume

$$\hat{V} = \int d^n \chi \int dg_I dg'_I \hat{\varphi}^\dagger(g_I, \chi^a) V(g_I, g'_I) \hat{\varphi}(g'_I, \chi^a)$$

- value of clock/rods scalar fields

$$\hat{X}^b \equiv \int d^n \chi \int dg_I \chi^b \hat{\varphi}^\dagger(g_I, \chi^a) \hat{\varphi}(g_I, \chi^a)$$

- momentum of clock/rods scalar fields

$$\hat{\Pi}_b = \frac{1}{i} \int d^n \chi \int dg_I \left[ \hat{\varphi}^\dagger(g_I, \chi^a) \left( \frac{\partial}{\partial \chi^b} \hat{\varphi}(g_I, \chi^a) \right) \right]$$

- value of matter scalar field

$$\hat{\Phi} = \frac{1}{i} \int dg_I \int d^4 \chi \int d\pi_\phi \hat{\varphi}^\dagger(g_I, \chi^\mu, \pi_\phi) \partial_{\pi_\phi} \hat{\varphi}(g_I, \chi^\mu, \pi_\phi)$$

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$$\hat{\Pi}_\phi = \int dg_I \int d^4 \chi \int d\pi_\phi \pi_\phi \hat{\varphi}^\dagger(g_I, \chi^\mu, \pi_\phi) \hat{\varphi}(g_I, \chi^\mu, \pi_\phi)$$

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- universe volume

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operators defined in full QG theory

used to define collective relational (**spacetime localized**) observables for effective continuum dynamics

as expectation values in "good clock+rods" condensate states

$$N(x^0, x^i) \equiv \langle \sigma_{\epsilon, \delta, \pi_0, \pi_x, x^\mu} | \hat{N} | \sigma_{\epsilon, \delta, \pi_0, \pi_x, x^\mu} \rangle$$

$$V(x^0, x^i) \equiv \langle \sigma_{\epsilon, \delta, \pi_0, \pi_x, x^\mu} | \hat{V} | \sigma_{\epsilon, \delta, \pi_0, \pi_x, x^\mu} \rangle$$

$$X^\mu(x^0, x^i) \equiv \langle \sigma_{\epsilon, \delta, \pi_0, \pi_x, x^\mu} | \hat{V} | \sigma_{\epsilon, \delta, \pi_0, \pi_x, x^\mu} \rangle \simeq x^\mu$$

$$\Pi(x^0, x^i) \equiv \langle \sigma_{\epsilon, \delta, \pi_0, \pi_x, x^\mu} | \hat{\Pi}_\nu | \sigma_{\epsilon, \delta, \pi_0, \pi_x, x^\mu} \rangle$$

$$\phi(x^0, x^i) \equiv \langle \sigma_{\epsilon, \delta, \pi_0, \pi_x, x^\mu} | \hat{\Phi} | \sigma_{\epsilon, \delta, \pi_0, \pi_x, x^\mu} \rangle$$

$$\Pi_\phi(x^0, x^i) \equiv \langle \sigma_{\epsilon, \delta, \pi_0, \pi_x, x^\mu} | \hat{\Pi}_\phi | \sigma_{\epsilon, \delta, \pi_0, \pi_x, x^\mu} \rangle$$



concrete example of cosmology from "quantum geometric" TGFT models valid for EPRL & BC models, possibly more

hydrodynamics eqns for cosmological observables (with some assumptions on states + approximations)

using:  $\tilde{\sigma}_j \equiv \rho_j \exp[i\theta_j]$  rewrite in standard hydrodynamic form (fluid density, phase)

homogeneous background + inhomogeneous perturbations (spacetime localization defined in relational terms)

$$\rho_j = \bar{\rho}_j + \delta\rho_j \quad \theta_j \equiv \bar{\theta}_j + \delta\theta_j \quad \bar{\rho} = \bar{\rho}(x^0, \pi_\phi) \quad \bar{\theta} = \bar{\theta}(x^0, \pi_\phi)$$

can also extract

effective dynamics for scalar cosmological perturbations

L. Marchetti, DO, '22; A. Jercher, L. Marchetti, A. Pithis, '23;  
R. Dehnil, S. Liberati, DO, to appear

can be recast in standard local QFT language)

n.b. localization is relational - non-trivial spatial dependence comes from non-trivial dependence of mean field perturbations on the relational rods

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background volume dynamics:

L. Marchetti, DO, '21    A. Jercher, DO, A. Pithis, 21

$$\left(\frac{V'}{3V}\right)^2 \simeq \left(\frac{2 \sum_j \int d\pi_\phi V_j \operatorname{sgn}(\rho') \rho_j \sqrt{\mathcal{E}_j - Q_j^2/\rho_j^2 + \mu_j^2 \rho_j^2}}{3 \sum_j \int d\pi_\phi V_j \rho_j^2}\right)^2 \quad \frac{V''}{V} \simeq \frac{2 \sum_j \int d\pi_\phi V_j [\mathcal{E}_j + 2\mu_j^2 \rho_j^2]}{\sum_j \int d\pi_\phi V_j \rho_j^2}$$

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DO, L. Sindoni, E. Wilson-Ewing, '16;  
L. Marchetti, DO, '20, '21

- very early times: very small volume - QG interactions subdominant

for large class of states:

$$\exists j / \rho_j(\chi) \neq 0 \forall \chi \longrightarrow$$

$$V = \sum_j V_j \rho_j^2$$

remains positive at all times  
(with single turning point)

and fluctuations remain under control)

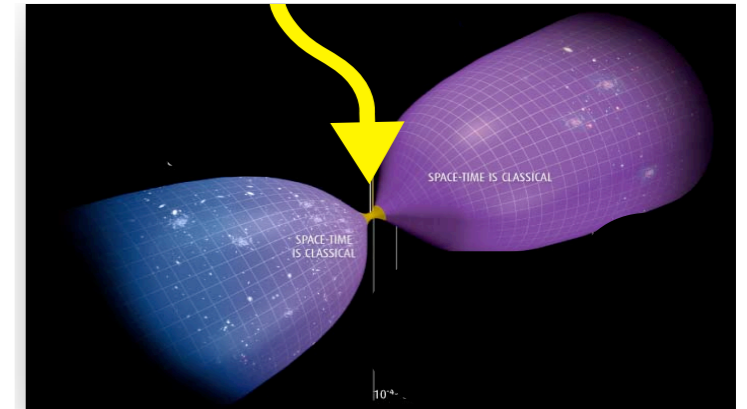
- intermediate times: large volume - QG interactions still subdominant

under some (rather mild) conditions  
on parameters of GFT model

(here written neglecting matter contribution)

$$\left(\frac{V'}{V}\right)^2 = \frac{V''}{V} = 12\pi\tilde{G}$$

quantum bounce  
(no big bang singularity)!



classical Friedmann dynamics in GR  
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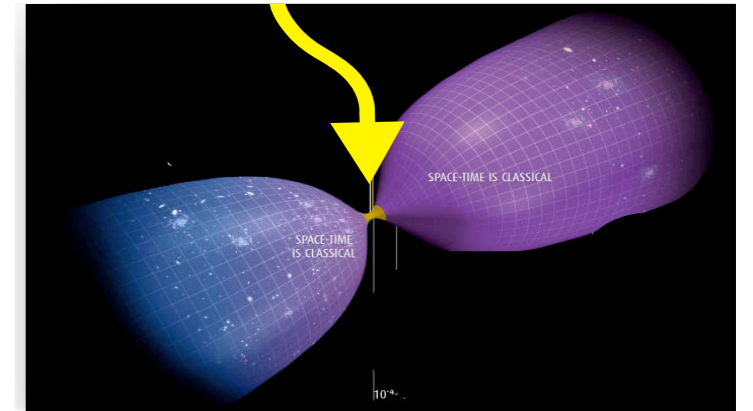
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classical Friedmann dynamics in GR  
(wrt relational clock, with effective  
Newton constant) - flat FRW

- late times: as universe expands, interactions become more relevant, until they drive evolution

—————→ accelerated cosmological expansion

X. Pang, DO, '21

- "phenomenological" approach:

- effective cosmological dynamics

$$w = 3 - \frac{2VV''}{(V')^2}$$

for "emergent matter" component (of QG origin)

order-6 interactions

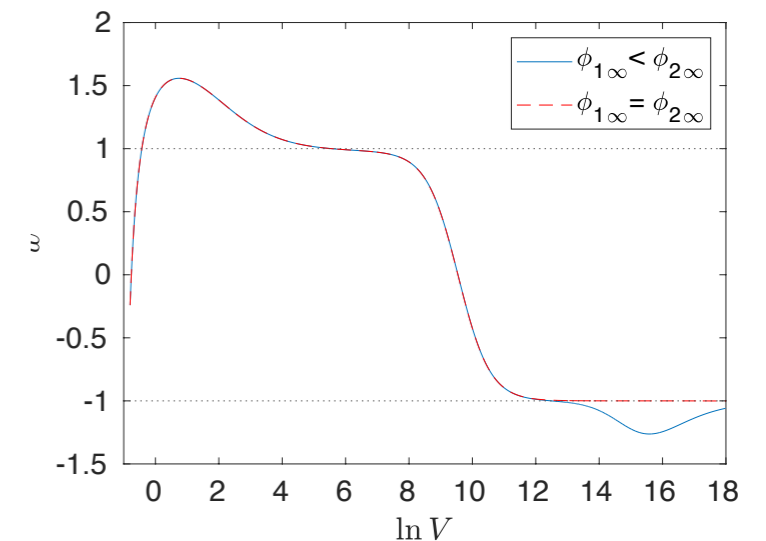
2 modes



effective phantom-like dark energy (of pure QG origin)

+ asymptotic De Sitter universe

X. Pang, DO, '21



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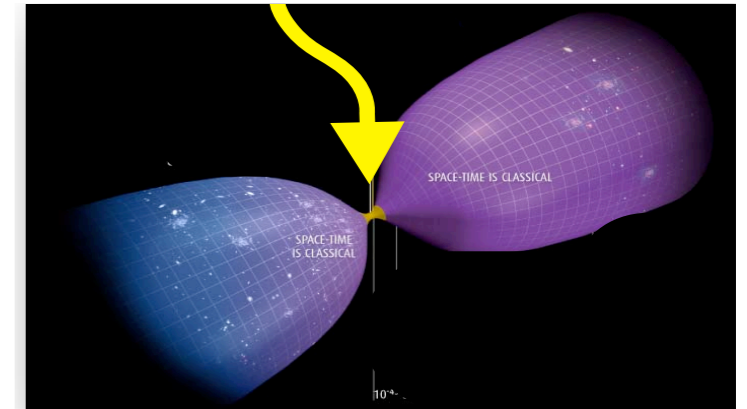
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on parameters of GFT model  
(here written neglecting matter contribution)

$$\left(\frac{V'}{V}\right)^2 = \frac{V''}{V} = 12\pi\tilde{G}$$

classical Friedmann dynamics in GR  
(wrt relational clock, with effective  
Newton constant) - flat FRW

- late times: as universe expands, interactions become more relevant, until they drive evolution

—————> accelerated cosmological expansion

X. Pang, DO, '21

- "phenomenological" approach:

- effective cosmological dynamics

$$w = 3 - \frac{2VV''}{(V')^2}$$

for "emergent matter" component (of QG origin)

order-6 interactions

2 modes



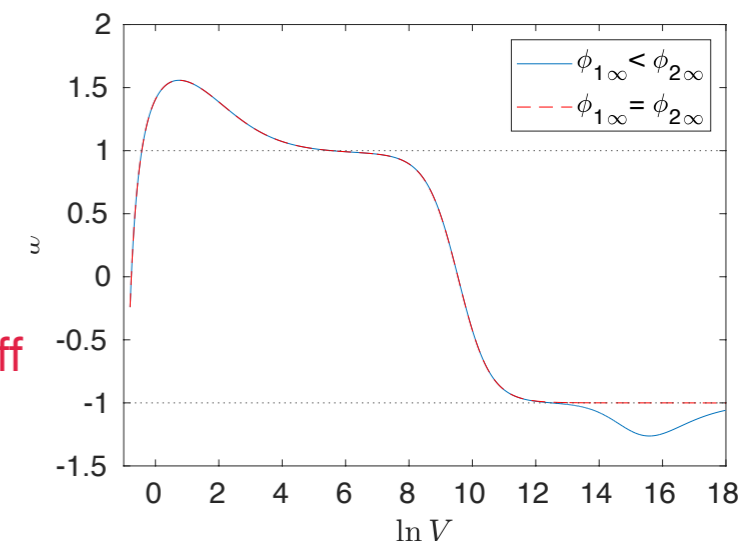
effective phantom-like dark energy (of pure QG origin)

X. Pang, DO, '21

+ asymptotic De Sitter universe

- early-time acceleration (inflation) of pure QG origin possible - but hard to switch off

however, QG affects dynamics of inflaton T. Landstätter, L. Marchetti, DO, to appear



**Foundations of TGFTs (and other  
"non-spatiotemporal QG")**

**and**

**Jaynes' maximal entropy principle**

## Possible foundations of TGFTs (and other "non-spatiotemporal QG")

how can the quantum dynamics be defined, from first principles?

(recall, lacking straightforward classical mechanics foundations as well as canonical quantization justification, due to absence of preferred temporal variable and due to non-local nature)

(also, TGFTs are not the result of quantizing, by any standard technique, classical GR)

- covariant (quantum statistical) path integral

treat TGFTs as statistical (field) systems, defined by a "equilibrium" probability distribution

probability distribution, in turn, defined by standard path integral in terms of "action"



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but how to choose it?                      and what is "equilibrium" in absence of time?

General problem in background independent (classical and) quantum gravity: what is "equilibrium" in absence of (preferred) temporal direction?

C. Rovelli, '12; G. Chirco, T. Josset,  
C. Rovelli, '15; I. Kotecha, '19

one strategy based on Jaynes' entropy maximization

I. Kotecha, DO, '17; G. Chirco,  
I. Kotecha, DO, '18

# TGFT (quantum) statistical mechanics

one strategy for identifying/constructing equilibrium states, applied to TGFT context:

$$\mathcal{H}_F = \mathcal{F}(\mathcal{H}_v) = \bigoplus_{V=0}^{\infty} \text{sym} \left\{ \left( \mathcal{H}_v^{(1)} \otimes \mathcal{H}_v^{(2)} \otimes \dots \otimes \mathcal{H}_v^{(V)} \right) \right\}$$

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M. Montesinos, C. Rovelli, '01; G. Chirco, I. Kotecha, DO, '18

Maximising  $S[\rho] = -\langle \ln \rho \rangle_\rho$  under

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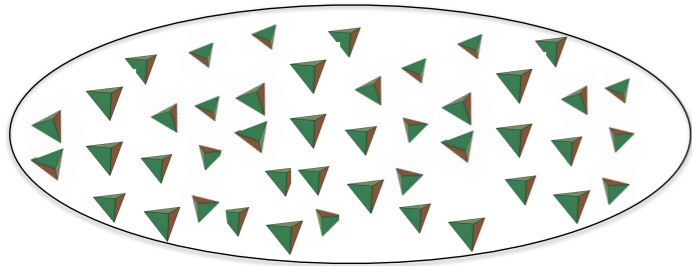
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$$|\psi\rangle = e^{-\frac{\|\psi\|^2}{2}} e^{\int_{SU(2)^4} d\vec{g} \psi(\vec{g}) \hat{\varphi}^\dagger(\vec{g})} |0\rangle, \quad \hat{\varphi}(\vec{g}) |\psi\rangle = \psi(\vec{g}) |\psi\rangle$$

- For dynamical constraint operator  $\hat{C}$ .

$$\begin{aligned} Z &= \text{Tr}_{\mathcal{H}_F} (e^{-\beta \hat{C}}) = \int [D\mu(\psi, \bar{\psi})] \langle \psi | e^{-\beta \hat{C}} | \psi \rangle \\ &= \int [D\mu(\psi, \bar{\psi})] (e^{-\beta \langle \psi | \hat{C} | \psi \rangle} + \langle \psi | : \text{po}(\hat{\varphi}, \hat{\varphi}^\dagger, \beta) : | \psi \rangle) \end{aligned}$$

various choices for C  
(determine TGFT model):  
geometric operators,  
dynamical constraints, ....

- Effective statistical field theory  $Z \approx Z_{\text{eff}} = \int [D\mu(\psi, \bar{\psi})] e^{-C_{\text{eff}}(\psi, \bar{\psi})}$

GFT partition function DO, '13

G. Chirco, I. Kotecha, DO, '18

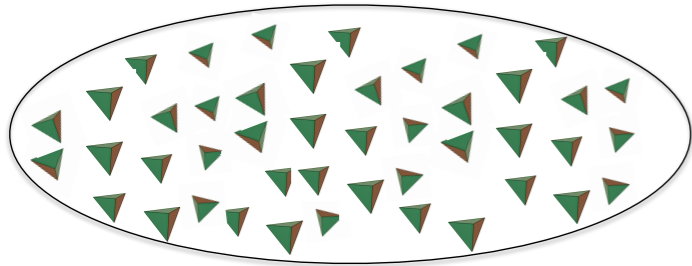
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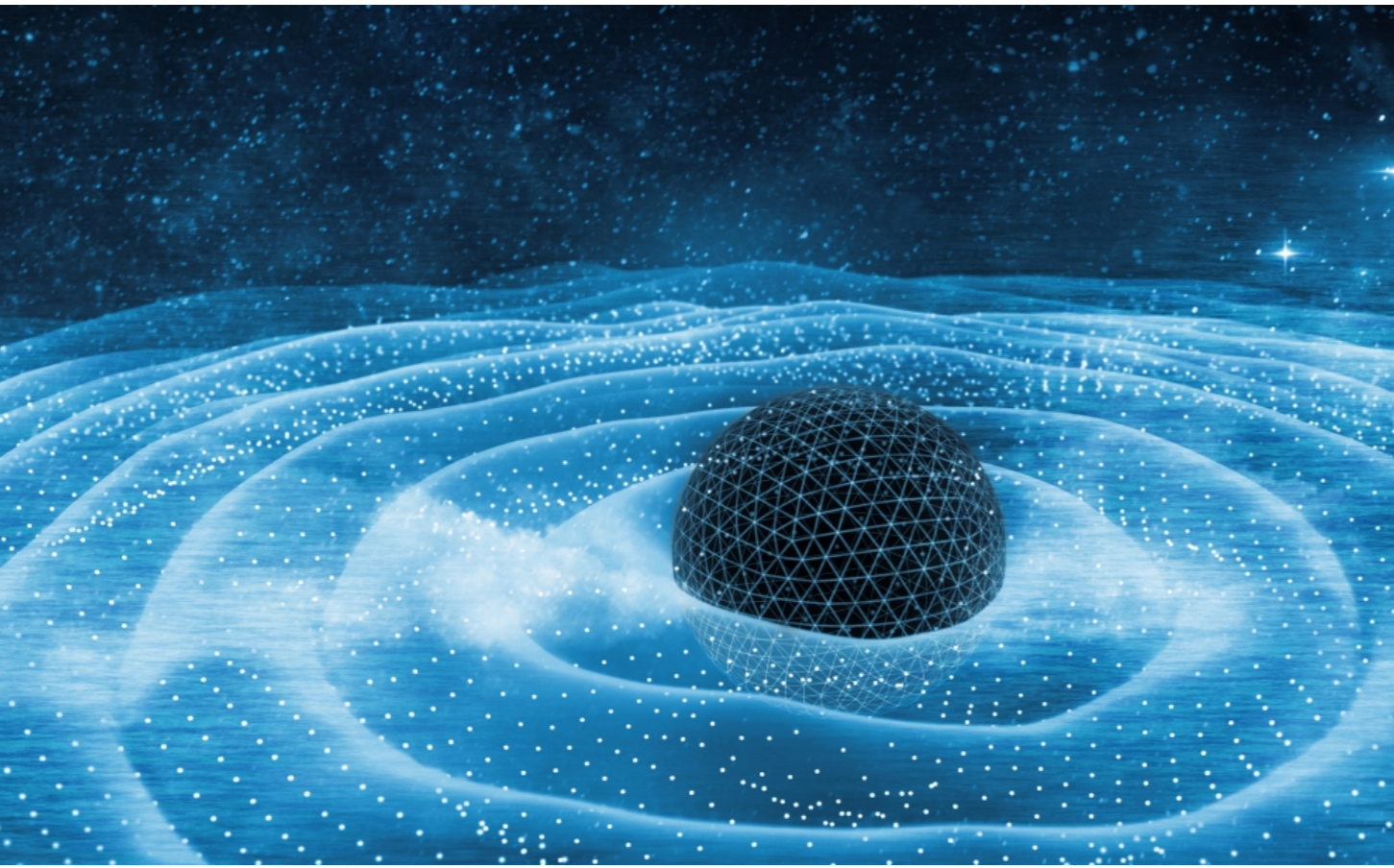
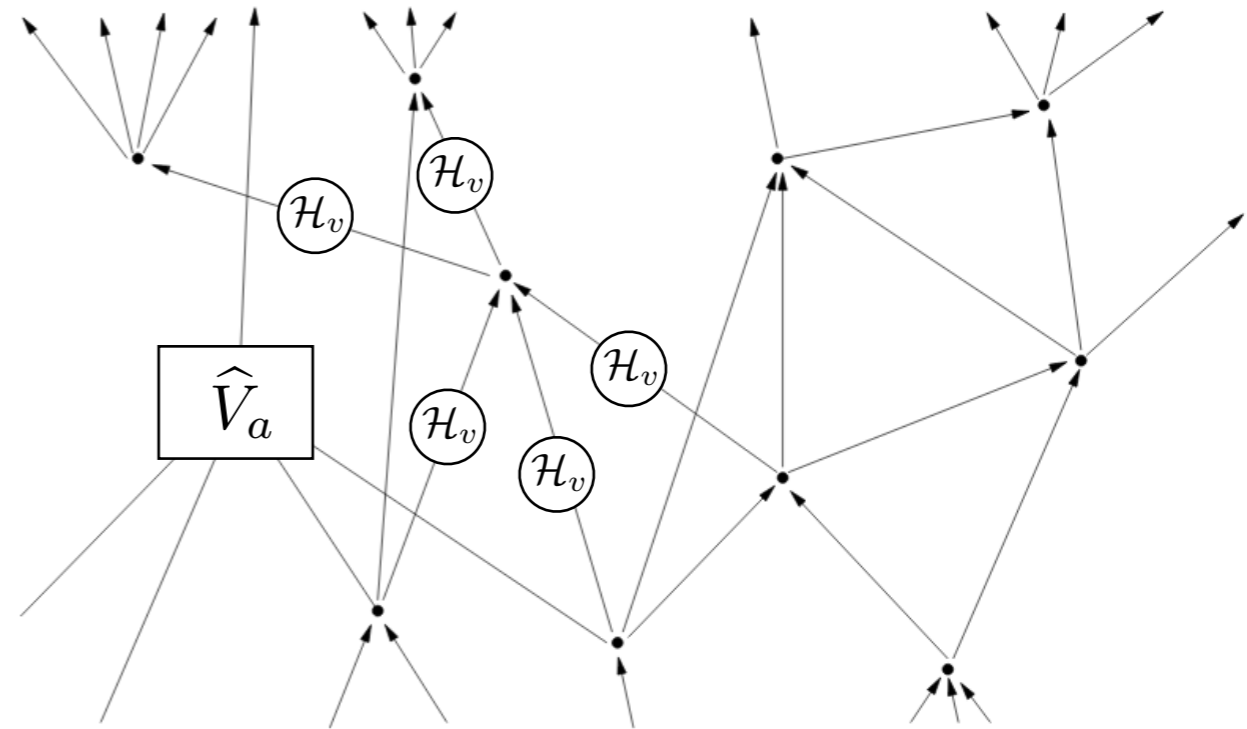
note: Jaynes' principle (and entropy) is epistemic: role of observer/epistemic agent

**remarks:**  
**foundational/philosophical issues in light**  
**of QG**  
**(and the role of agency)**

- (quantum) information and computation
- interpretation of Quantum Mechanics
- laws of nature

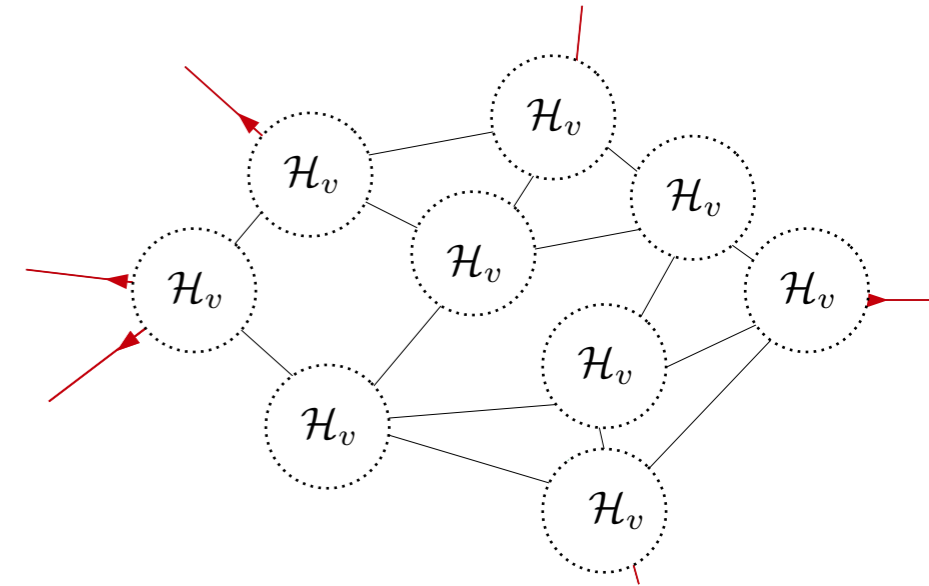


# QG, (quantum) information and computation



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- both semiclassical considerations and QG formalisms suggest that
  - spacetime and gravity as we know them are not fundamental, but emergent, collective notions
  - the universe is a (peculiar, background independent) quantum many-body system of pre-geometric "entities"
- several QG formalisms (eg TGFTs) have combinatorial and algebraic quantum structures as quantum states: quantized simplicial structures & spin networks



- these quantum states can be framed as quantum circuits

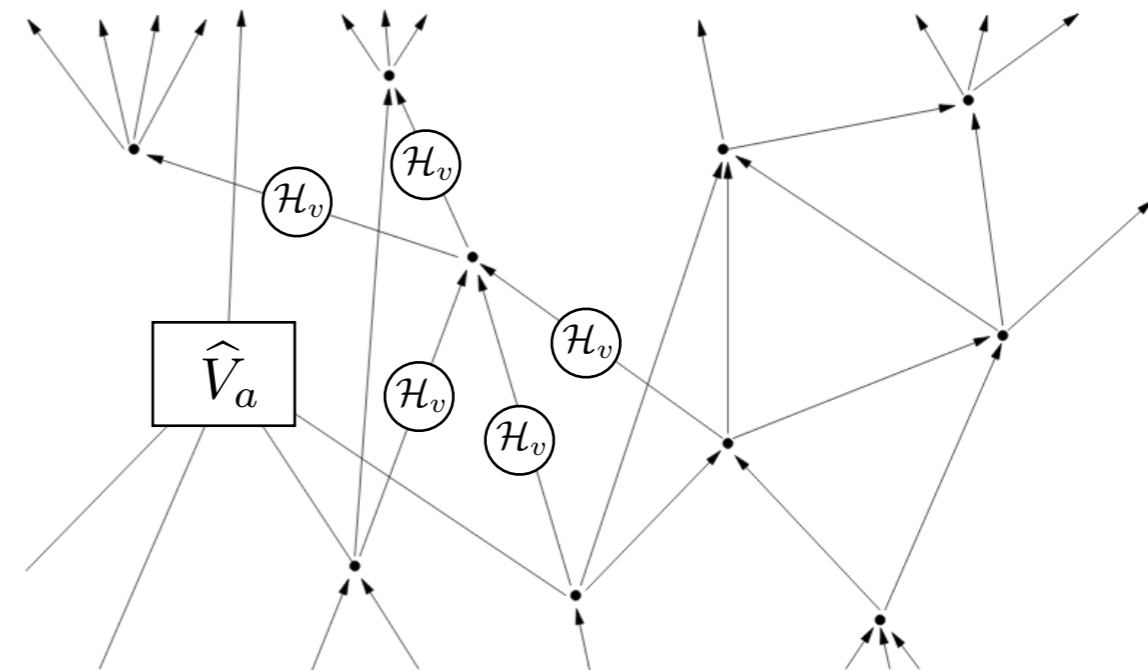
G. Chirco, E. Colafranceschi, DO, '21a,'21b

E. Colafranceschi, S. Langenscheidt, DO, '22 + to appear

Q. Chen, E. Livine, '21; G. Czelusta, J. Mielczarek, '20, '23

- in the same QG formalisms (eg TGFTs), possible dynamical processes take the form of spin foam models (or algebraic versions of lattice gravity path integrals)

- spin foam models can be recast as quantum causal histories
- quantum causal histories can be framed as quantum circuits



F. Markopoulou, '99; E. Livine, DO, '02; E. Hawkins, F. Markopoulou, H. Sahlmann, '03

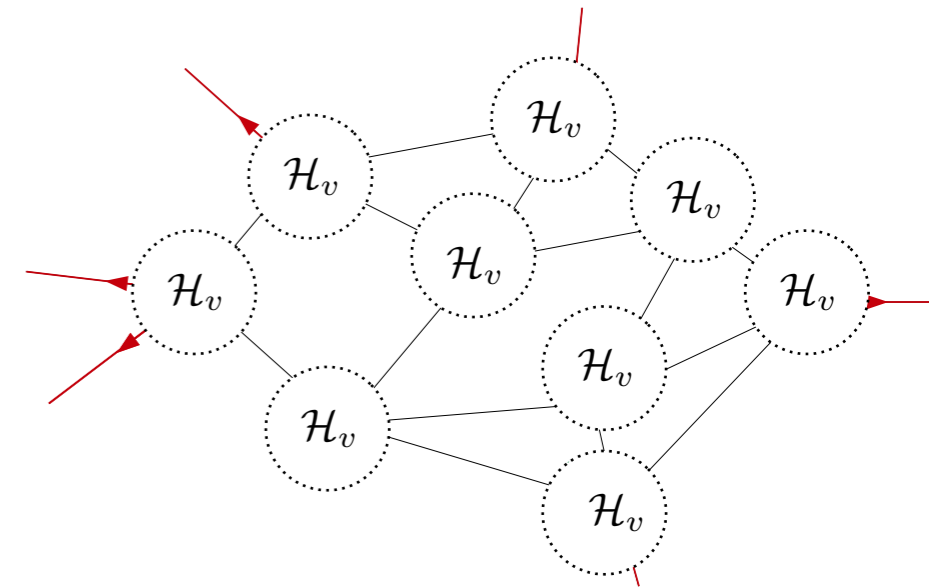
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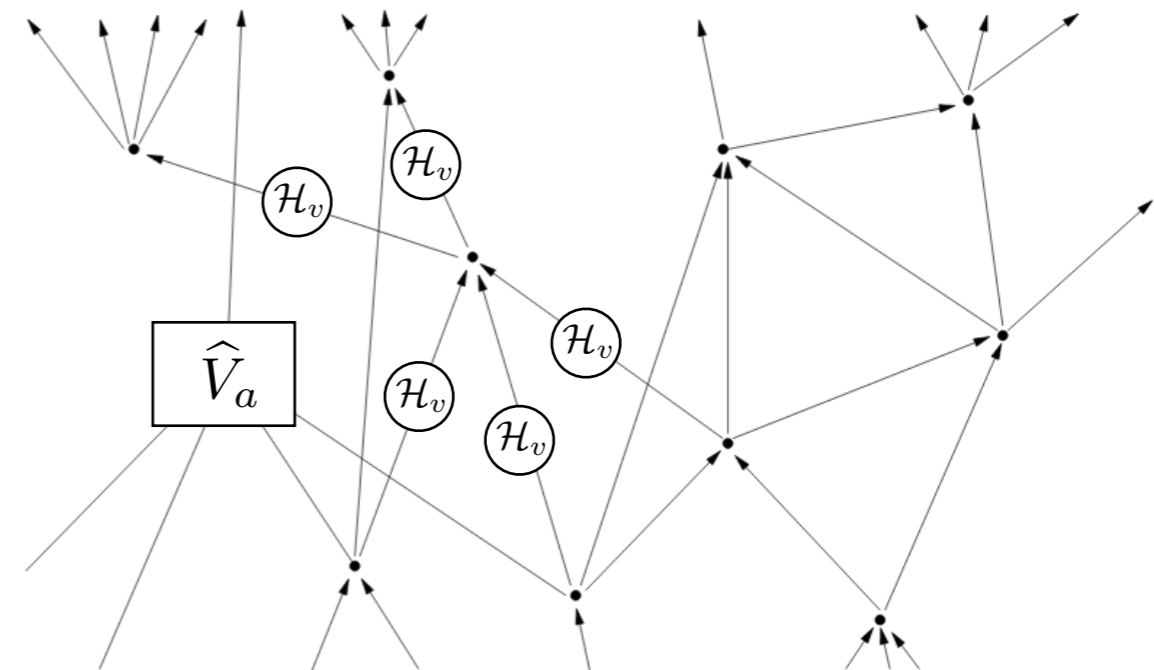
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so: is the universe a quantum computer?

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standpoint and general perspective: an epistemic view on physical laws and the role of agency (see later)

- laws of nature are the product of intelligent agents; their role is irreducible and not negligible (outside idealizations)
- epistemic nature of laws and role of intelligent agents has concrete implications for (our understanding and formulation of) fundamental physics
  - resonances with (and inclinations towards) epistemic perspectives on QM
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in fundamental QG:

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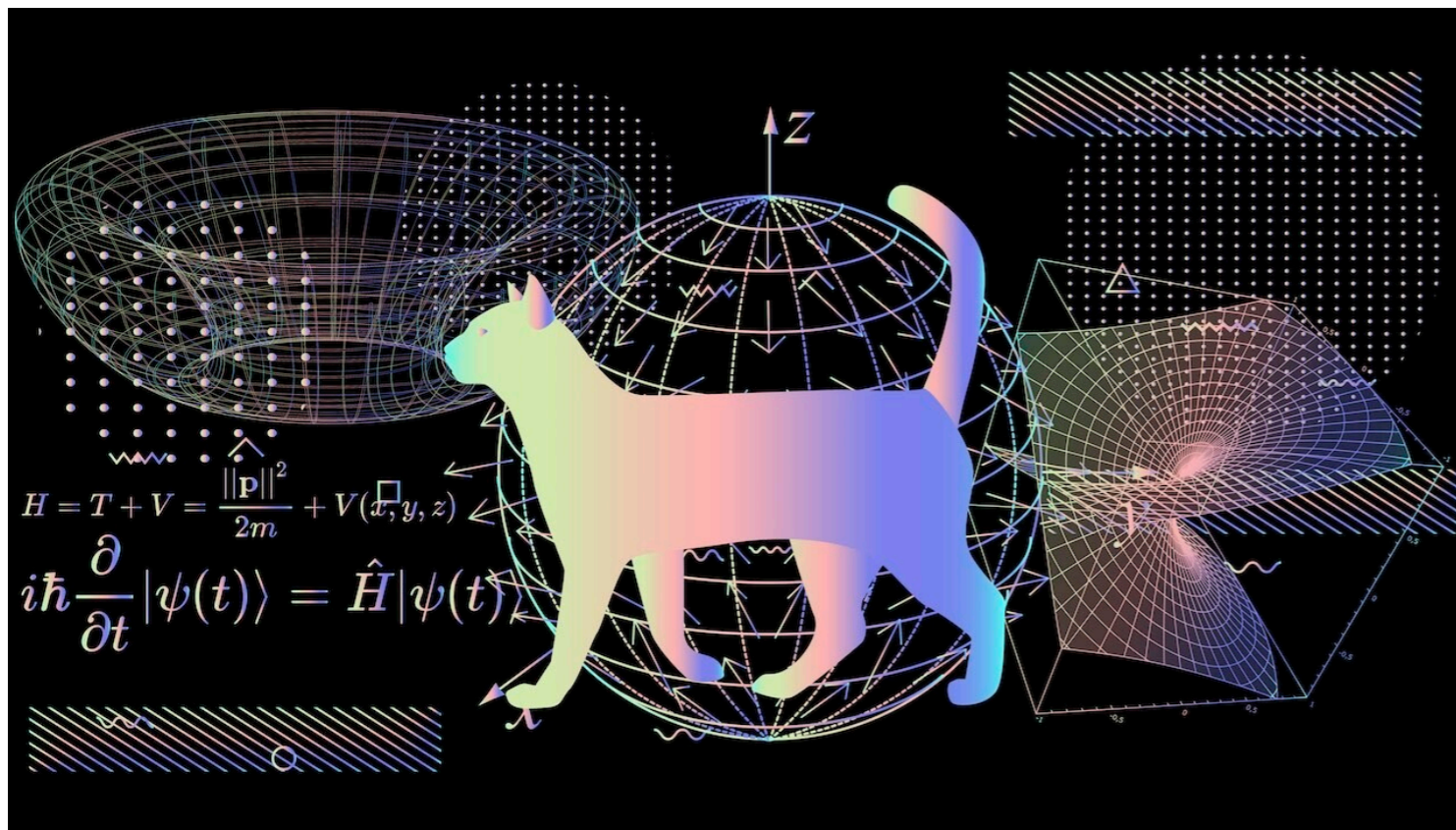
—————→ the universe is (largely) what we think it is, and we think like computers

—————→ the quantum (non-spatiotemporal) universe is naturally modeled as a quantum computer

# Foundations and interpretations of Quantum Mechanics

the foundational issues in Quantum Mechanics

and how Quantum Gravity changes them



QG requires abandoning/generalizing (one or more) basic principles of QM and QFT

locality, unitarity, local Lorentz symmetry?

probably worse in "emergent spacetime" scenarios

but QG generalization  
will necessarily impact  
also QM interpretation!

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topics in quantum foundations of interest for QG

indefinite causality

quantum reference frames

generalised probability theories

beyond unitary quantum evolution

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# Key issues and QM interpretations

1. measurement problem (what is the collapse of wavefunction?)

2. nature of quantum states (are they real?)

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dynamical collapse models

wavefunction collapse is physical process,  
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there are only relative  
(to observer) facts

epistemic-pragmatist perspectives  
(or "neo-Copenhagen")

(Relational QM, Müller's interpretation, Healey's  
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J. Pienaar, '21; A. Barzegar, DO, '22

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underlying ontology is not spatiotemporal

observer very far from object (no direct experience)  
epistemic aspects (not operational) are even more dominant

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# Laws of nature

## what are they?

very long-standing issue in philosophy (phil. science, epistemology, metaphysics, ...)

(Armstrong, Ayer, Callender, Cartwright, Cohen, Dretske, Giere, Hüttermann, Lewis, Maudlin, Mill, Psillos, Ramsey, Skyrms, Van Fraassen, .....)

vague notion: "general relations among properties of physical systems"

are they objective and intrinsic to the world or epistemic in nature?

and how does QG change the story?

S. Hartmann, DO, in prog

V. Lam, DO, in prog.



# Humeanism

"laws as patterns of facts in the world (and in spacetime)"

- **Ontological picture**: fundamental basis of non-modal facts, on which laws (and everything else...) supervene
- **Humean basis** (D. Lewis): distribution of fundamental intrinsic properties over spacetime.
  - Spacetime relations as 'world-making' (or 'gluing') relations.



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but which regularities are laws and why?

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## Best systems

laws as propositions in the best systematizations of regularities

D. K. Lewis: *Counterfactuals*

"[...] a contingent generalization is a law of nature if and only if it appears as a theorem (or axiom) in each of the true deductive systems that achieves a **best combination** of **simplicity** and **strength**."

but how to "measure" simplicity and strength? are they subjective?

**essential epistemic elements**



## Humeanism

"laws as patterns of facts in the world (and in spacetime)"

- **Ontological picture**: fundamental basis of non-modal facts, on which laws (and everything else...) supervene
- **Humean basis** (D. Lewis): distribution of fundamental intrinsic properties over spacetime.
  - Spacetime relations as 'world-making' (or 'gluing') relations.



## Regularity theory

"... all that is required for there to be a law in nature is the existence of *de facto* regularities. In the most straightforward case, the constancy consists in the fact that events, or properties, or processes of different types are invariably conjoined with one another." (A.J. Ayer: 'What is a Law of Nature?')

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## Primitivism/dispositionalism

- **Ontological picture**: some (irreducible) primitive modality gives rise to ("produces") the spatiotemporal distribution of particular facts.
  - Primitivism (Maudlin): fundamental physical laws are ontological primitives
  - Dispositionalism: laws are grounded in the fundamentally dispositional or causal nature of properties.



## an strong epistemic view on laws is close to law antirealism:

S. Hartmann, DO, in prog

### Why laws are not real

Van Frassen, Cartwright,  
Giere, ....

because that's the simplest solution of the conceptual problems raised by assuming they exist

because they are simply not factual (they do not even represent observed facts)

(scientific theories are collections of models, all "laws" actually used by scientists are approximate and ad hoc rules tailored to specific, limited situations, with no real claim of generality or fundamentality)

### can we be content with this blunt anti-realist view on laws?

- only provided one can account for the many functions laws fulfil in science, without assuming their existence "out there"
- this may require a different understanding of scientific explanations of natural phenomena, not metaphysically loaded, possibly more limited (empirical adequacy); scientific theories understood as "guiding clues" for belief about the world

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QG poses several new challenges to existing accounts of laws

# Humeanism

- **Humean basis:** if spacetime is not fundamental, the traditional Humean "basis" is not fundamental
- More general challenge from quantum theory (Maudlin): quantum entanglement relations do not supervene on the Humean basis (are "non-local")
- **Crucial difficulty:** what provides and how to characterise the Humean basis in a context without spacetime?
  - The non-spatio-temporal characterisation of the Humean basis will be based on fundamental QG entities
  - Some QG approaches suggest **quantum entanglement relations (between QG entities)** as 'gluing relations'

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## Primitivism/ dispositionalism

- These non-Humean conceptions operate against a primitive temporal and causal background
  - Maudlin (2007, 182): "the total state of the universe is, in a certain sense, derivative: it is the product of the operation of the laws on the initial state"
- **difficulty: how to articulate a non-temporal nomic production without spacetime?**
  - QG 'processes' that could instantiate it: spin foam / GFT transition amplitudes; primitive combinatorial/algebraic structures endowed with fundamental dispositions.
- difficulty: notion of 'production' seems to involve some ("causal") asymmetry
  - some form of 'ordering' in QG amplitudes should be present ("proto-causality")

## QG challenges to agent-first (epistemic) accounts (or law anti-realism)

- an epistemic view on laws could be more flexible to adapt to the absence of spacetime at the fundamental level
- **key challenge:** build a QG theory with strong explanatory power, despite being remote from experience (thus also far from operationalism) and underdetermined by observations
- its laws will be grounded in its epistemic virtues, and so will be its suggested ontology
- strongly relying (concerning non-directly observable QG entities) on epistemic tools of abstraction, imagination, counterfactuals, hypothetical reasoning, analogies





**Thank you for your attention!**