#### Supernova Theory

(and Cosmological Distances)

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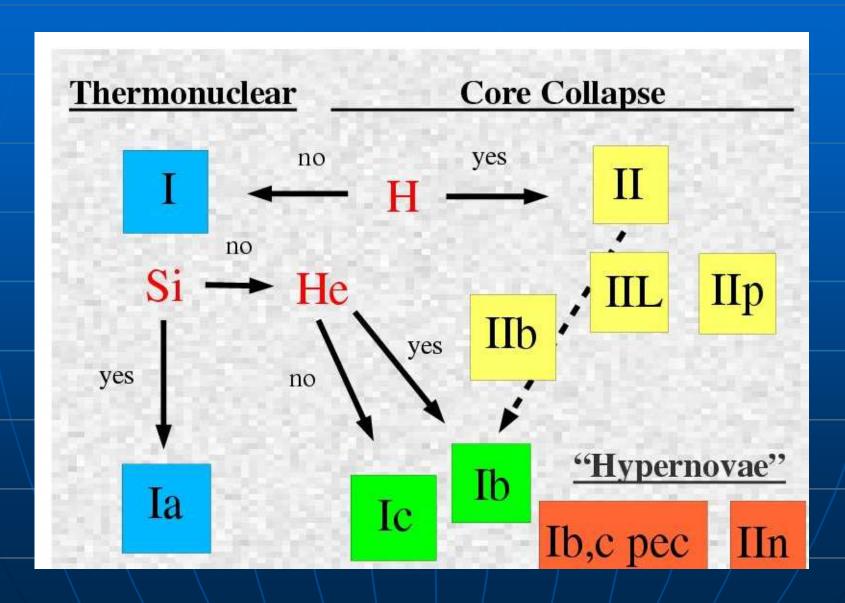
Scuola Nazionale di Astrofisica, Sant'Agata, May 8 – 13, 2005

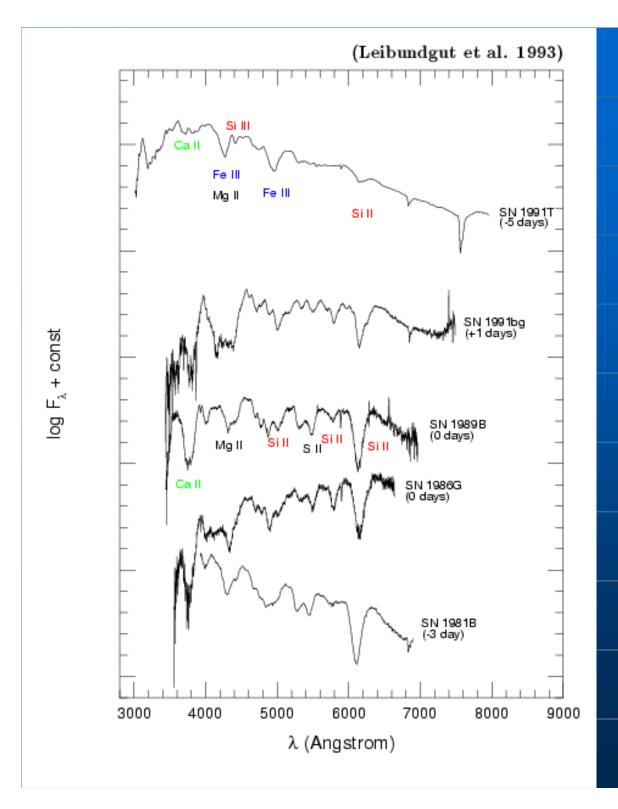


#### Outline of the lectures

- n Supernova types and phenomenology
- Models of core-collapse supernovae (Type II; Type Ib,c; GRB's)
- Models of thermonuclear supernovae(Type Ia)
- Luminosity distances and supernovacosmology

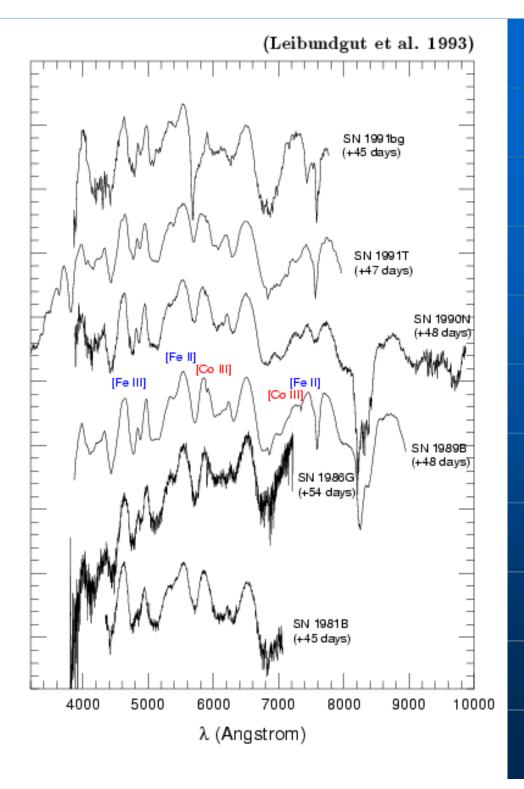
#### Supernova classification



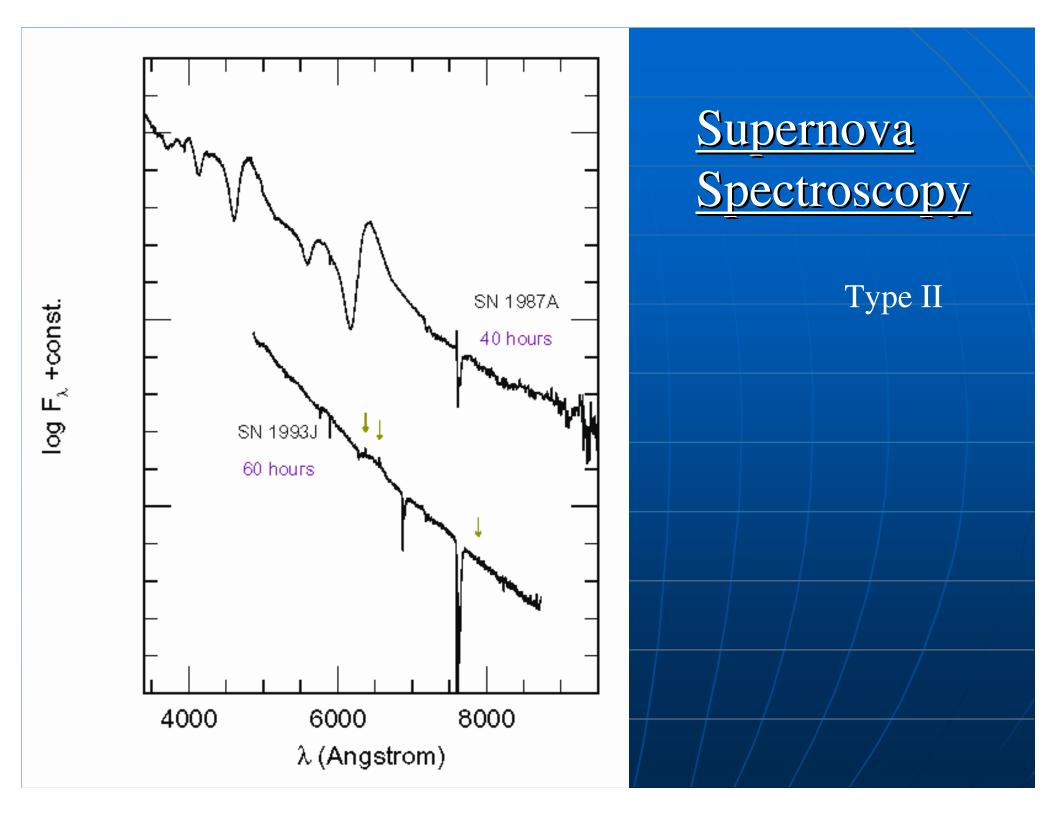


Type Ia

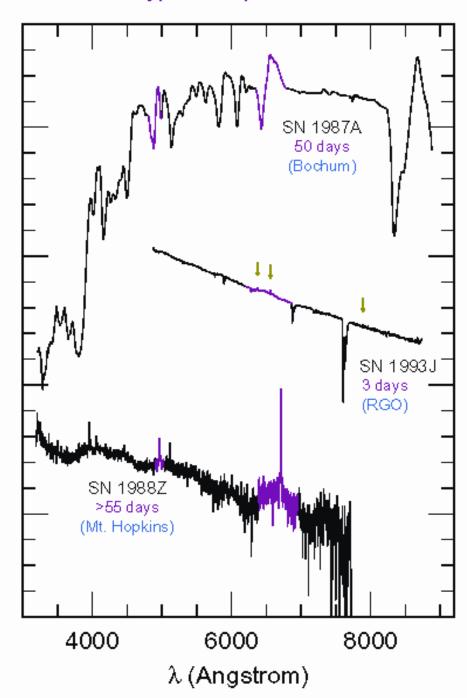
Type Ia



Type Ia

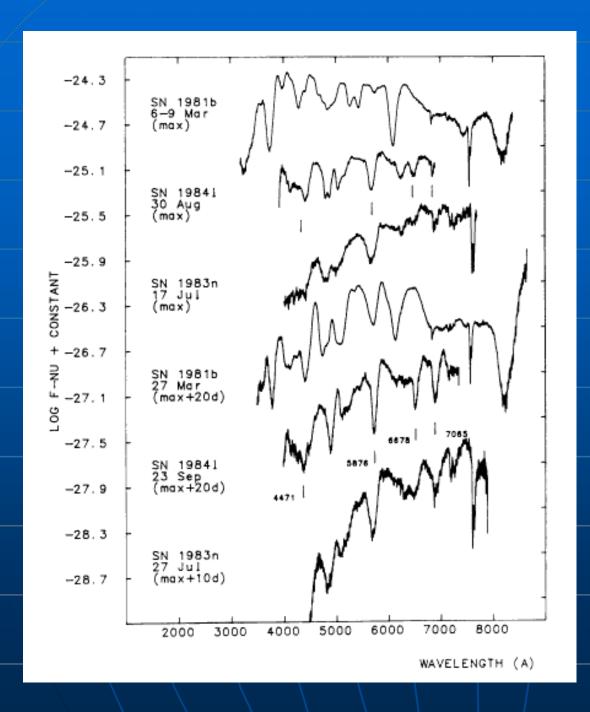


#### Type II Supernovae

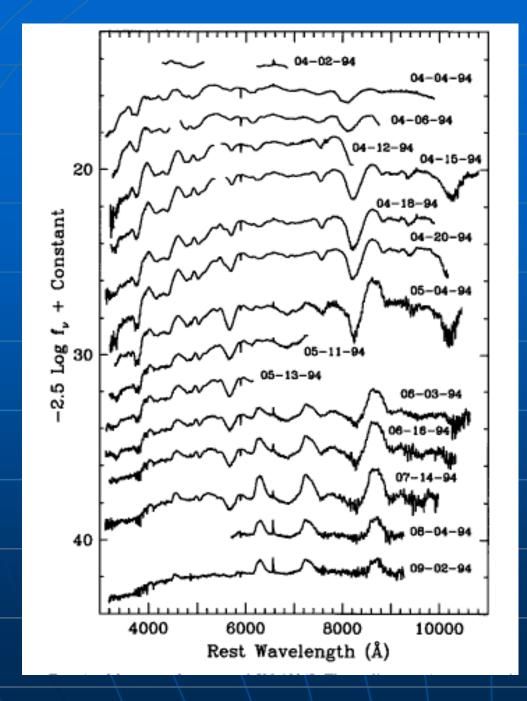


#### Supernova Spectroscopy

Type II

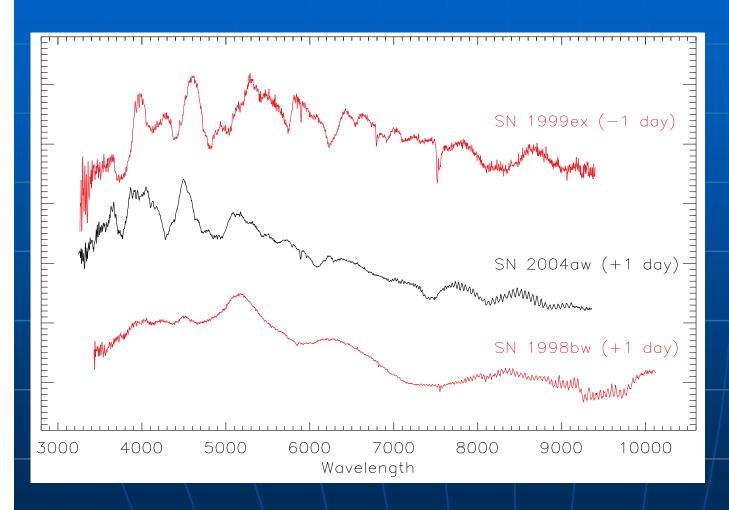


Type Ib



Type Ic

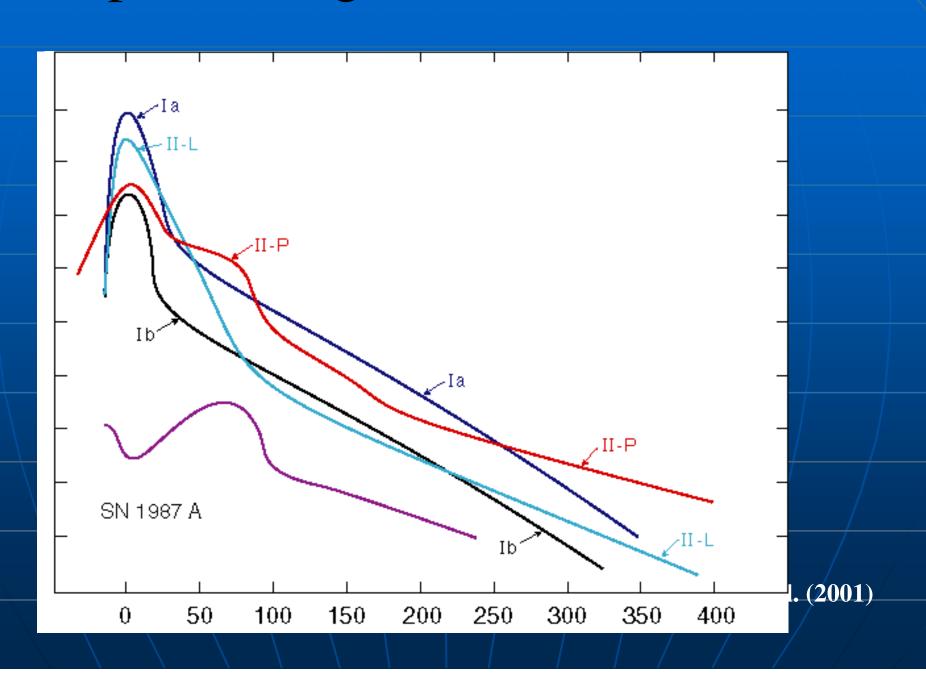
SN 1994I (Filippenko et al.)

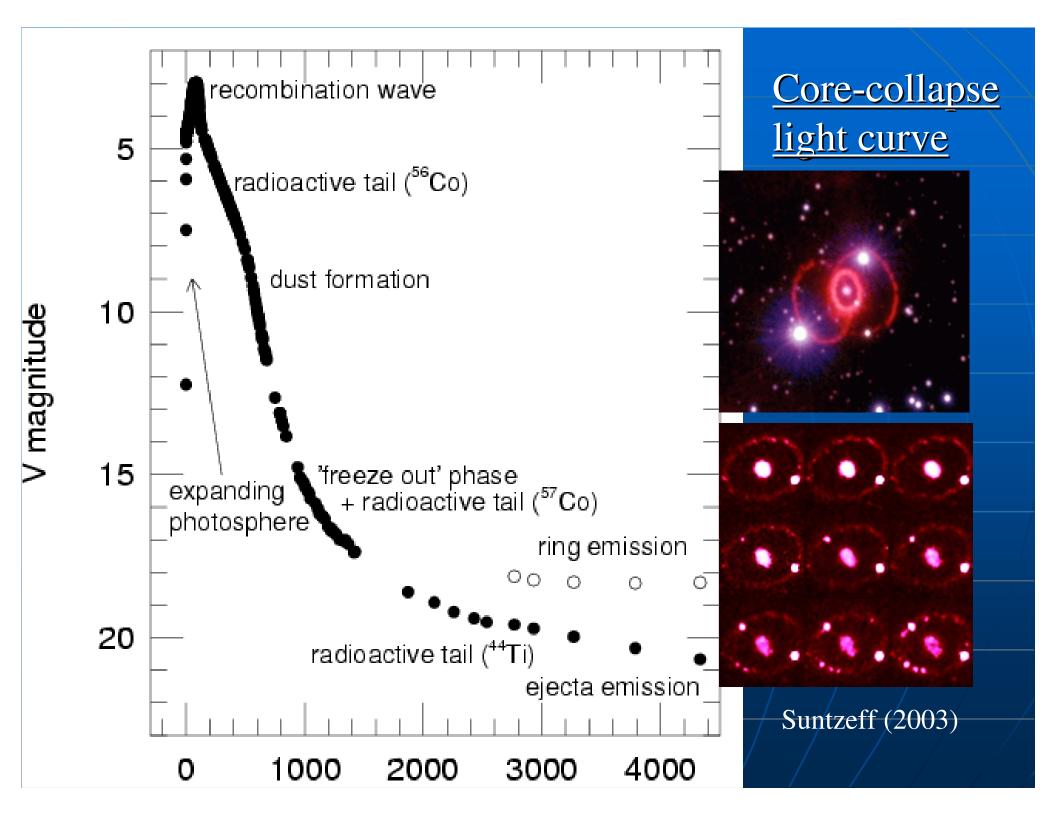


Type Ic pec

(Taubenberger et al. 2005)

#### Supernova light curves





#### Energy sources for light curves

#### n shock

- breakout
- kinetic energy

#### n cooling

• due to expansion of the ejecta

#### n radioactivity

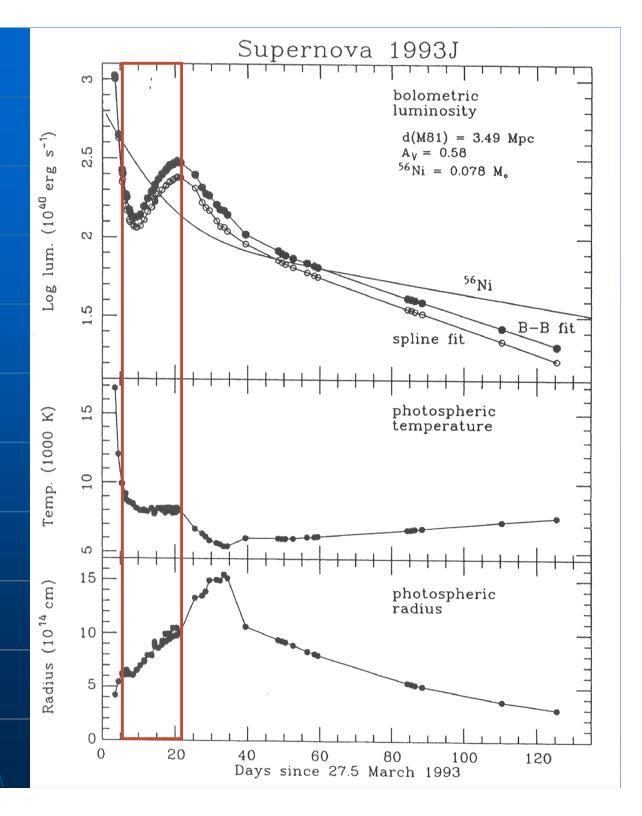
nucleosynthesis

#### n recombination

• of the shock-ionised material

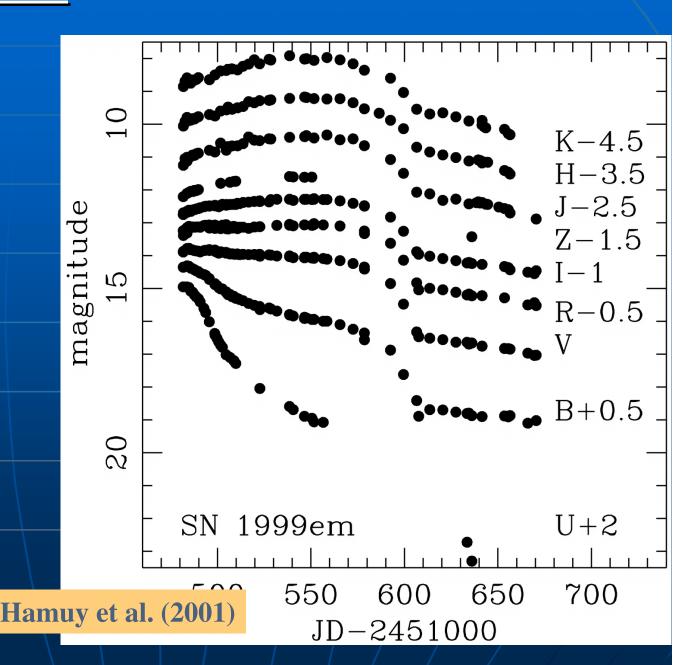
#### Expansion

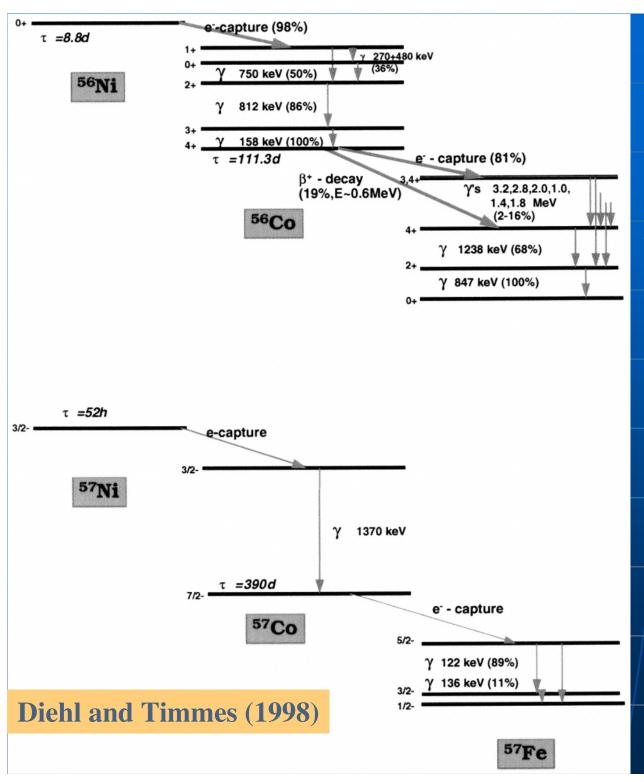
- n Brightness increase
  - increased surface area
  - slowtemperaturedecrease



#### Recombination

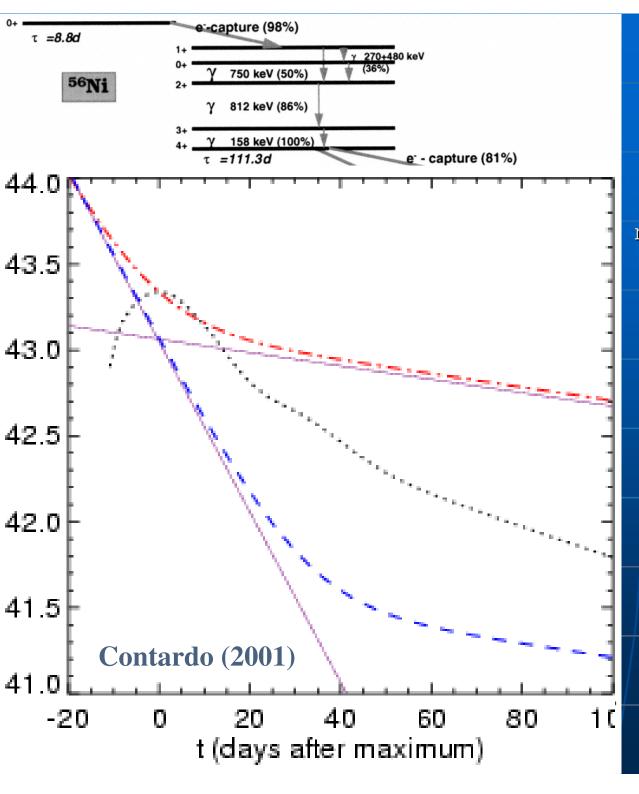
- n Balance of the recombination wave and the expansion of the ejecta
  - leads to an extended plateau phase





#### Radioactivity

- Isotopes of Ni and other elements
  - conversion of γrays and
    positrons into
    heat and optical
    photons



#### Radioactivity

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    positrons into
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#### Supernova types: Summary

#### n Thermonuclear SNe

- from low-mass stars (<8M)
- highly evolved stars (white dwarfs)
- explosive C and O burning
- binary systems required
- complete disruption

#### n Core-collapse SNe

- high mass stars (>8M)
- large envelopes (still burning)
- burning due to compression
- single stars (binaries for SNe Ib/c?)
- neutron star

#### Core-collapse Supernovae

(In part "borrowed" from a lecture by Ewald Müller)



#### Prototype:

Crab nebula with pulsar (constellation Orion)

Remnant of a supernova observed in 1054

#### 30 Doradus region in the Large Magellanic Cloud

(d ~ 160 000 light years)



Supernova 1987A 7:35 UT 23.2.1987 Blue Supergiant Sanduleak 69.202

#### A few observational facts

(core collapse supernovae, i.e. SNe II, Ib, Ic)

very bright events:

 $L \sim 10^{10} L_{sun}$ 

fast expanding ejecta:

 $v \sim 10^4 \, \text{km/s}$ 

energies: electromagnetic:

 $\sim 10^{49} \text{ erg}$ 

kinetic:

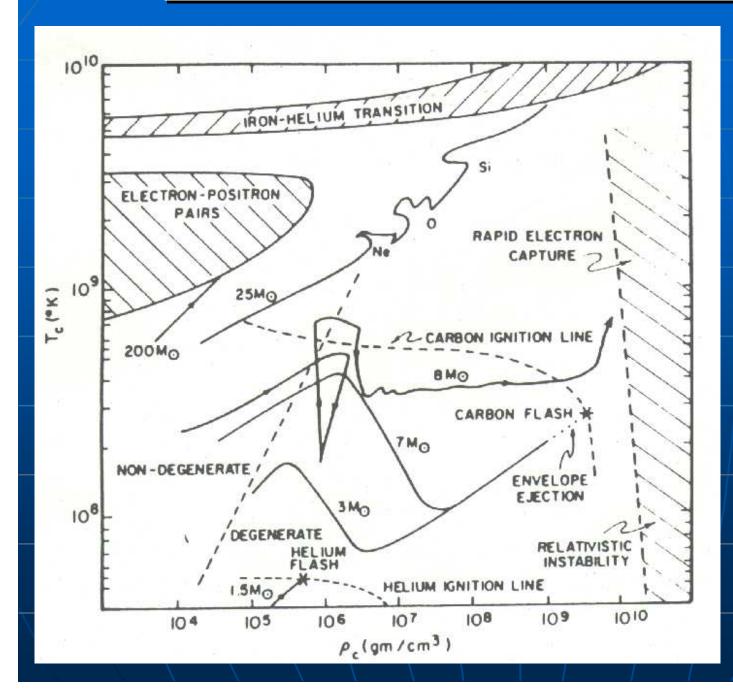
 $\sim 10^{51} \text{ erg}$ 

neutrinos (SN1987A):  $\sim 3.10^{53}$  erg

progenitor star distroyed (SN 1987A, SN 1993J)

compact remnant (as far as we know)

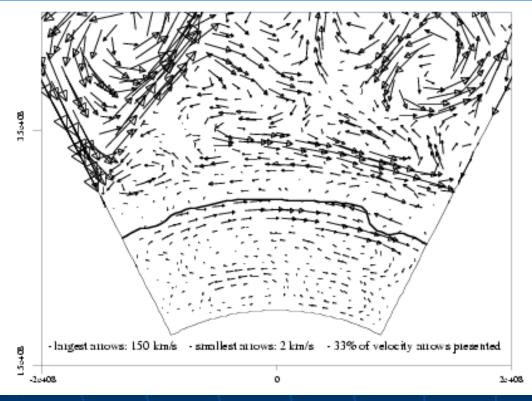
#### Evolution of massive stars

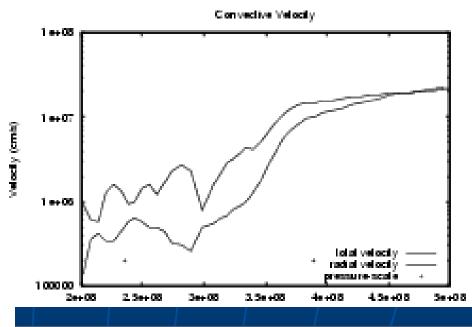


For models of the explosions:

"Fe"-core
masses and their
entropies
have to be
known!

#### Problems with massive star evolution: Non-local time-dependent convection!

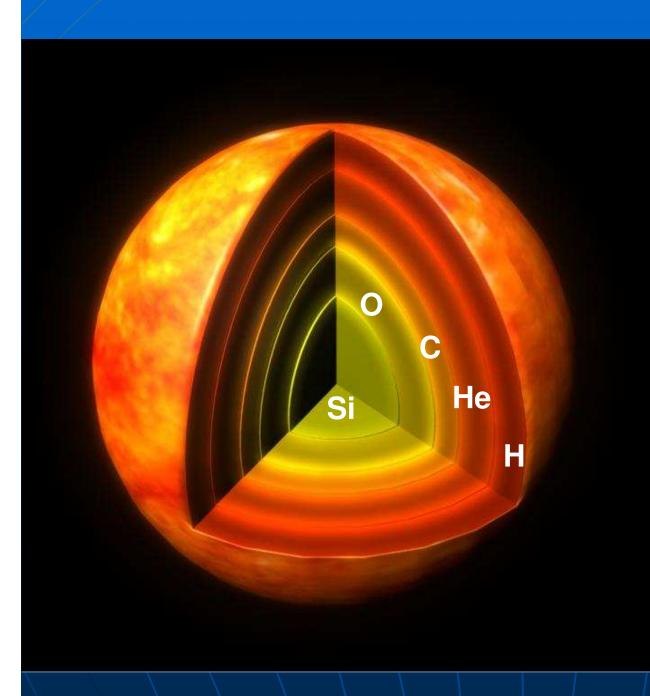




Radius (cm)

Hydrostatic O-burning (Asida & Arnett, 2000)

(See also Brummell et al. 2002, Rogers et al. 2003, ....)



Onion-like structure of a <u>presupernova</u> star several million years after its birth:

mass: 10 ... 10<sup>2</sup> M<sub>sun</sub> radius: 50 ... 10<sup>3</sup> R<sub>sun</sub>

- shells of different composition are separated by active thermonuclear burning shells
- core Si-burningleads to formationof central iron core

Note: figure not drawn to scale!

#### Energy sources for a core collapse supernova

#### Gravitational binding energy

Formation of a compact object of ~1 solar mass with a radius ~10km

$$\rightarrow$$
 E<sub>b</sub> ~ 3 x 10<sup>53</sup> (M/M<sub>sun</sub>)<sup>2</sup> (R/10km) <sup>1</sup>erg

Fe-Ni core:  $\rho \sim 10^{10} \text{ g/cm}^3$ ,  $T \sim 10^{10} \text{ K}$ 

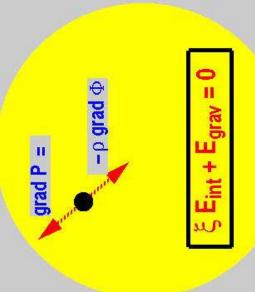
- → P<sub>e</sub> (relativistic degenerate Fermi gas)
- → maximum mass (Chandrasekhar)

#### Core becomes unstable due to:

- a) electron captures
- b) photo-disintegrations

## Evolution towards gravitational collapse

stellar evolution mostly hydrostatic, i.e pressure and gravitational forces are in equilibrium



virial theorem

 $\xi := 3P/\rho u$  ideal gas:  $P = (\gamma-1) \rho u \longrightarrow \xi = 3(\gamma-1)$  relativ. Fermi gas:  $P = 1/3 \rho u \longrightarrow \xi = 1$ 

total energy:

W := Eint + Egrav = 
$$(1-\xi)$$
 Eint =  $(\xi-1)/\xi$  Egrav

# Evolution towards gravitational collapse

gas: finite temperature --> star radiates

### energy conservation:

$$\frac{dW}{dt} + L = 0$$

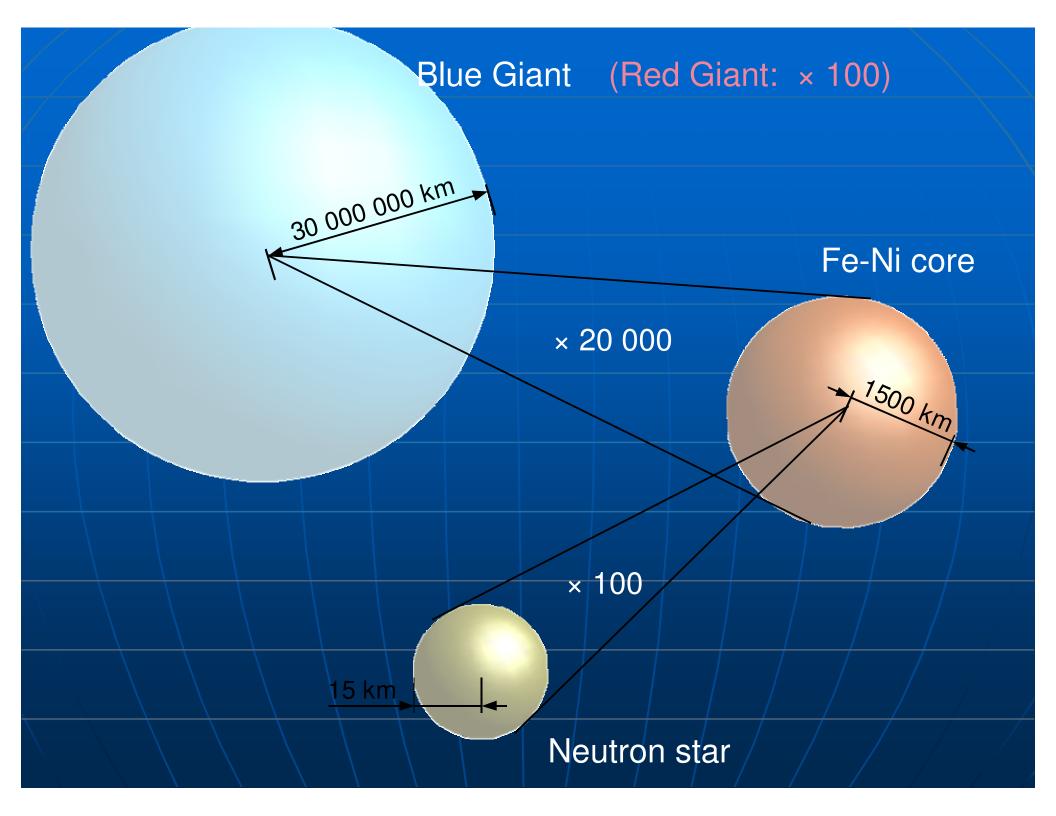
$$L = (\xi-1) dE_{int} / dt = -(\xi-1)/\xi dE_{grav} / dt$$

contraction with  $\gamma = 5/3$  ( $\xi = 2$ ):

50% of liberated energy are radiated away

50% of liberated energy heat the star

## --- star has negative specific heat!



#### Simulations of core collapse supernovae challenging, because of:

a) neutrino transport (fermions, multi-flavor)

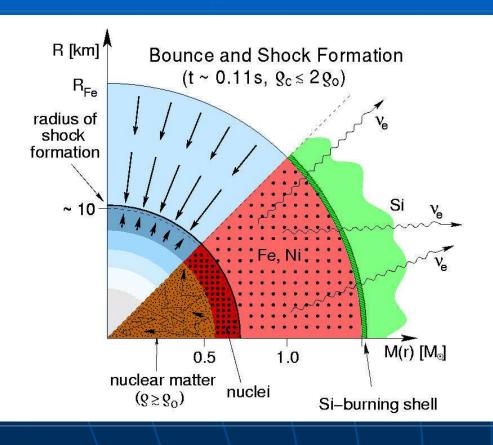
(semi-transparent region: Boltzmann solver)

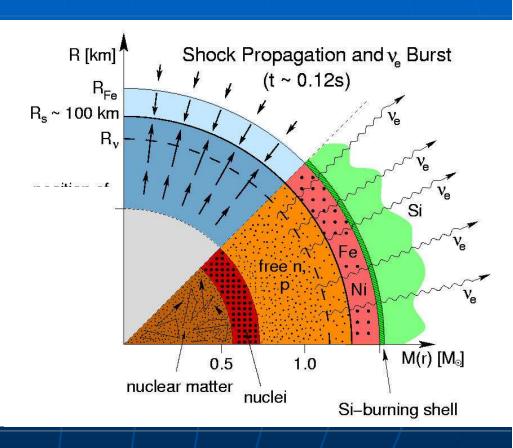
- b) very different time and length scales
  - → adaptive mesh refinement (AMR)

c) multi-dimensional flow problem

#### Core collapse supernovae:

- Prompt explosion mechanism does <u>not</u> work (explored during the 1970's and 1980's; commonly accepted early 1990's)

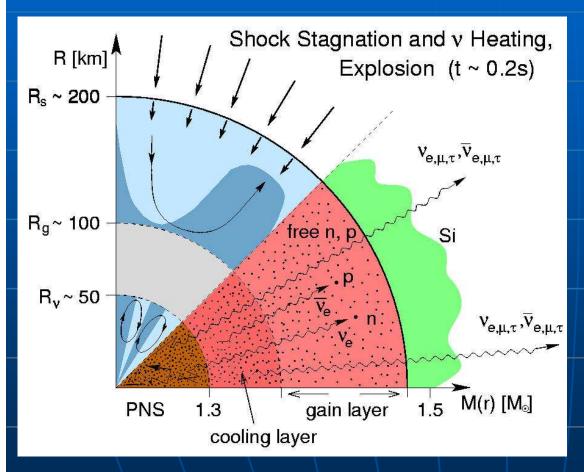


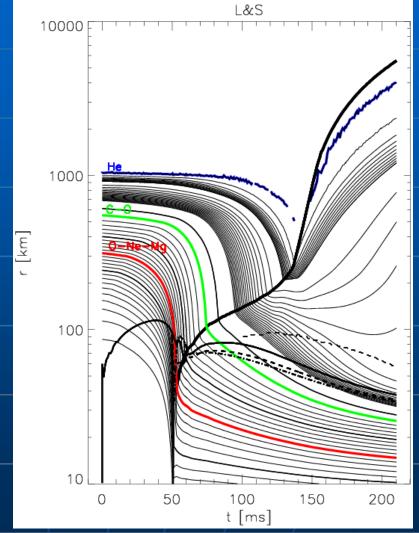


- Shock wave forms close to sonic point ( M ~ 0.7 M<sub>sun</sub> ) initial energy: (5 ... 8) x 10<sup>51</sup> erg
- Severe energy losses during shock propagation (8 MeV/nucleon or 1.6 x 10<sup>51</sup> erg/0.1M<sub>sun</sub>)

- <u>Current paradigm</u>: neutrino-driven delayed explosions (discovered through computer simulations by Wilson '82, and

first analyzed by Wilson & Bethe '95)

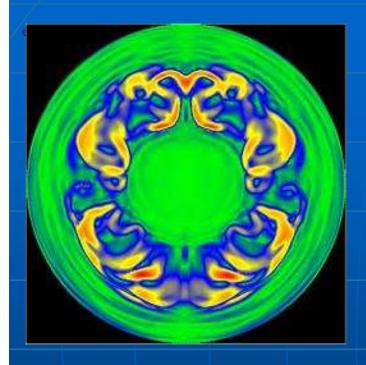




In its simplest form: Seems to work for low-mass core only! (Kitaura et al '05)

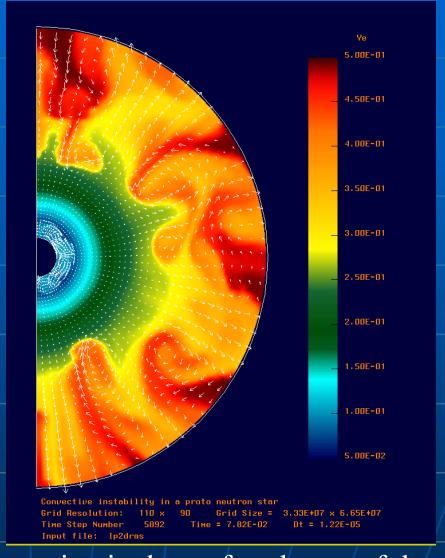
- Ø Observations imply: non-radial flow and mixing are common in core collapse supernovae
- Ø Theoretical models based on delayed explosion mechanism predict non-radial flow and mixing due to:
- Ledoux convection inside the proto-neutron star (deleptonization and neutrino diffusion)
- convection inside neutrino heated hot bubble
   (neutrino energy deposition behind the shock)
- Rayleigh-Taylor instabilities in stellar envelope (non-steady shock propagation; hot bubble)

#### Core collapse supernovae need multidimensional modeling!



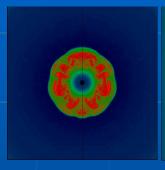
Ledoux convection inside proto-neutron star due to negative lepton and entropy gradients (Keil, Janka & Müller '96)

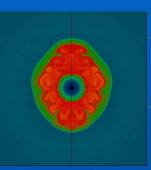
- asymmetric v-emission (few sec) and flow (~100 s?)

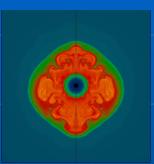


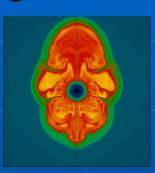
Convection in the surface layers of the proto-neutron star and in the hot bubble 78 ms after core bounce (Janka & Müller '96)

### State-of-the-art hydrodynamic simulations with Boltzmann v-transport, realistic EOS, relativistic gravity, and realistic progenitors

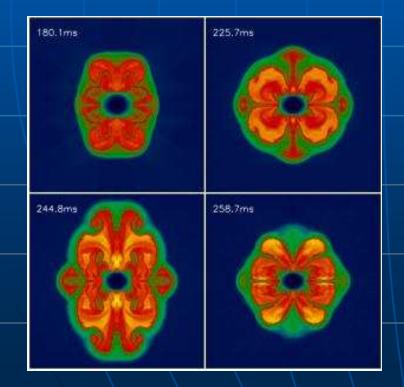








Snapshots, 2D run of a <u>non</u> rotating axisymmetric 11.2 M<sub>sun</sub> progenitor (Buras, Rampp & Janka 2003)

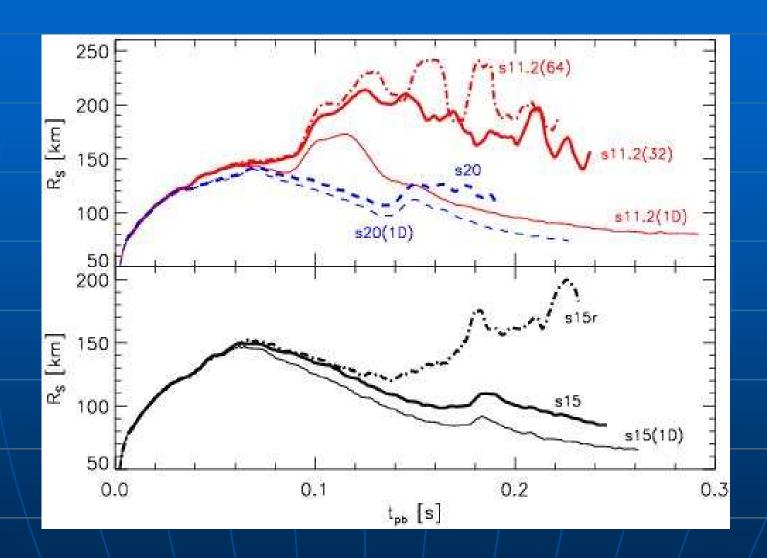


Snapshots, 2D run of a <u>rotating</u> ( $b_{initial} = 0.05\%$ ,  $\omega_{i,c} = 0.5s^{-1}$ ; Heger etal 2003)

axisymmetric 15  $M_{sun}$  progenitor

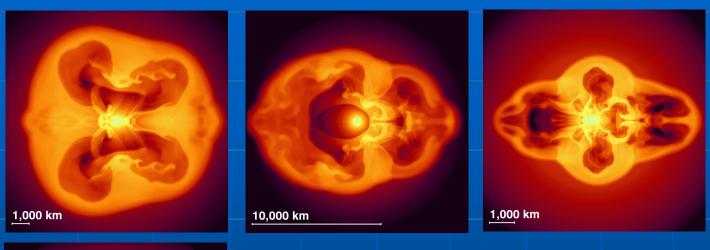
(Buras, Rampp, Janka & Kifonidis 2003)

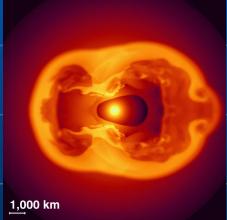
#### Core Collapse Supernovae: But ...



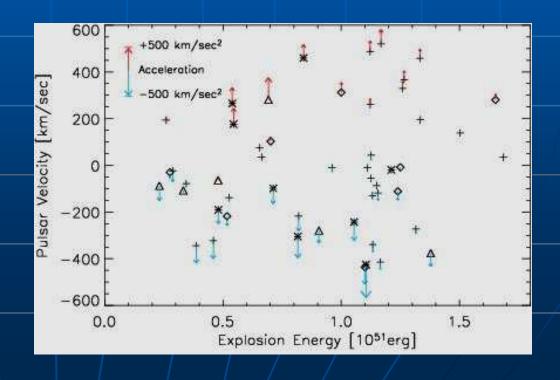
Buras et al. (2003): No explosions (in 2-D)!!!

#### 2D simulations show growth of dipolar (l=1) mode in post-shock layer → neutron star kicks (Scheck et al. 2003)

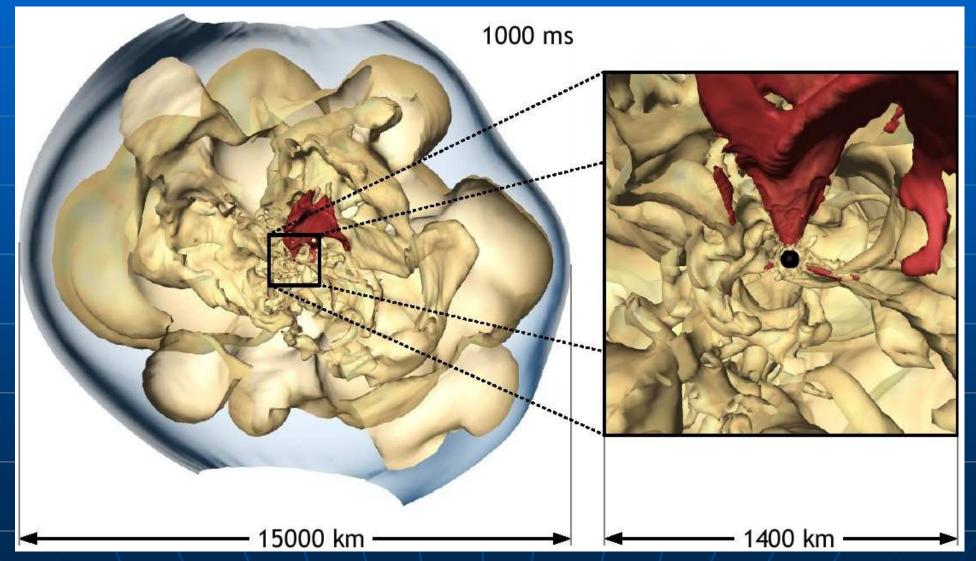




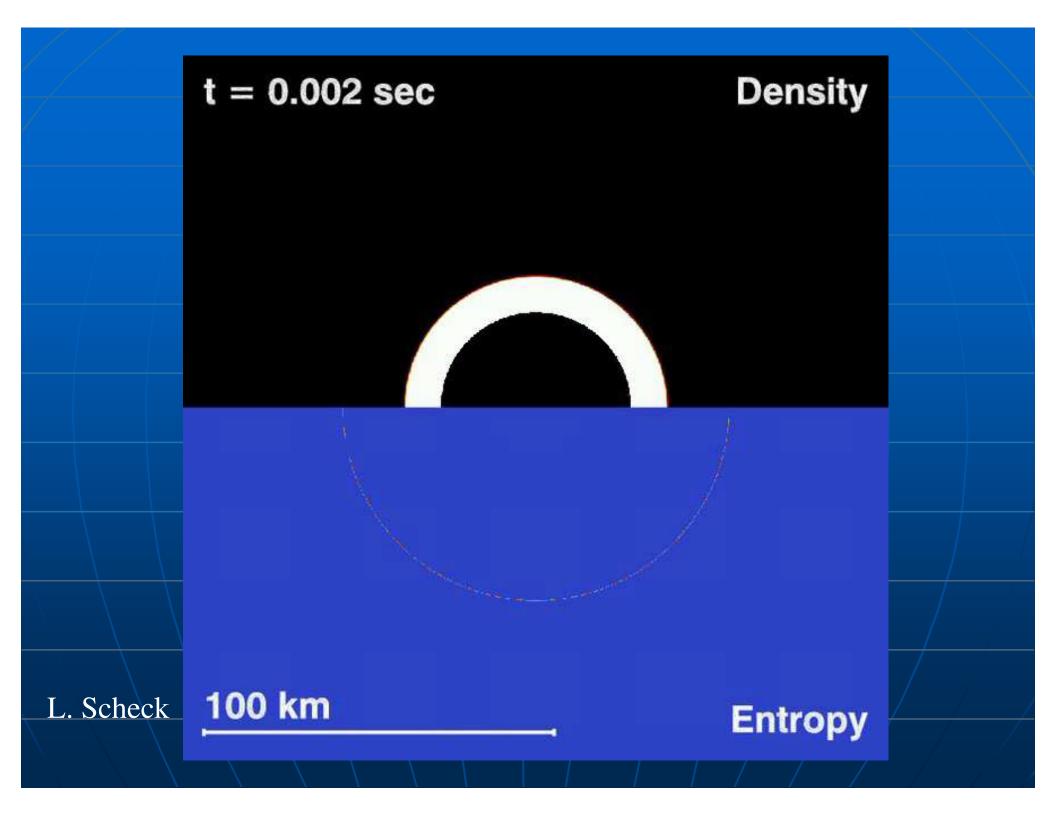
Density distribution 1 sec after core bounce

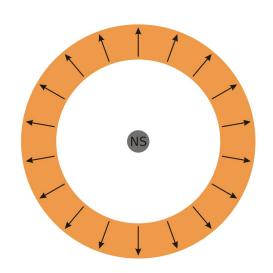


Global dipolar oscillations of the post-shock layer also seen in recent 3D simulations neglecting (Blondin et al. '03) or simplifying (Scheck et al. '04) the treatment of v-transport

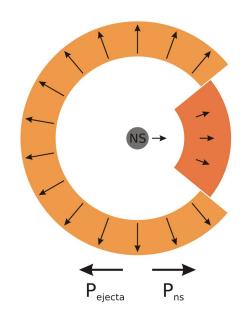


3D core collapse simulation: shock, Y = const & downflow to NS (Scheck 2004)

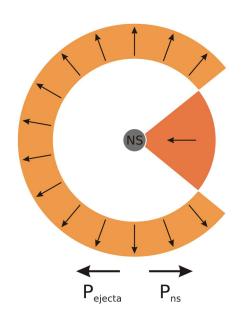




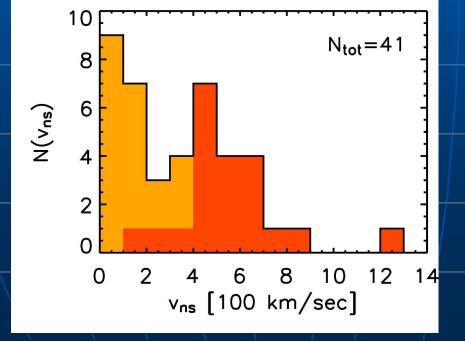
spherical explosion: no kick



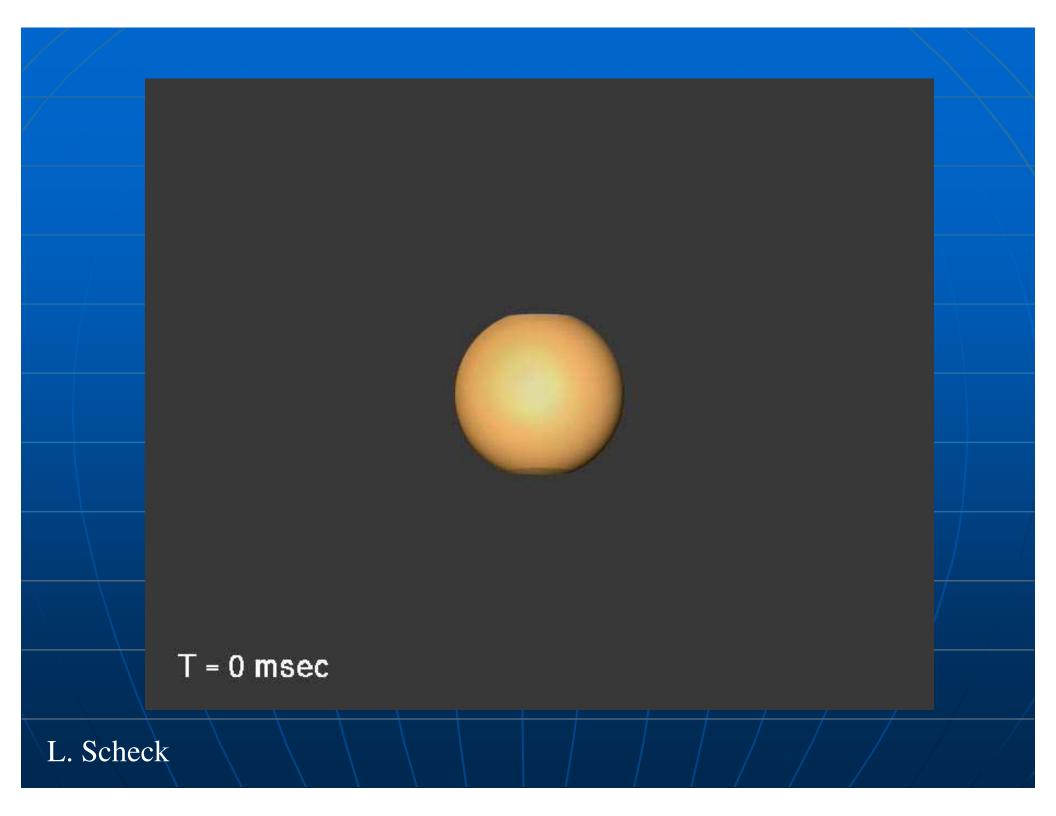
anisotropic explosion: kick due to gravitational acceleration



anisotropic explosion: kick due to anisotropic accretion

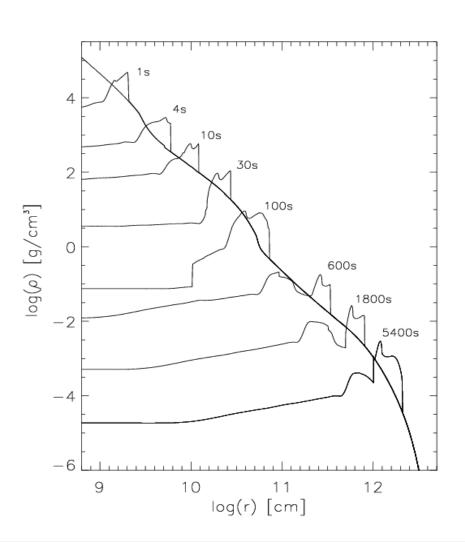


Large set of simulations shows bi-modal kick velocity distribution (Scheck 2005)

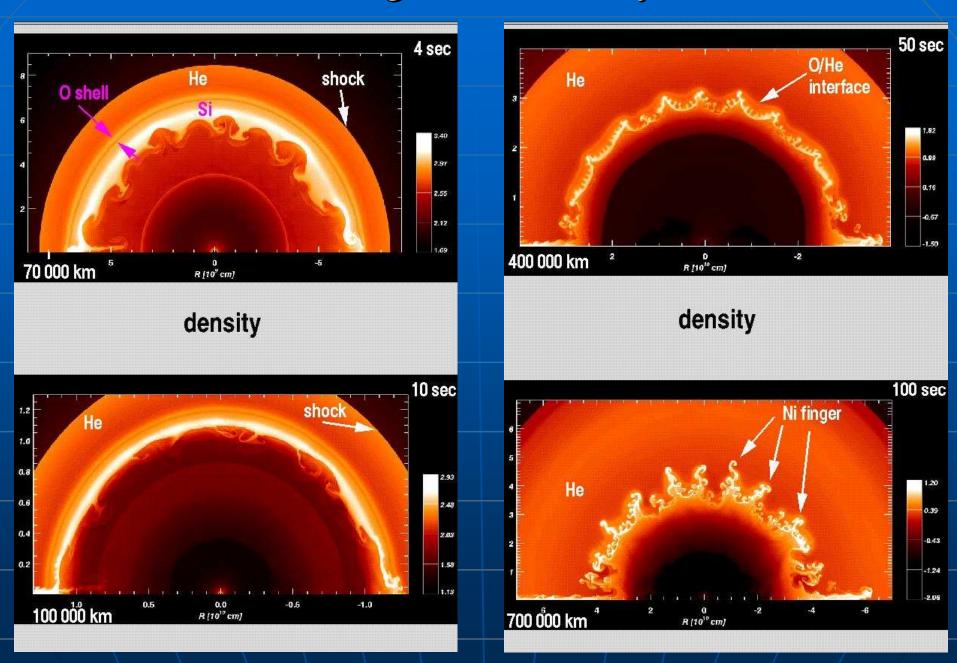


#### Core Collapse Supernovae: Further evolution

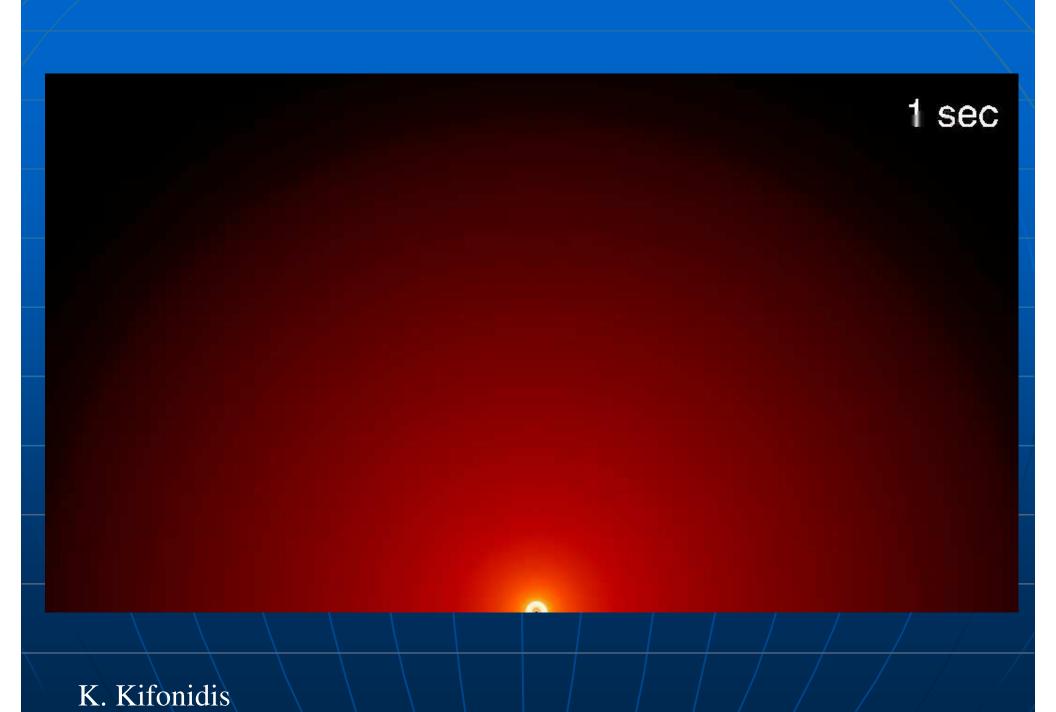
<u>Unsteady</u> shock propagation through stellar envelope --> Rayleigh Taylor unstable regions



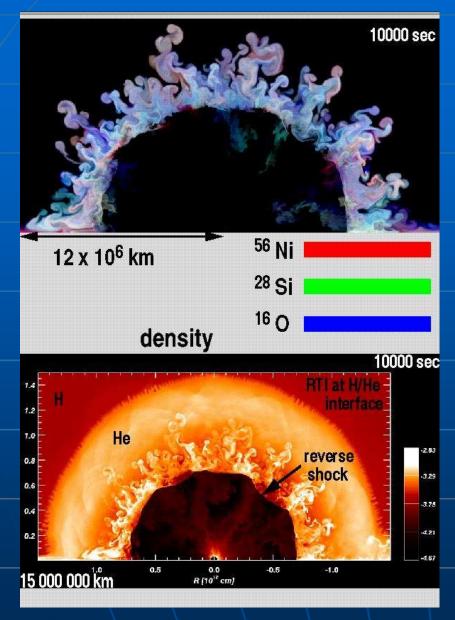
#### Instabilities, mixing and nucleosynthesis



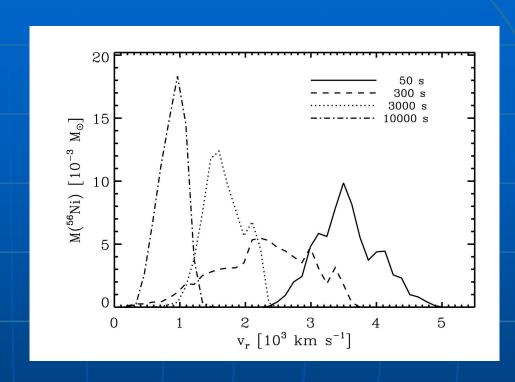
AMR simulation of shock propagation through stellar envelope (Kifonidis, Plewa, Janka & Müller 2003



#### Instabilities, mixing and nucleosynthesis

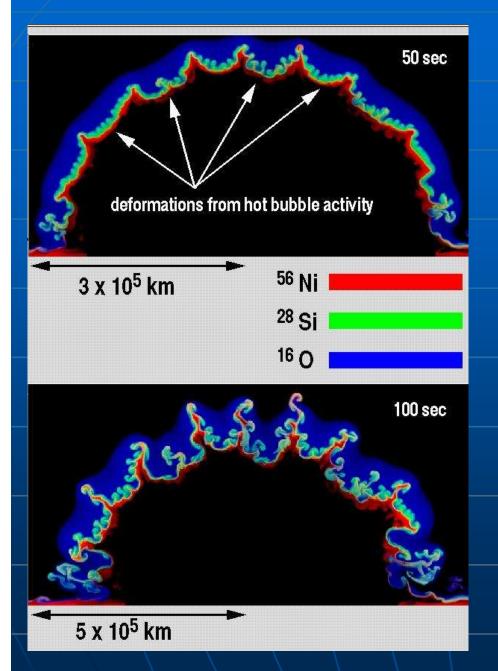


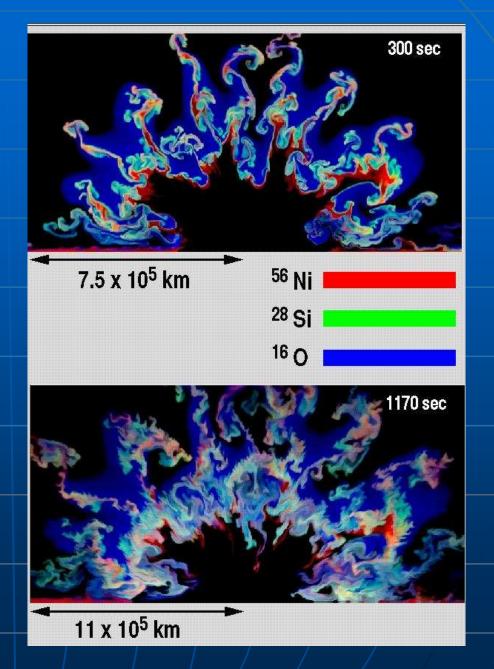
AMR simulation of shock propagation through stellar envelope (Kifonidis, Plewa, Janka & Müller 2003)



- results of simulations in accordance with observations of SNe Ib/Ic
- simulations <u>do not reproduce</u> large velocities of Fe/Ni observed in <u>SN 1987A</u>

#### Instabilities, mixing and nucleosynthesis (cont.)





AMR simulation of shock propagation through stellar envelope (Kifonidis, Plewa, Janka & Müller 2003

#### Summary (Part I)

- q Core-collapse supernova explosions are triggered by neutrino interactions with matter and hydrodynamic instabilities and/or rotation (magnetic fields?).
- q Even the best models available predict weak explosions only.
- q What is the missing physics?
- "Artificially" triggered explosion models predict nuclear abundances in fair agreement with observations.

## Thermonuclear (Type Ia) Supernovae



## Example:

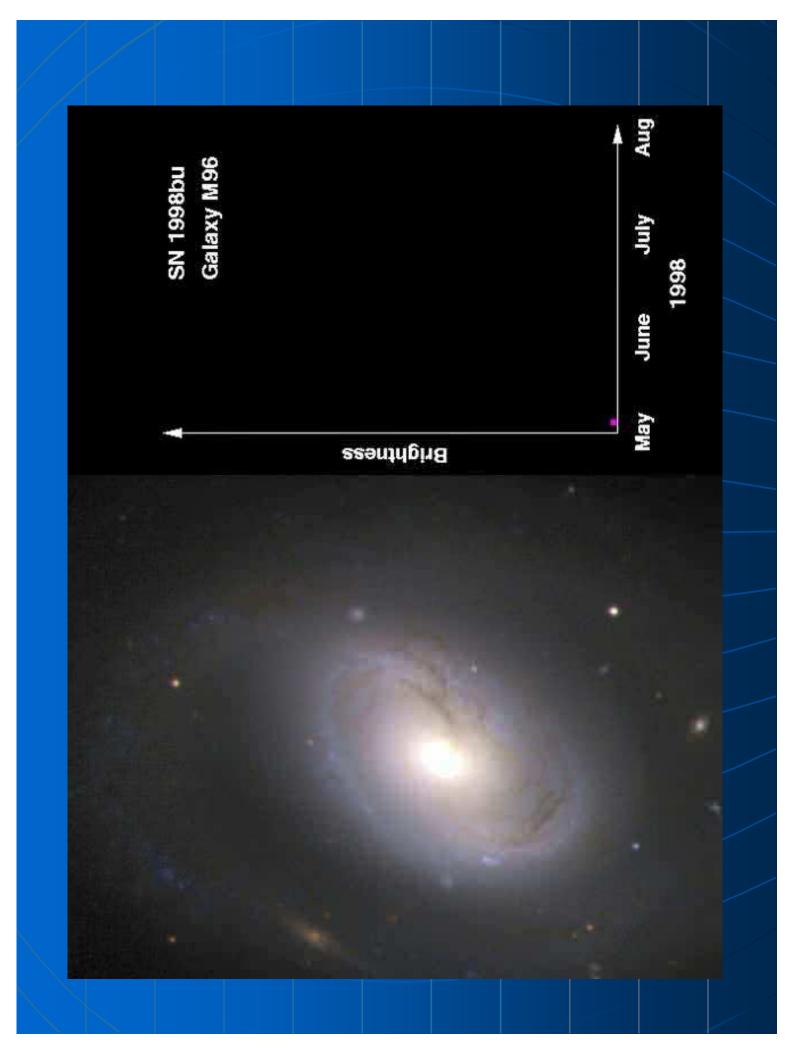
SN 2002bo in NGC 3190;

Discovered: March 9, 2002

B-maximum: March 22, 2002



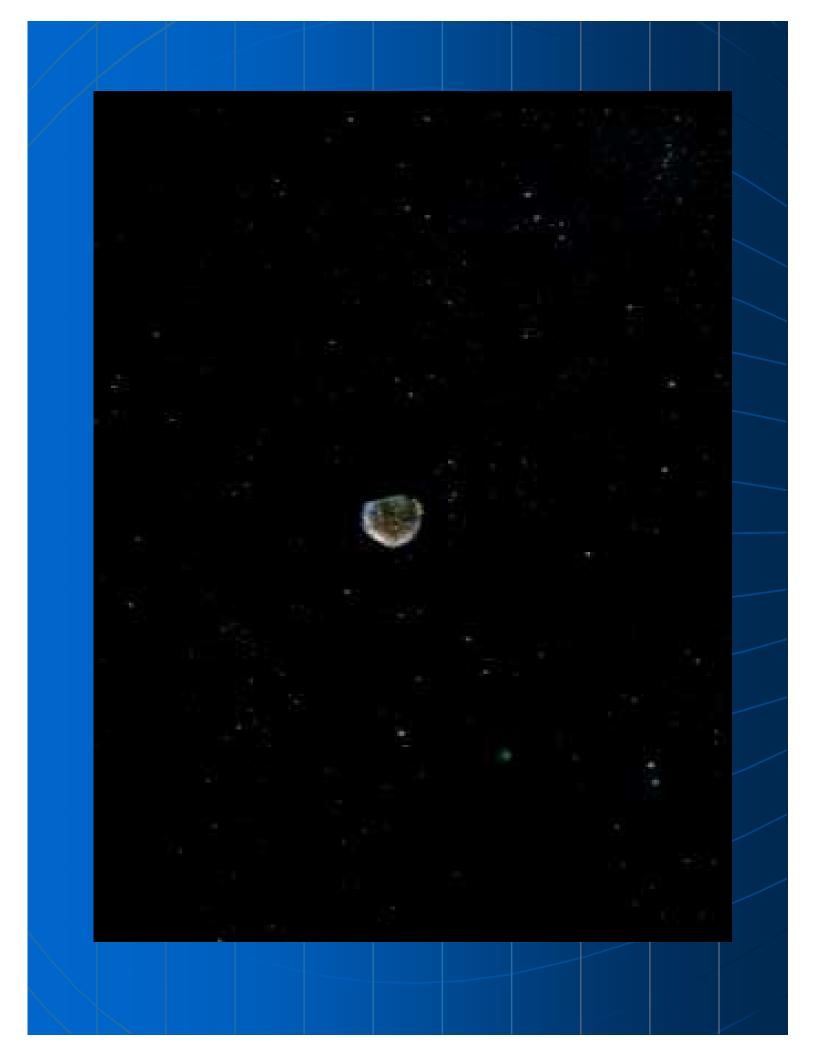




## The "standard model"

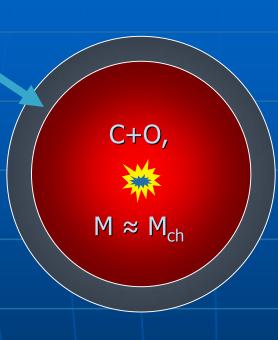


- White dwarf in a binary system
- Growing to the Chandrasekhar mass by mass transfer



## The "standard model"

He (+H) from binary companion



Density  $\sim 10^9 - 10^{10}$  g/cm

Temperature: a few 10<sup>9</sup> K

Radii: a few 1000 km

n Explosion energy:

n Fusion of n C+C, C+O, O+On  $\Rightarrow$  "Fe"

n Laminar burning n velocity:

n  $U_L \sim 100$  km/s  $<< U_S$ 

n

n Too little is burned!

## The physics of turbulent combustion

- Everydays experience:
   Turbulence increases the burning velocity.
- ✓ In a star:Reynoldsnumber ~ 10<sup>14</sup>!
- Physics of thermonuclear burning is very similar to premixed chemical flames.



#### A couple of definitons:

Kolmogorov (length) scale

$$\eta := (v^3/\epsilon)^{1/4}$$

(Turbulent) Reynolds number

$$Re := v'/s_L \cdot 1/l_F$$

(Turbulent) Damköhler number

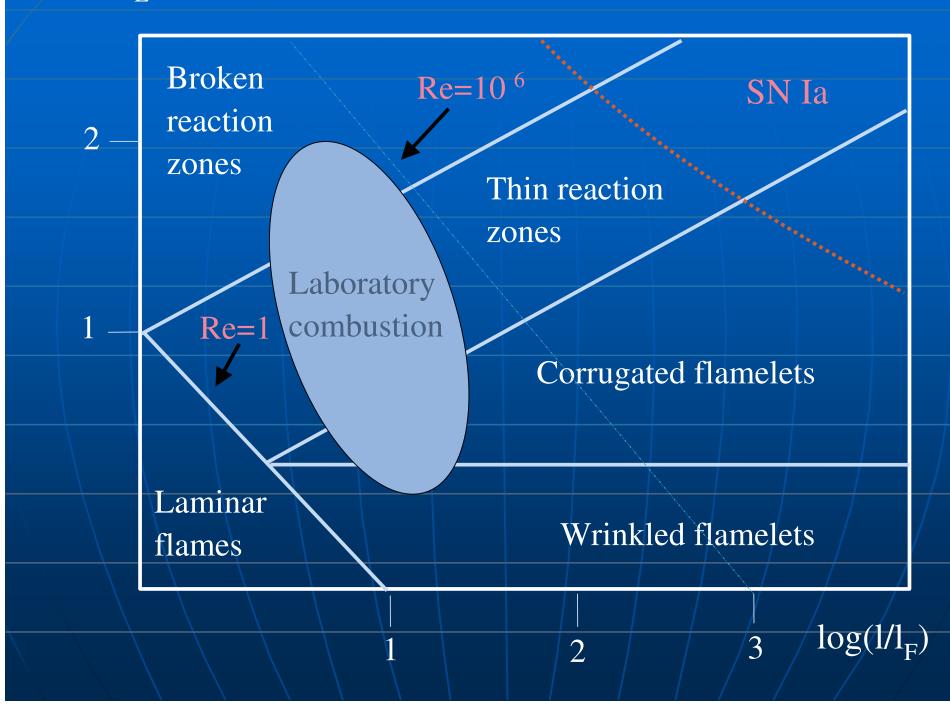
$$Da := s_L/v' \cdot 1/l_F$$

(Turbulent) Karlovitz number

$$Ka := l_F^2/\eta^2$$

$$\Rightarrow$$
 Re = Da<sup>2</sup> · Ka<sup>2</sup>

#### $\log(v'/s_L)$



## Burning regimes of pre-mixed flames

#### 1. Cellular burning, wrinkled flamelets

$$u_{cell} = s_L [1+\epsilon(\mu)] ; \mu = \rho_b/\rho_u ,$$
 
$$\epsilon(\mu) \approx 0.41 (1-\mu)^2$$

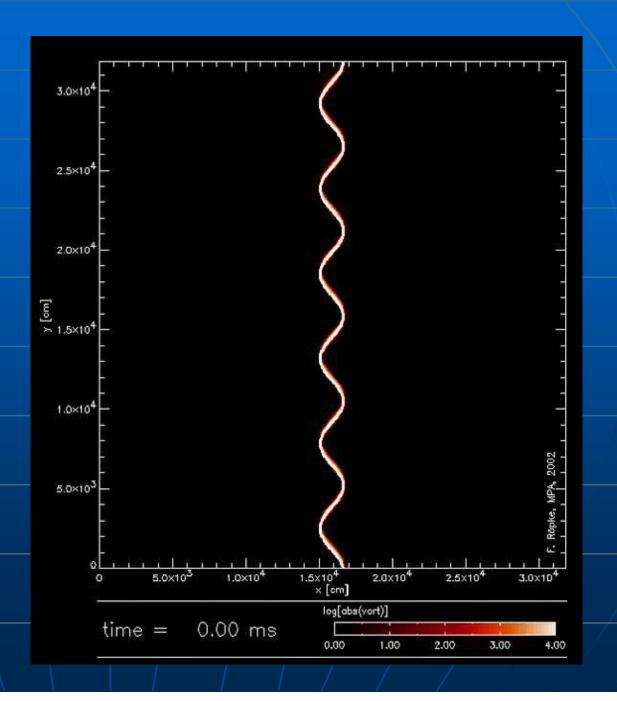
Or: "Fractal model"

$$u_{cell}(1) = s_L (1/l_{crit})^{D-1}$$

The Landau-Darrieus instability and its interaction with turbulence:

Quiescent fuel

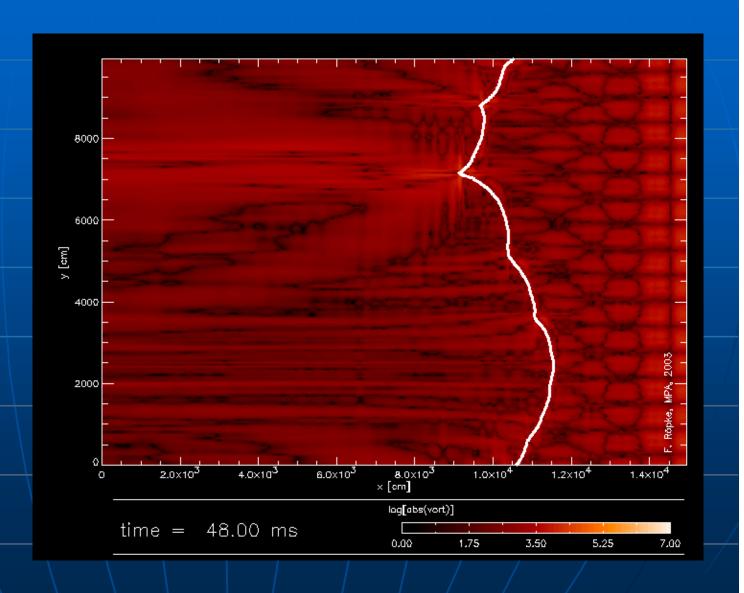
(Röpke et al., 2003a)



# The Landau-Darrieus instability and its interaction with turbulence:

Weak vortical flow

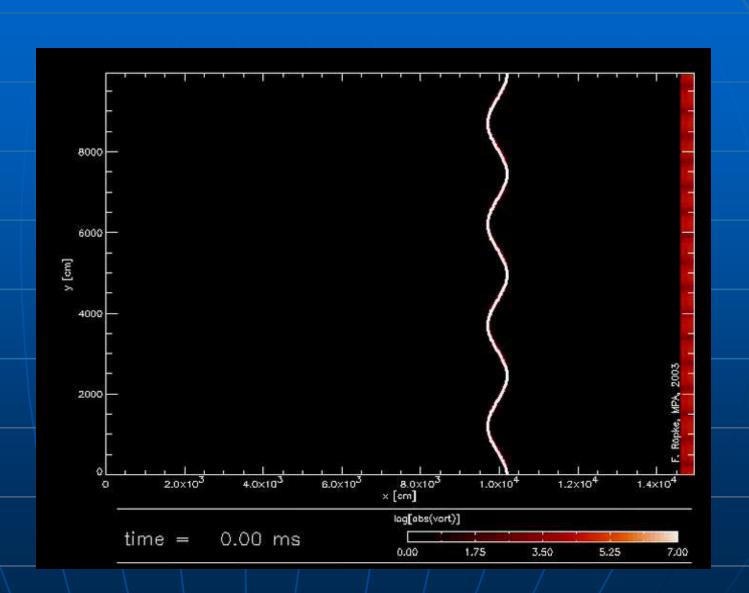
(Röpke et al., 2003b)



# The Landau-Darrieus instability and its interaction with turbulence:

Strong vortical flow

(Röpke et al., 2003b)



## Burning regimes of pre-mixed flames

#### 2. The corrugated flamelet regime

Transition at the "Gibson scale":

$$v(l_{Gibs}) = u_{cell}(l_{Gibs})$$

In the limit of strong turbulence:

$$s_{turb}$$
 (1)  $\approx$  v'(1), 1 >  $l_{Gibs}$  (independent of  $s_L!!!$ )

$$d_{turb} \approx 1$$
 ("turbulent flame brush")

#### Fully developed turbulence?

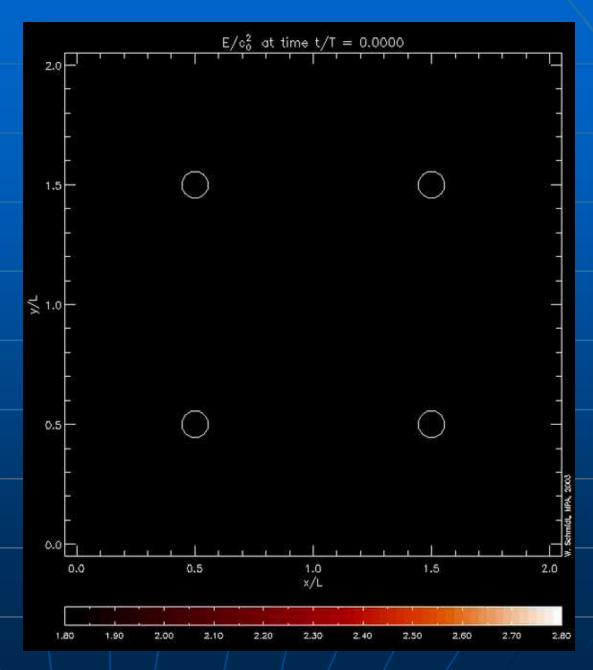
3-D "direct"
numerical simulations
of flames moving in
white dwarf matter:
Energy

$$\rho = 2.9 \cdot 10^9 \, \text{gcm}^{-3}$$

$$V/s_{lam} = 4$$

$$V/c_0 = 0.043$$

(Schmidt et al., 2004)



#### Fully developed turbulence?

3-D "direct"
numerical simulations
of flames moving in
white dwarf matter:
Vorticity

$$\rho = 2.9 \cdot 10^9 \, gcm^{-3}$$

$$V/s_{lam} = 4$$

$$V/c_0 = 0.043$$

(Schmidt et al., 2004)

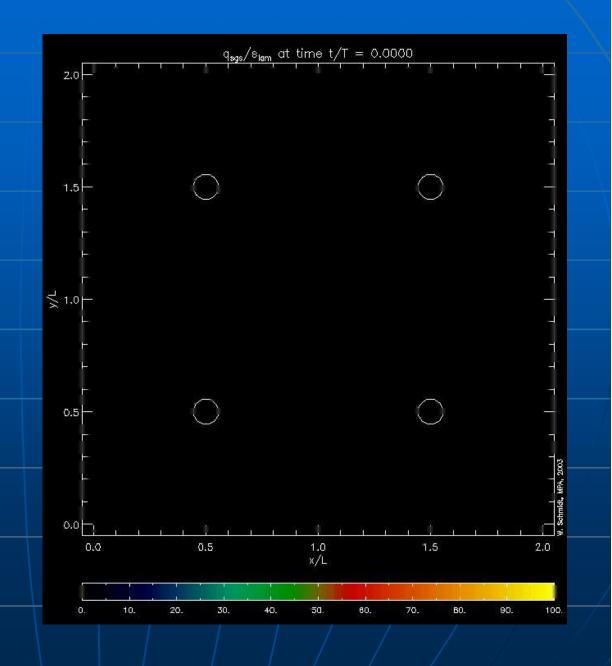
#### Fully developed turbulence?

3-D "direct"
numerical simulations
of flames moving in
white dwarf matter:
Subgridscale energy

$$\rho = 2.9 \cdot 10^9 \, gcm^{-3}$$

$$V/s_{lam} = 4$$

$$V/c_0 = 0.043$$



(Schmidt et al., 2004)

## Burning regimes of pre-mixed flames

#### 3. The distributed-burning

Turbulent eddies interact with the flame:

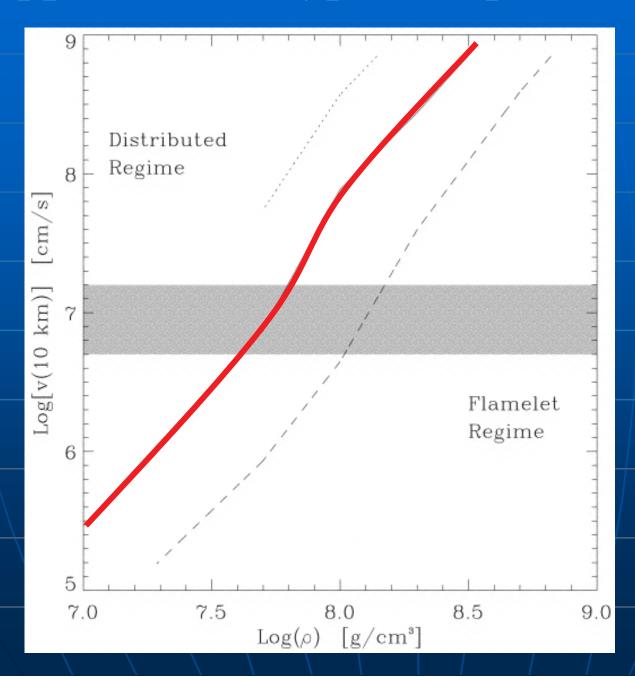
$$l_F \ge l_{Gibs}$$

Rough estimate ("Damköhler scaling"):

$$s_{turb}/s_L \approx const (D_t/D)^{1/2}$$
 (dependent on  $s_L !!!!$ )

$$const = O(1)$$

## Application to type Ia supernova



Niemeyer & Woosley (1997)

## Burning regimes of pre-mixed flames

#### 4. The Rayleigh-Taylor regime

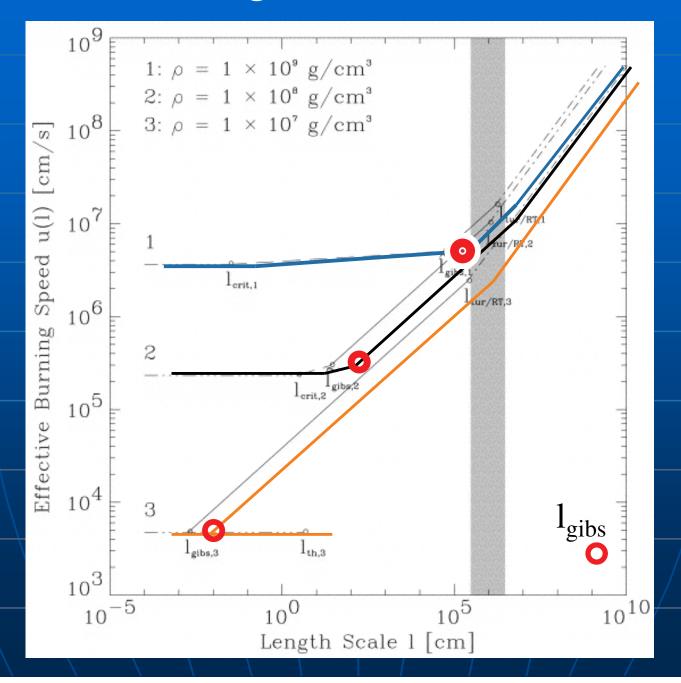
$$v_{RT} = B \sqrt{(g_{eff} 1)}$$
;  $B \approx 0.5$ ;  $g_{eff} = At \cdot g$ 

Sharp-Wheeler model:

$$r_{sw} \approx 0.05 g_{eff} t^2$$
;  $v_{sw} \approx 0.1 g_{eff} t$ ;

$$1_{\text{tur/RT}} \approx 10^6 \text{ cm}$$

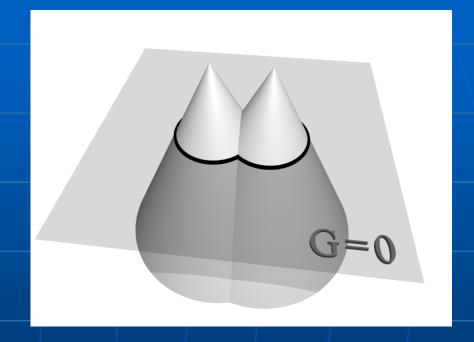
## Effective burning velocities in SN Ia



Niemeyer & Woosley (1997)

#### How to model thermonuclear flames?

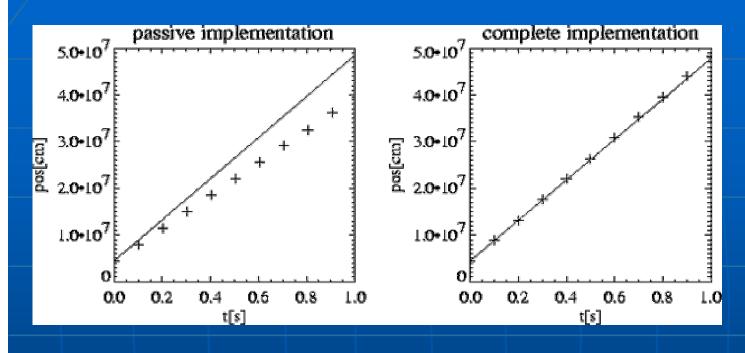
- The "flames" cannot be resolved numerically.
- The amplitutes of turbulent velocity fluctuations in the length scale of the flame are determined on the integral scale.



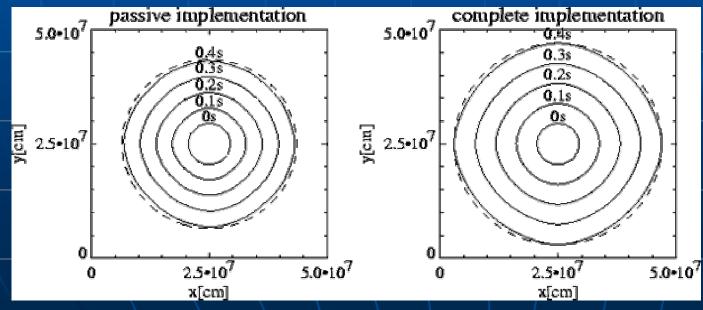
$$\partial G/\partial t = -\mathcal{D}_f \nabla G$$

$$\mathcal{D}_f = \mathbf{v_u} + \mathbf{s_{tur}} \mathbf{n}; |\nabla G| = 1$$

#### Some test of the code



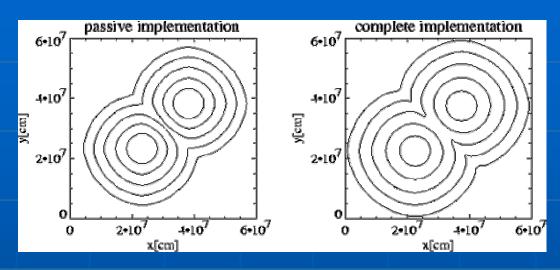
Planar flame



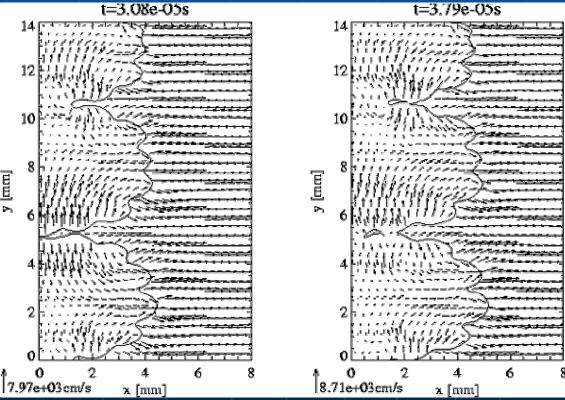
Circular flame

Reinecke et al. (1999)

## Some test of the code (ctn.)



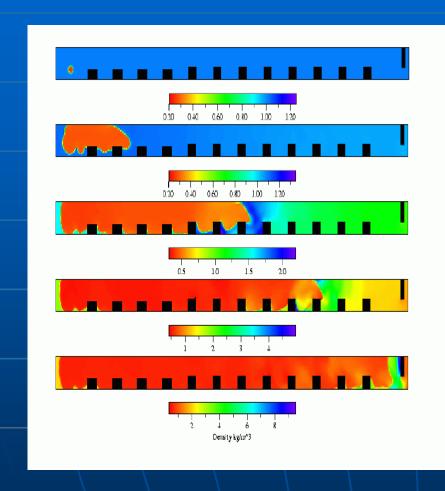
Merging circular flames

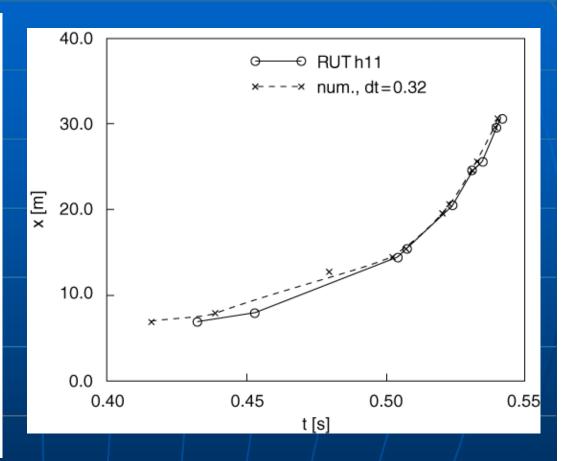


Hydrogen-in-air flames

Reinecke et al. (1999)

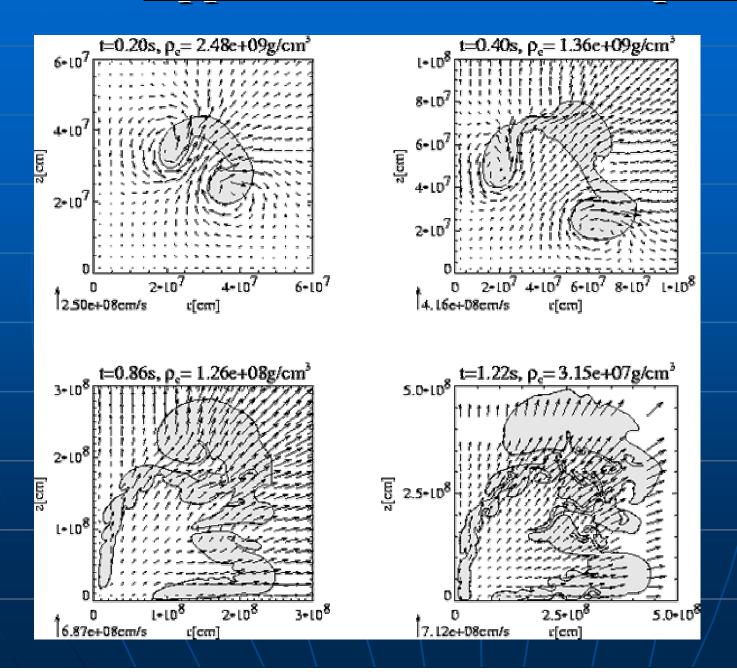
#### Other laboratory flames





The method can reproduce terrestrial experiments well! (Smiljanowski et al. 1997)

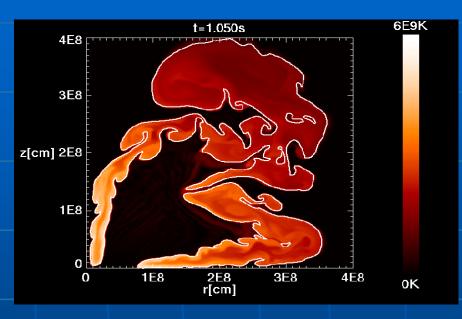
### Application to the SN Ia problem

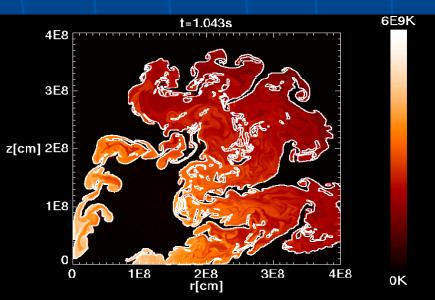


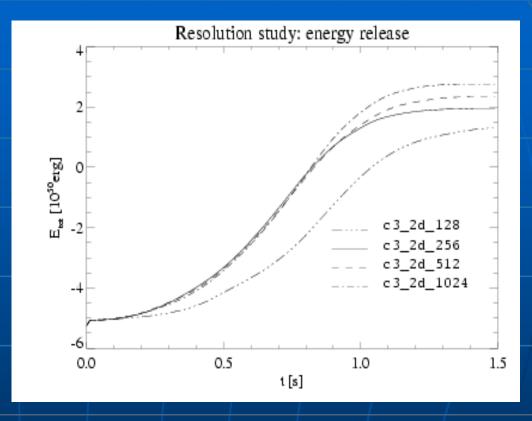
One rising blob (in 2D)

Reinecke et al. (1997)

#### Convergence tests in 2D

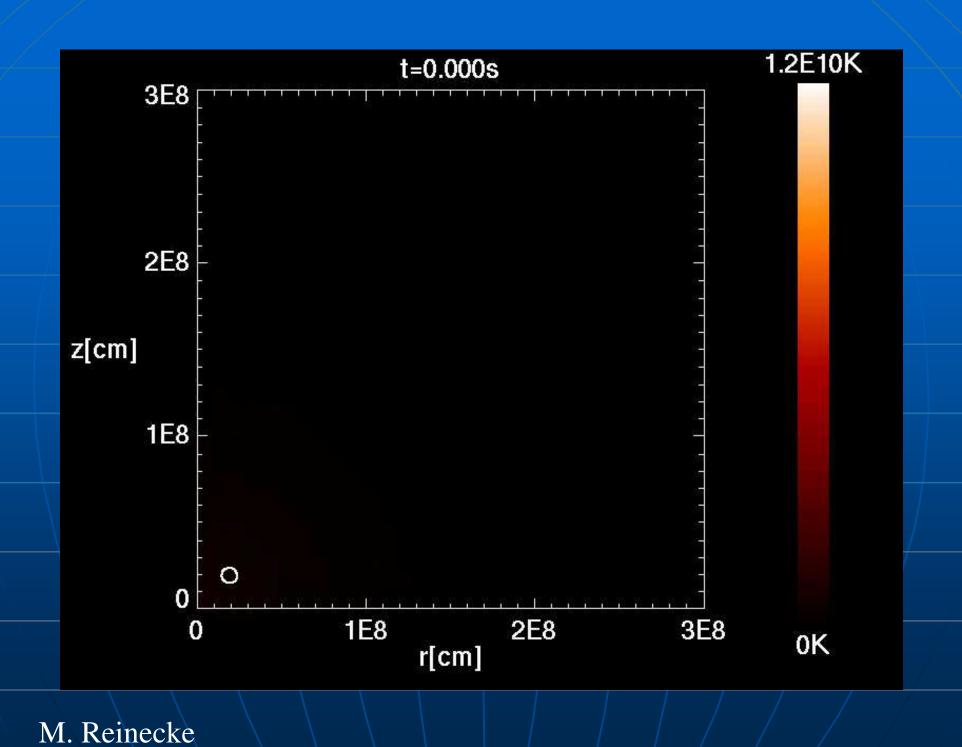




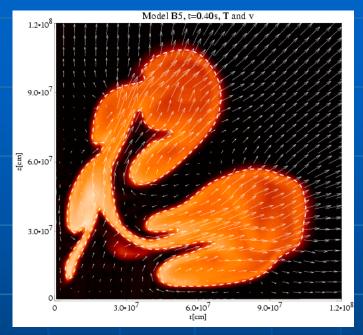


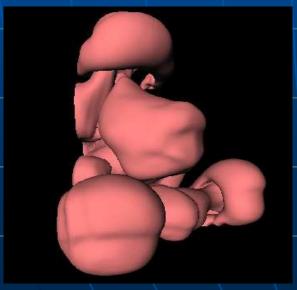
Global results are independent of the numerical resolution!

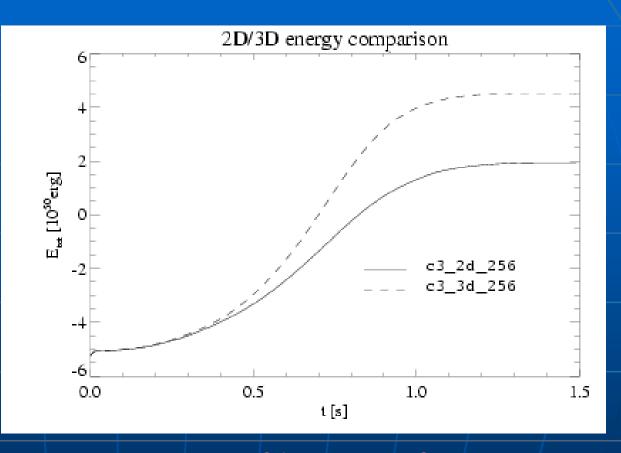
Reinecke et al. (1999, 2002)



#### $2D \Rightarrow 3D$



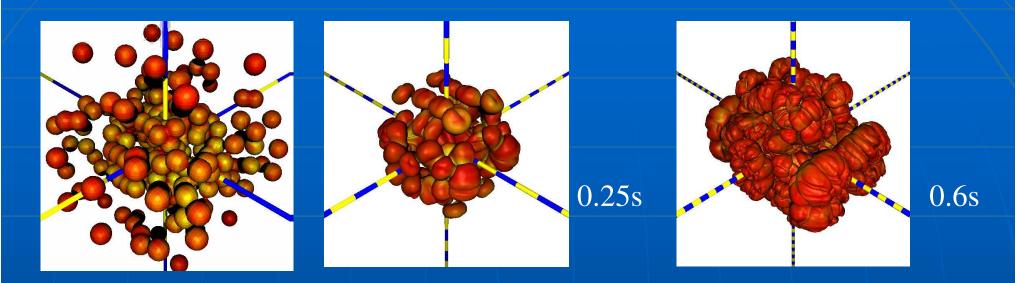


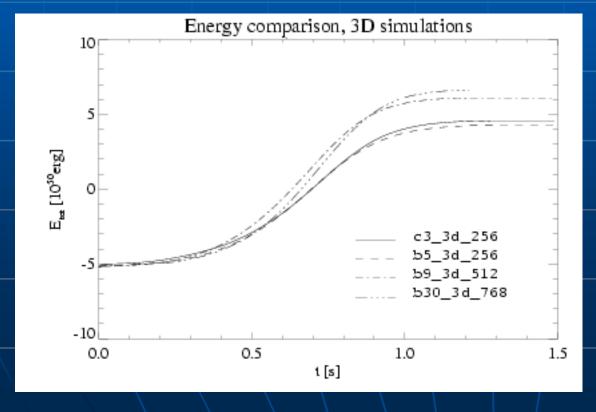


# Because of larger surface area: More energy is produced!

Reinecke et al. (2001) (See also Gamezo et al., 2003)

#### 3D models: The best we could do until now!

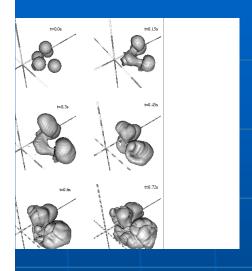




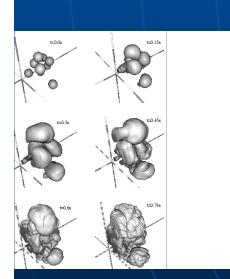
Mod b30\_3d

(Reinecke et al., 2003)

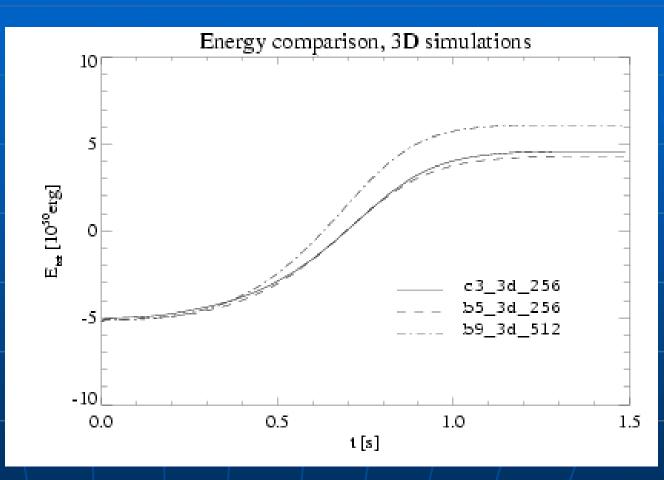
# Modeling Flames in 3D: Dependence on initial conditions?



Mod. b5\_3d



Mod. b9\_3d



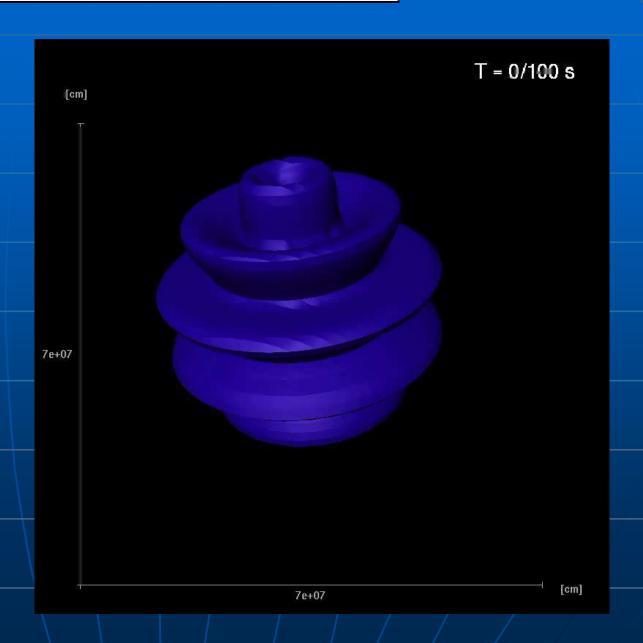
Moderate dependence on initial conditions!

(Reinecke et al., 2002)

## Recent modifications of the code:

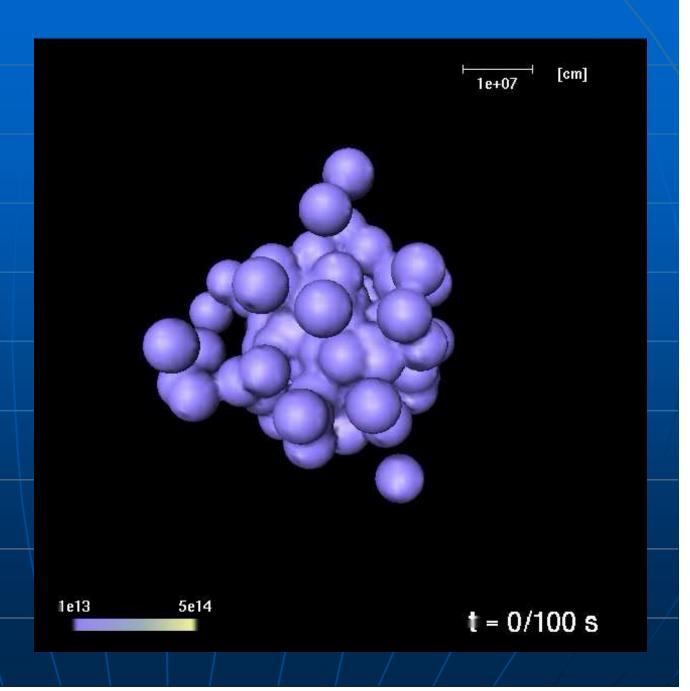
1. Moving grid

Röpke (2004)



#### 2. Full star (" $4\pi$ ")

Röpke & Hillebrandt (2004)



# Some open questions and challenges

#### q Ignition conditions:

How do WDs reach M<sub>Ch</sub>? Center/off-center ignition? One/multiple "points"?

#### q Combustion modeling:

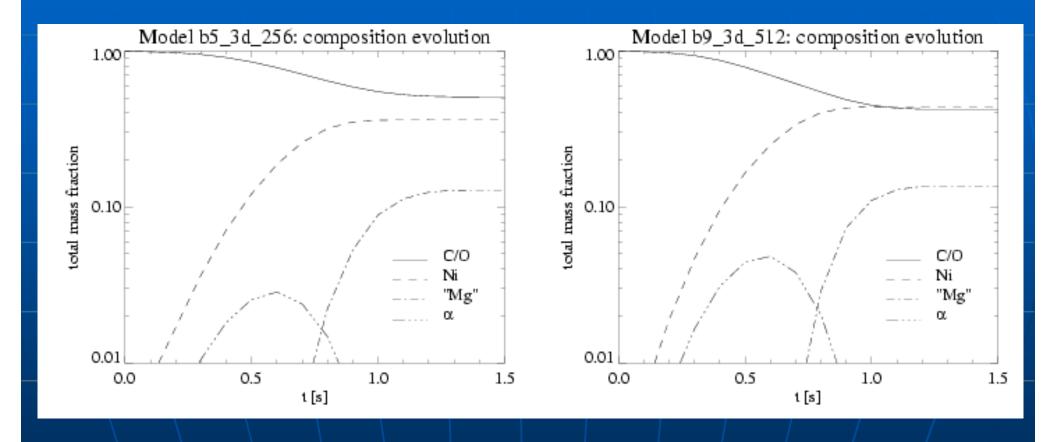
Sub-grid scale models? Interaction of nuclear flames with turbulence; "distributed burning"; "active turbulent combustion"?

Deflagration/detonation transition?

#### q "Full-star" models:

Composition? Rotation? Light curves? Spectra?

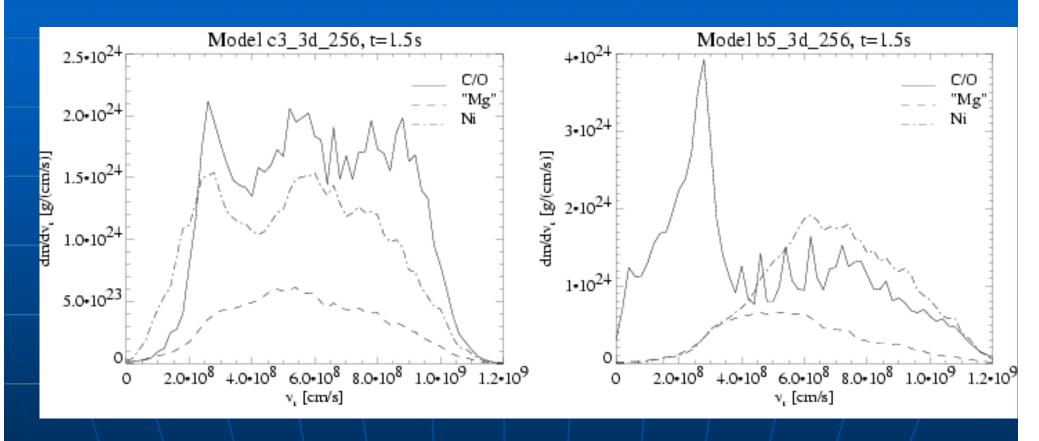
#### Observable Predictions: Chemical composition?



Significant amounts of unburned C and O!

(Reinecke et al., 2002)

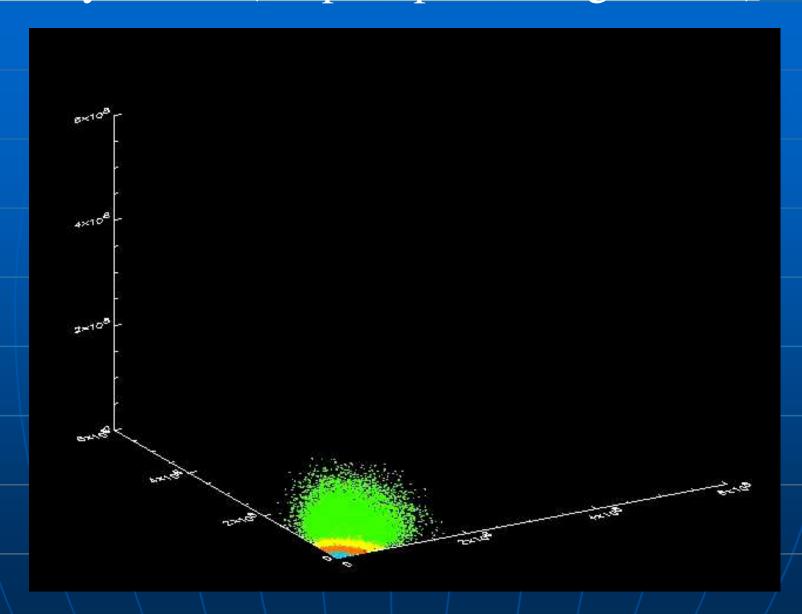
# Observable Predictions: Chemical composition in velocity space?



Velocity distribution sensitive to ignition conditions!

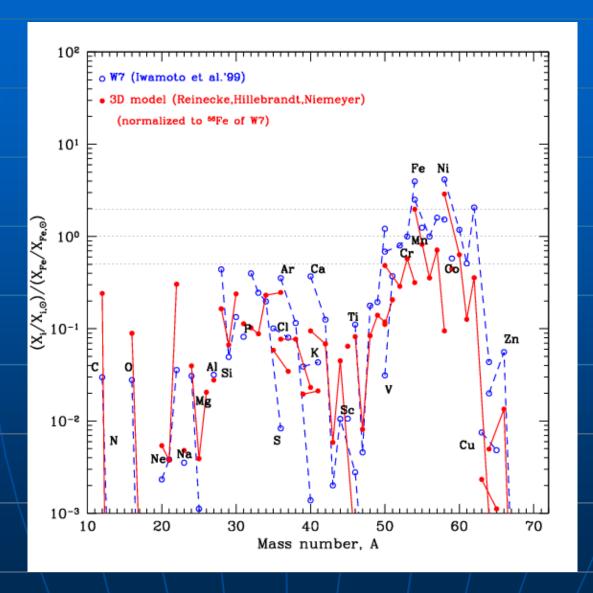
(Reinecke et al., 2002)

## Nucleosynthesis (in 'post-processing' mode)



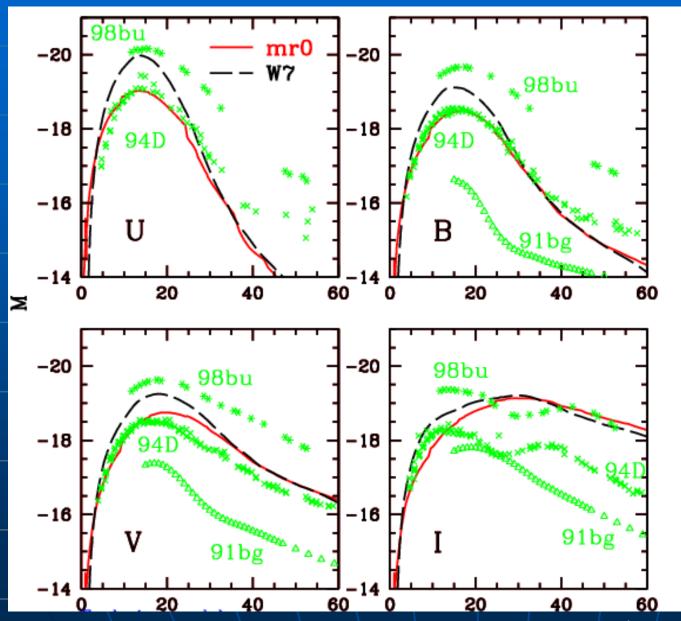
(Travaglio et al., 2004)

#### Chemical composition: What is different from W7?



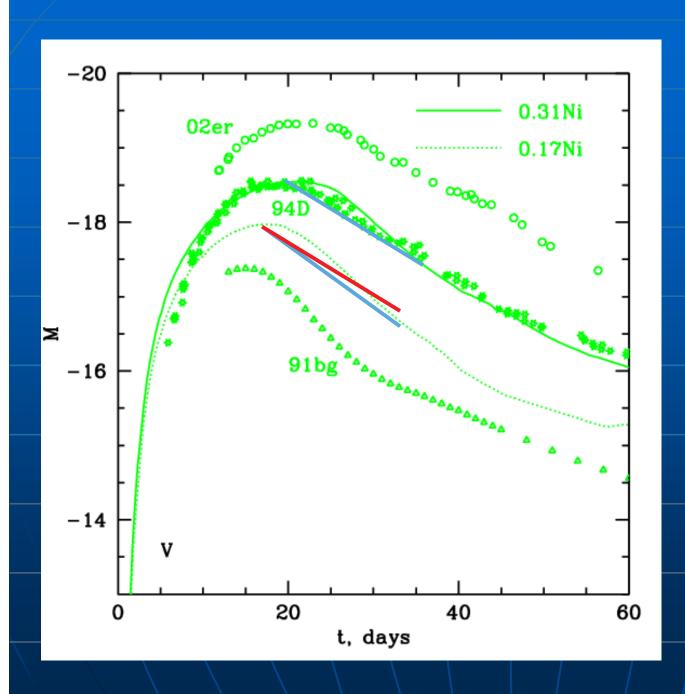
(Off-center ignited 3D model, Travaglio et al., 2003)

#### Colour light curves .....



Sorokina & Blinnikov (2002)

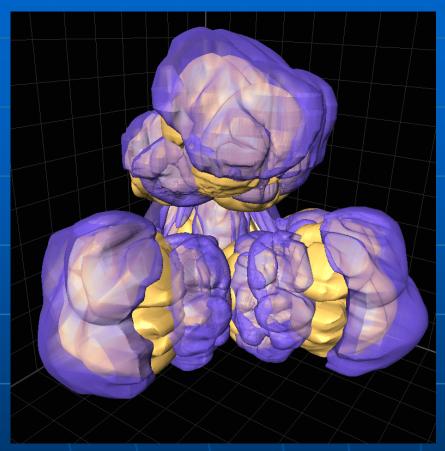
... and model predictions.



#### <u>Prediction from</u> <u>Theory :</u>

Light-curve shape / luminosity correlation?

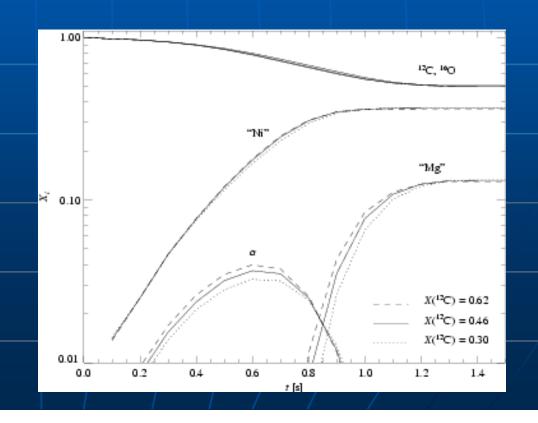
# Dependence on the initial C/O ratio?

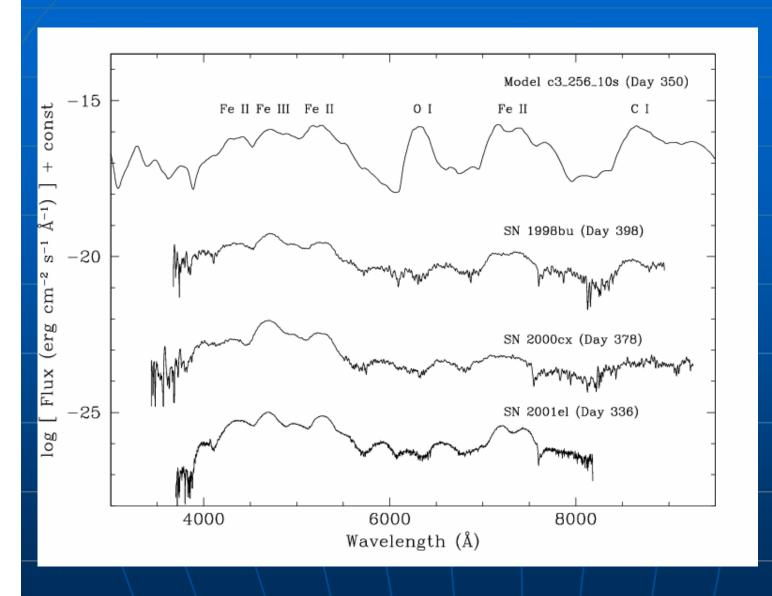


X(12C)	E <sub>nuc</sub> (10 <sup>50</sup> erg)	$M(Ni) (M_{\circ})$	$M_{lpha}^{ m max} \ (M_{\circ})$
0.30	8.85	0.5178	0.0458
0.46	9.46	0.5165	0.0518
0.62	9.97	0.5104	0.0564

Ni-mass (luminosity)
independent of initial C/O!

(Röpke & Hillebrandt, 2004)





<u>But:</u>
<u>Nebular</u>
<u>spectra!</u>

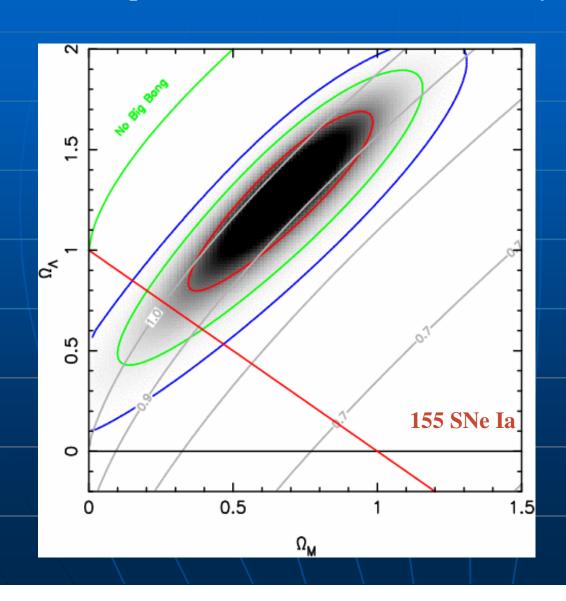
Too much low-velocity oxygen!

# Summary (Part II)

- "" "Parameter-free" thermonuclear models of type Ia supernovae, based on Chandrasekhar-mass C+O white dwarfs explode with about the right energy.
- n They allow to predict light curves and spectra, depending on physical parameters!
- n They can explain (most of?) the observed properties well.
- n The diversity may be due to randomness in the ignition conditions, (C/O), and metallicity.

# Supernovae and Cosmology: The Quest for Precise Luminosity Distances

(In part "borrowed" from a lecture by Bruno Leibundgut)



(Tonry et al. 2003)

#### Distances in the local universe

q Assume a linear expansion ("Hubble law")

$$v = cz = H_0 \cdot D$$

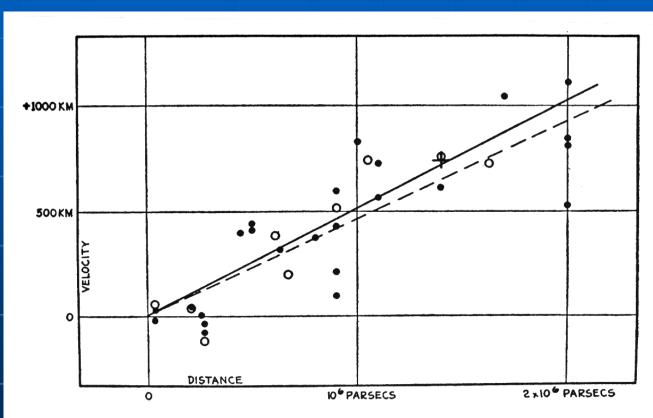
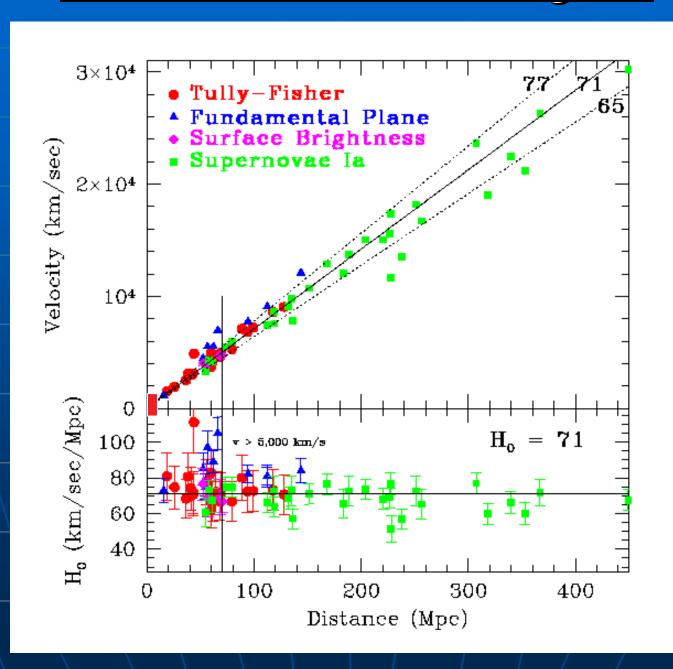
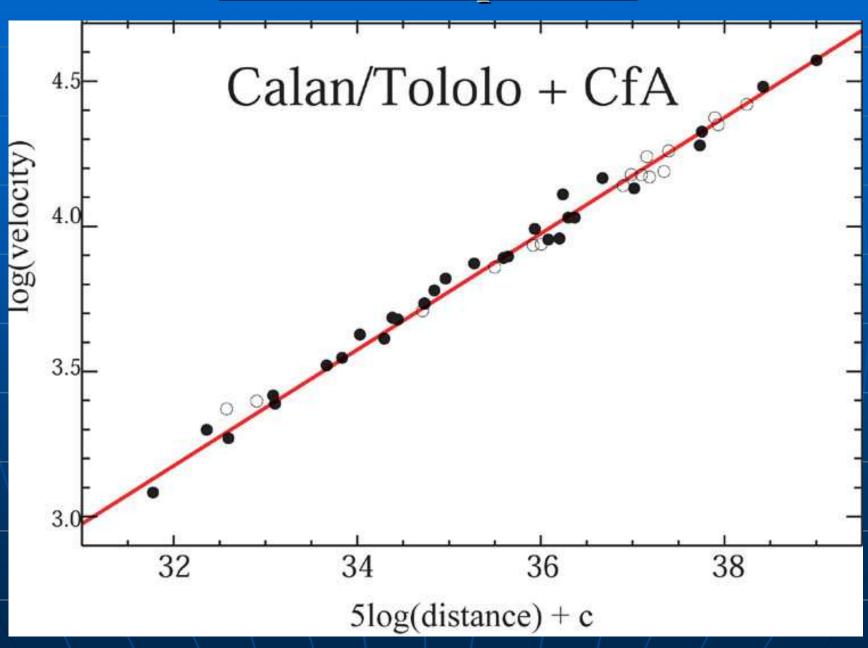


Fig. 9. The Formulation of the Velocity-Distance Relation.

#### A "modern" Hubble diagram



### Universal expansion



#### Distances in the local universe

n Assume a linear expansion (Hubble law):

$$v = cz = H_0 \cdot D$$

- n Use the distance modulus
  - $m-M=5\log(D/10pc)-5$
- n Distances of a 'standard candle' (M=const.)

$$m = 5\log(z) + b$$

$$b = M + 25 + 5\log(c) - 5\log(H_0)$$

#### The Hubble constant

- n Sets the absolute scale of cosmology
  - replaces these annoying h's in all the theorists talks
- Measure redshifts and distances in the nearby universe
  - Supernovae can do this in two ways:
    - n Expanding photosphere method of core-collapse SNe
    - n accurate (relative) distances from SN Ia

### Expanding Photosphere Method

n Baade (1926), Schmidt et al. (1993), Eastman et al. (1996), Hamuy et al. (2001)

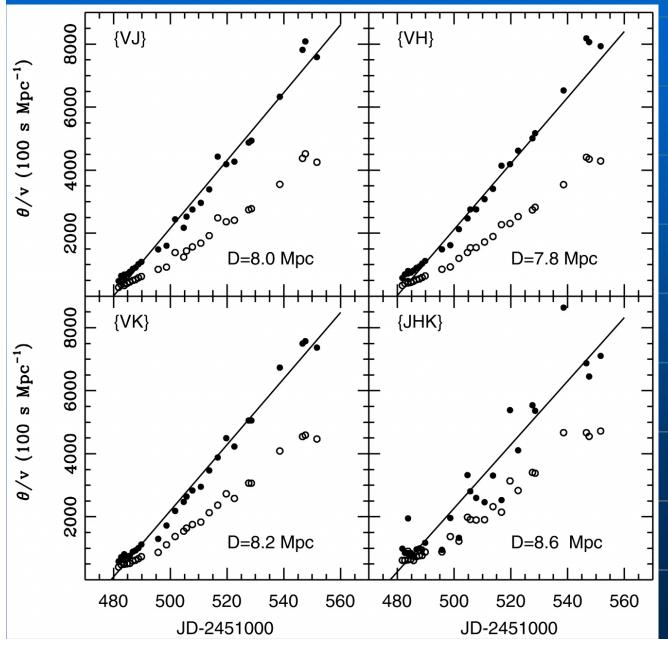
ø Assume homologous expansion:  $R(t) = R_0 + v(t-t_0)$ 

ø Photometric angular diameter

$$\Theta = \frac{R}{D} = \sqrt{\frac{f_{\lambda}}{\zeta_{\lambda}^{2} \pi B_{\lambda}(T) 10^{-0.4A(\lambda)}}}$$

#### Distances from EPM

(SN 1999em, Hamuy et al. 2001)



$$\frac{\Theta_i}{v_i} \approx \frac{t_i - t_0}{D}$$

**Slope** gives the distance

Intercept the size of the progenitor and/or time of explosion

# Distances from EPM

- Note that this distance measurement is completely independent of any other astronomical object!
  - no distance ladder
- n Assumption:
  - massive envelope that creates a photosphere
  - spherical symmetry
    - Ł not true for many core collapse supernovae
  - correction factors for deviation from black body spectrum
    - E model dependent

#### EPM so far

#### n Limitations

- needs large and extensive data sets
- difficulties to get into the Hubble flow
- distances only to galaxies with supernovae
  - n difficult to build large sample

#### n Promise

- completely independent distance measurements
  - n checks on the Cepheid distance scale

#### Distances with Type Ia Supernovae

- Use the Hubble diagram  $(m-M \ vs. \log z)$ Ø  $m-M=5\log(z)+25+5\log(c)-5\log(H_0)$
- n Note that the slope is given here.

Hubble constant can be derived when the absolute luminosity *M* is known

$$\emptyset \log H_0 = \log(z) + 5 + \log(c) - 0.2(m-M)$$

#### Hubble constant from SNe Ia

- n Calibrate the absolute luminosity
  - through Cepheids
    - n 'classical distance ladder'
      - depends on the accuracy of the previous rungs on the ladder
      - LMC distance, P I Grelation, metallicities
    - n HST program (Sandage, Tammann)
    - n HST Key Programme (Freedman, Kennicutt, .....)
  - through models
    - n extremely difficult

# Absolute Magnitudes of SNe Ia

SN	Galaxy	m-M	M <sub>B</sub>	$M_V$	M <sub>I</sub>	$\Delta m_{15}$
1937C	IC 4182	28.36 (12)	-19.56 (15)	-19.54 (17)	-	0.87 (10)
1960F	NGC 4496	31.03 (10)	-19.56 (18)	-19.62 (22)	-	1.06 (12)
1972E	NGC 5253	28.00 (07)	-19.64 (16)	-19.61 (17)	-19.27 (20)	0.87 (10)
1974G	NGC 4414	31.46 (17)	-19.67 (34)	-19.69 (27)	-	1.11 (06)
1981B	NGC 4536	31.10 (12)	-19.50 (18)	-19.50 (16)	-	1.10 (07)
1989B	NGC 3627	30.22 (12)	-19.47 (18)	-19.42 (16)	-19.21 (14)	1.31 (07)
1990N	NGC 4639	32.03 (22)	-19.39 (26)	-19.41 (24)	-19.14 (23)	1.05 (05)
1998bu	NGC 3368	30.37 (16)	-19.76 (31)	-19.69 (26)	-19.43 (21)	1.08 (05)
1998aq	NGC 3982	31.72 (14)	-19.56 (21)	-19.48 (20)	-	1.12 (03)
Straight mean			-19.57 (04)	-19.55 (04)	-19.26 (0 6	5)
Weighted mean			-19.56 (07)	-19.53 (06)	-19.25 (0 9	)

(Saha et al. 1999)

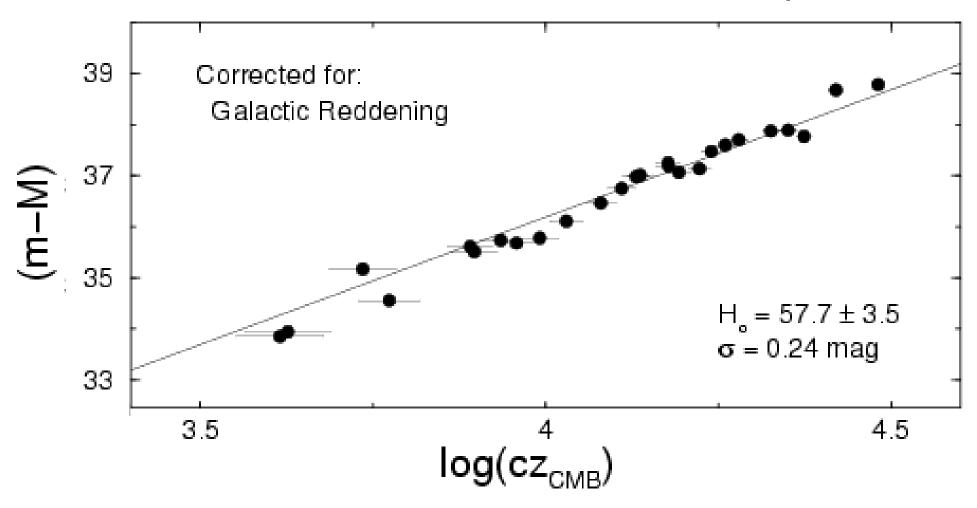
# Testing the SNe Ia as distance indicators

- Hubble diagram of SNe Ia in the local, linear expansion, Hubble flow
- n Calibration through "primary" distance indicators
- n Theoretical models

#### Nearby SNe Ia

Phillips et al. (1999)





# <u>Light curve shape – luminosity</u>

#### $_{\rm n}$ $\Delta m_{15}$ relation

Phillips (1993), Hamuy et al. (1996), Phillips et al. (1999)

#### n MLCS

Riess et al. (1996, 1998), Jha et al. (2003)

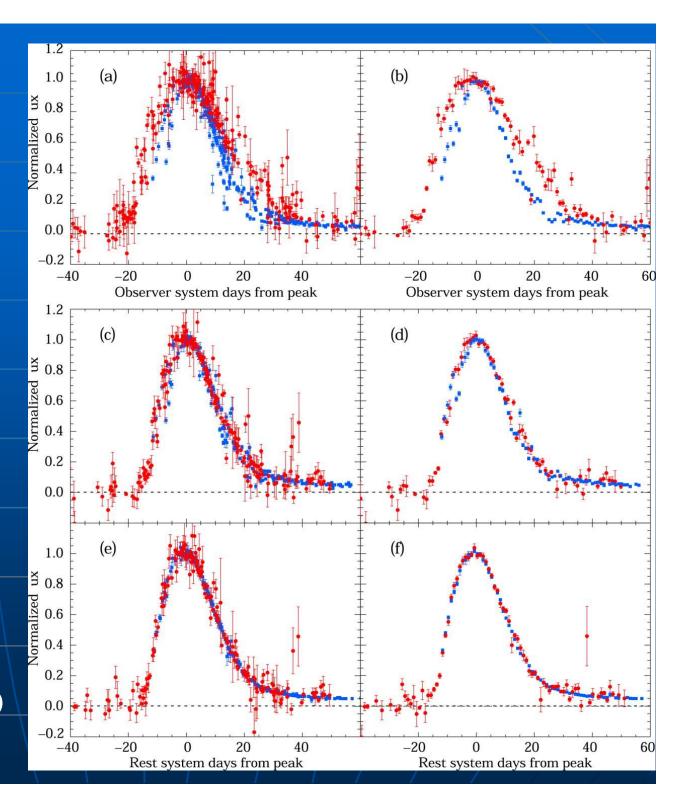
#### n stretch

Perlmutter et al. (1997, 1999), Goldhaber et al. (2001)

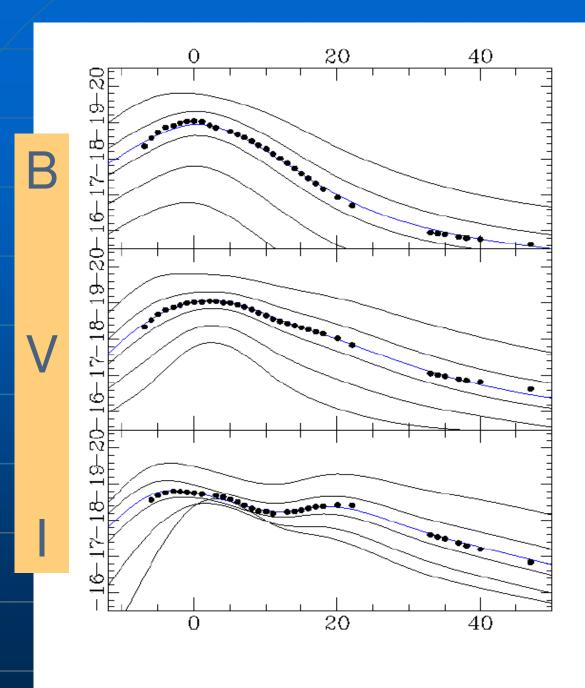
#### n MAGIC

Wang et al. (2003)

# The principles of the calibration



(Goldhaber et al. 2001)

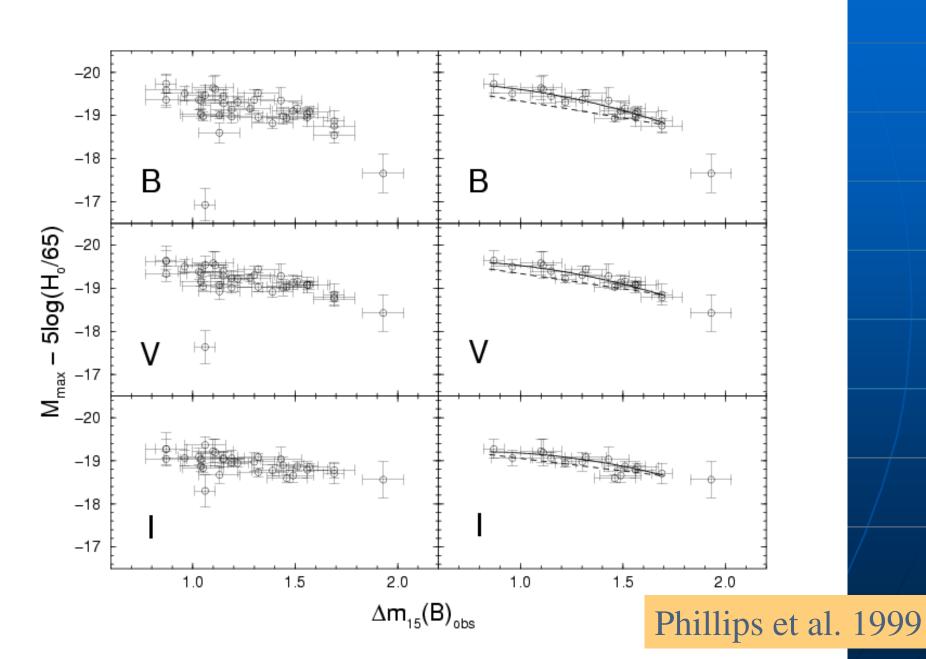


The SN Ia luminosity can be normalised:

Bright = slow
Dim = fast

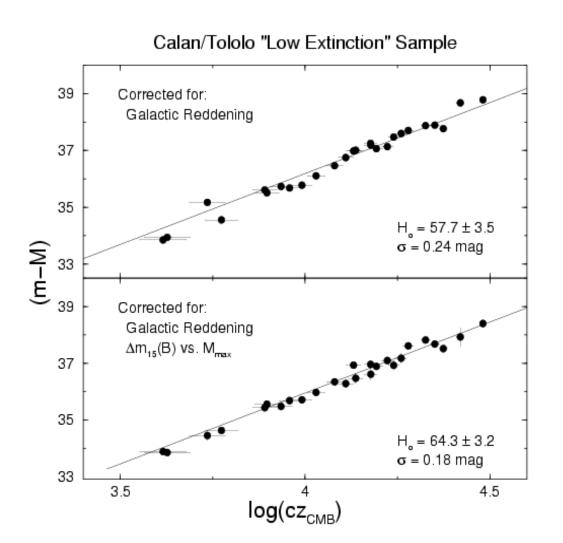
(Riess et al. 1996)

# Correlations



# Normalisation of the peak luminosity

#### Phillips et al. 1999



In Using the luminosity-decline rate relation one can normalise the peak luminosity of SNe Ia

Reduces the scatter!

#### n Luminosity vs. decline rate

• Phillips 1993, Hamuy et al. 1996, Riess et al. 1996, 1998, Perlmutter et al. 1997, Goldhaber et al. 2001

#### n Luminosity vs. rise time

• Riess et al. 1999

#### n Luminosity vs. color at maximum

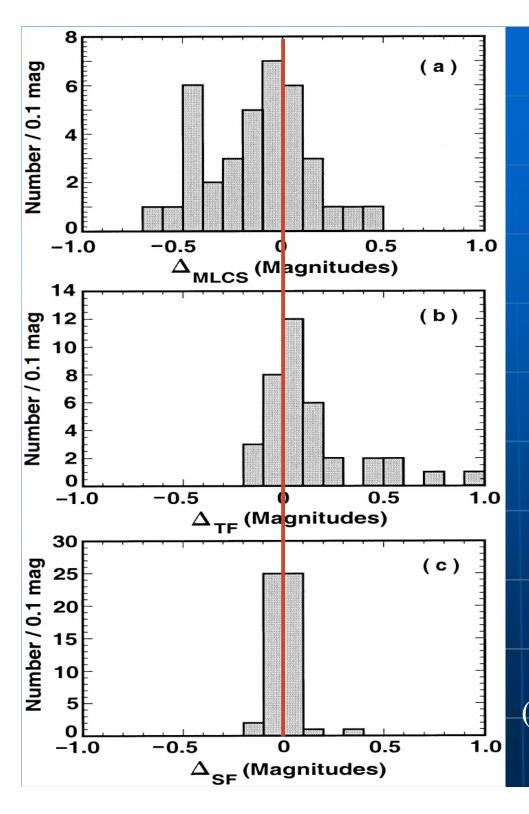
• Riess et al. 1996, Tripp 1998, Phillips et al. 1999

#### n Luminosity vs. line strengths and line widths

• Nugent et al. 1995, Riess et al. 1998, Mazzali et al. 1998

#### n Luminosity vs. host galaxy morphology

• Filippenko 1989, Hamuy et al. 1995, 1996, Schmidt et al. 1998, Branch et al. 1996



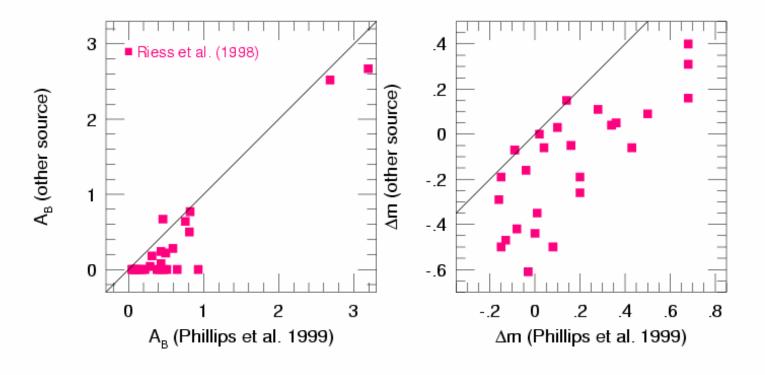
Riess et al. 1998

Phillips et al. 1999

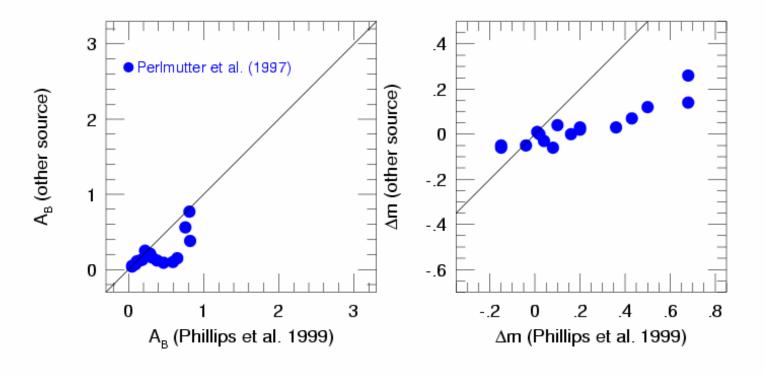
Perlmutter et al. 1997

(Drell et al. 2000)

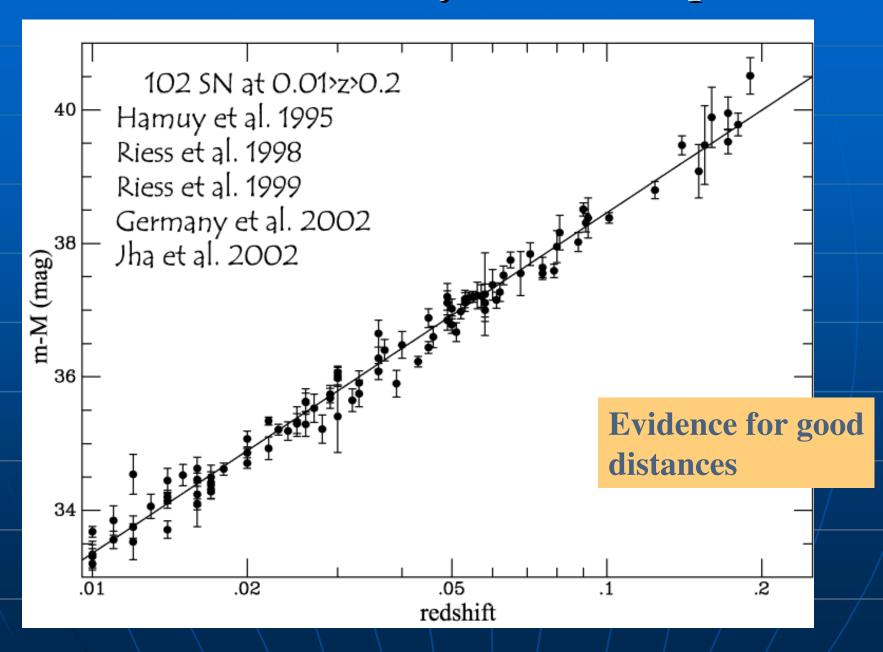








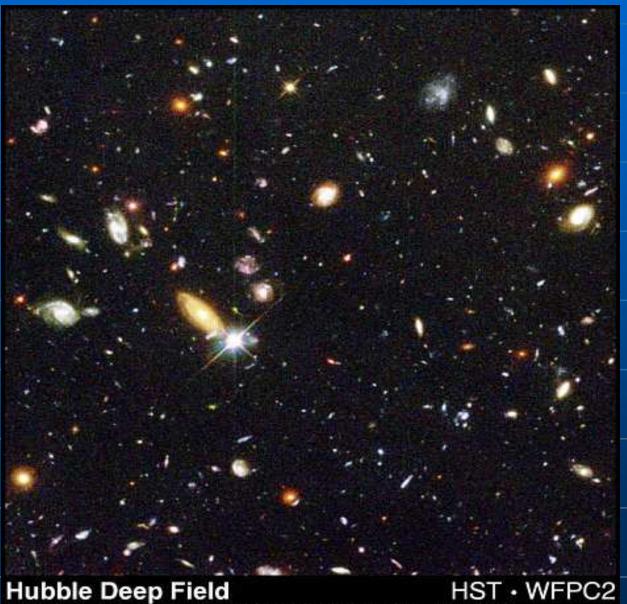
### The nearby SN Ia sample



### Hubble constant from SNe Ia

- n Extremely good (relative) distance indicators
  - distance accuracy around 10%
- <sup>n</sup> Uncertainty in H<sub>0</sub> mostly from the LMC and the Cepheid P-L relation

# Very distant supernovae



Supernovae are very rare, ~ 1 SN per 100 years and galaxy.

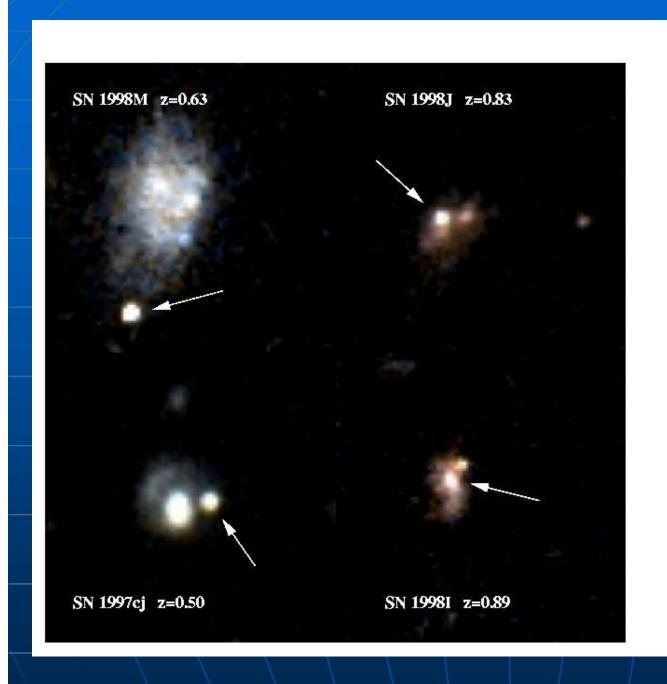
One has to observe very many galaxies!

PRC96-01a · ST Scl OPO · January 15, 1996 · R. Williams (ST Scl), NASA

# Supernova 1998ba Supernova Cosmology Project (Perimutter, et al., 1998) (as seen from Hubble Space Telescope) 3 Weeks Before as seen from Difference

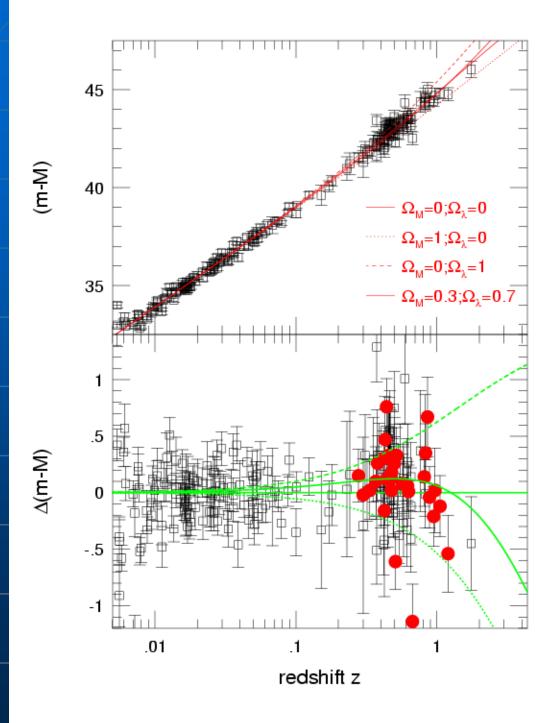
#### Search strategy:

- 1. Repeated scanning of a certain field.
- 2. Electronic readout of the data.
- 3. Follow-up observations, e.g., HST, VLT,



# Supernovae are routinely detected at redshifts Z > 0.4:

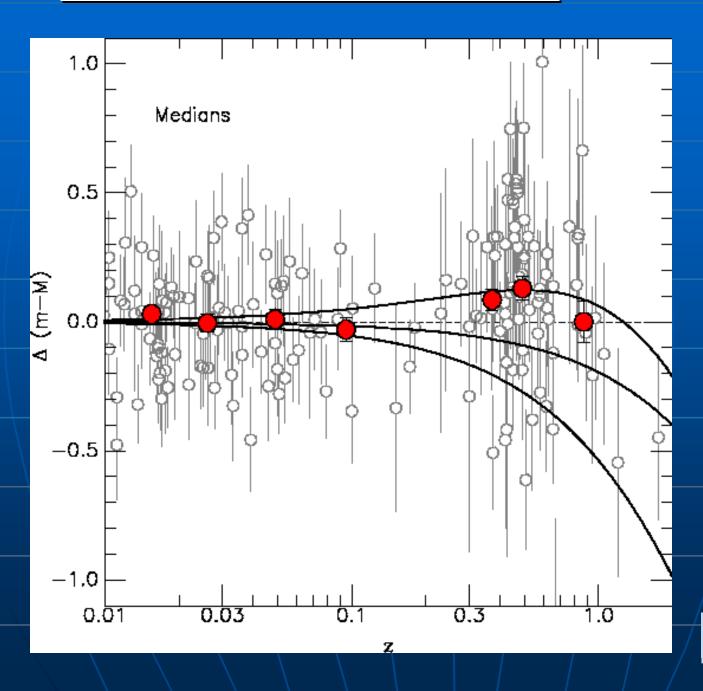
- What is the intrinsic scatter in luminosities?
- Are they different from the local sample?
- Do we understand the differences?



# Supernovae at high redshifts

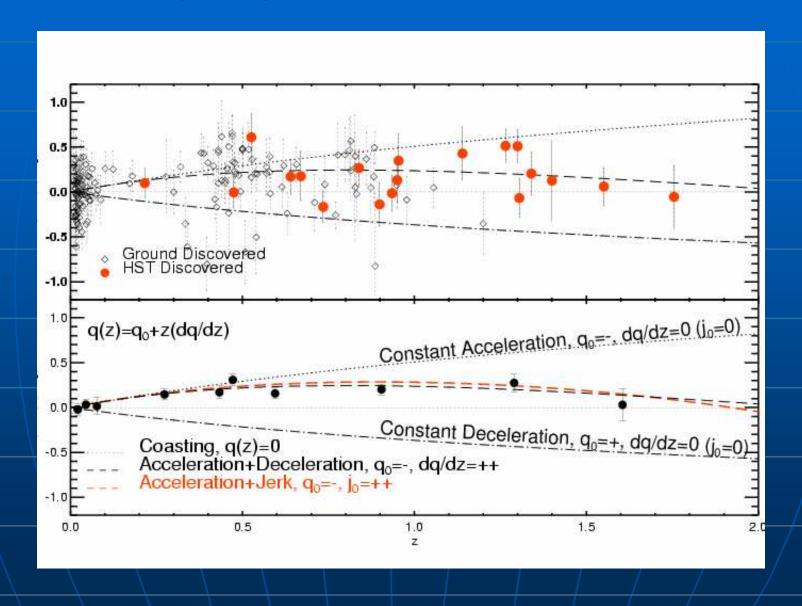
Tonry et al. 2003

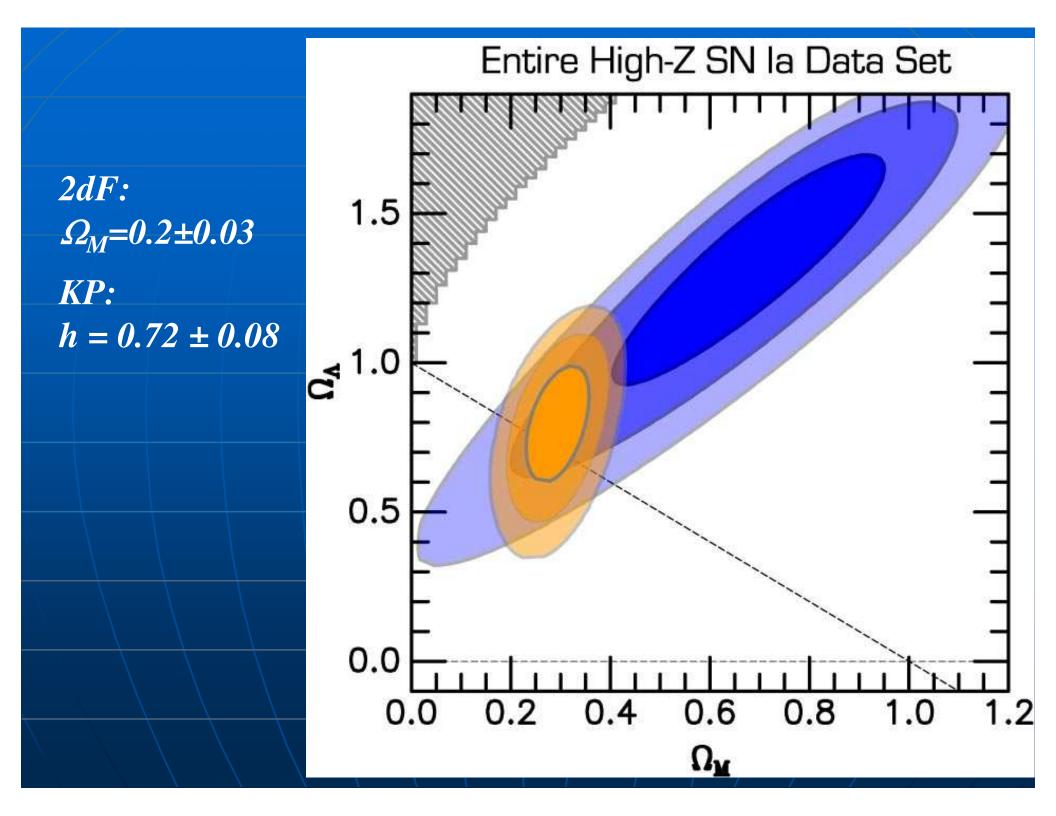
# 209 SNe Ia and medians



Tonry et al. 2003

# Very high redshift SNe Ia





#### n General luminosity distance

$$D_{L} = \frac{(1+z)c}{H_{0}\sqrt{|\Omega_{\kappa}|}} S \left\{ \sqrt{|\Omega_{\kappa}|} \int_{0}^{z} \left[ \Omega_{\kappa} (1+z')^{2} + \sum_{i} \Omega_{i} (1+z')^{3(1+\omega_{i})} \right]^{-1/2} dz' \right\}$$

• with 
$$\Omega_k = 1 - \sum_i \Omega_i$$
 and  $\omega_i = \frac{p_i}{\rho_i c^2}$ 

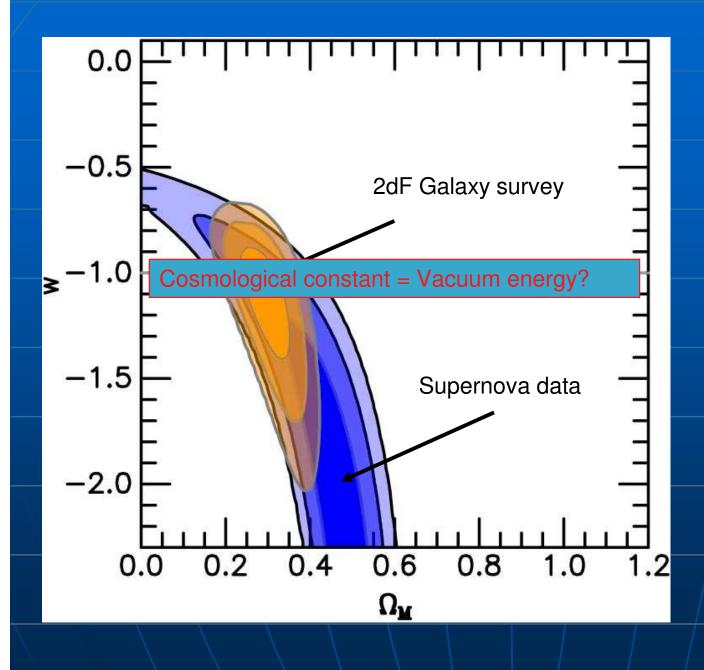
$$\omega_i = \frac{p_i}{\rho_i c^2}$$

$$\omega_{\rm M} = 0$$
 (matter)

$$\omega_{\rm R} = \frac{1}{3}$$
 (radiation)

$$\omega_{\Lambda}$$
= -1 (cosmological constant)

 $(w \equiv \omega !)$ 



# Cosmology and Typ Ia supernovae

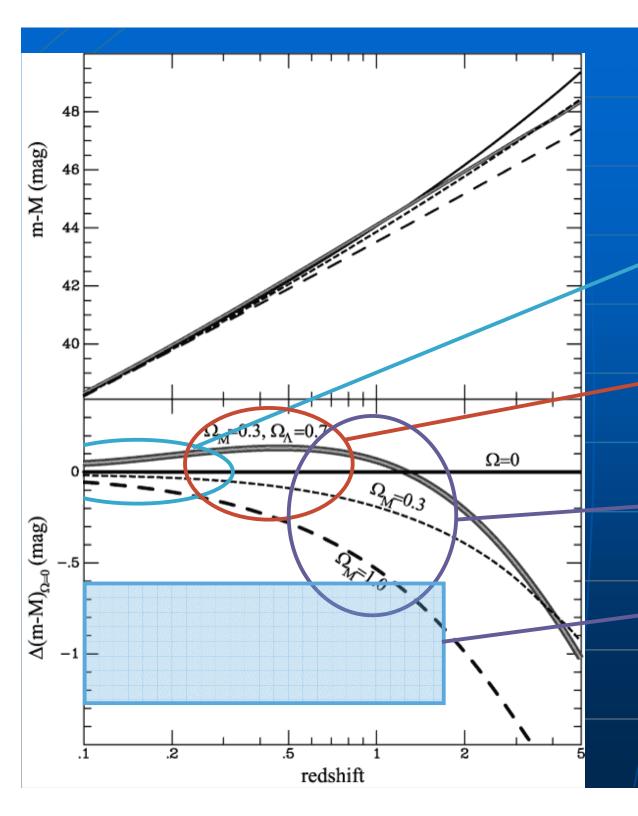
The "equation of state" of the Universe:

$$p = w\rho$$

$$\ddot{a} \sim (\rho + 3p)$$

$$w \leftarrow 1/3$$
:

Acceleration!



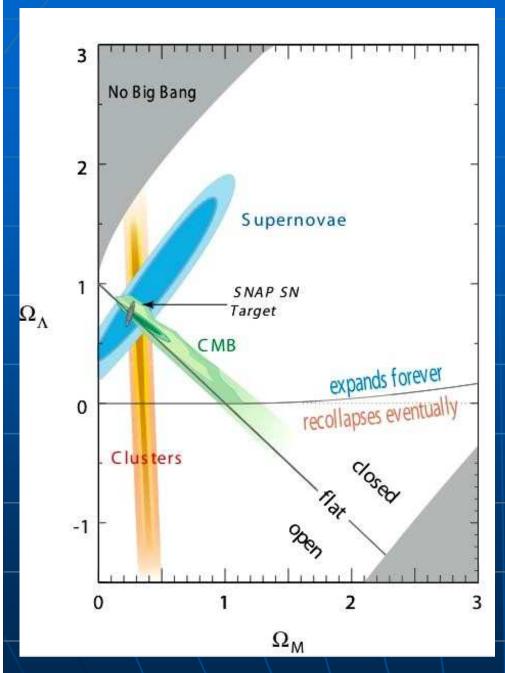
#### **SN Projects**

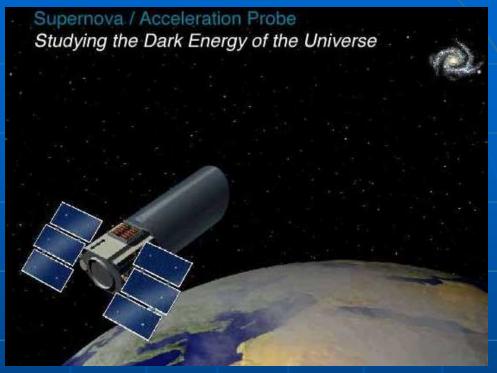
SN Factory Carnegie SN Projekt

ESSENCE CFHT Legacy Survey

High-z SN Search (GOODS)

SNAP (Supernova acceleration Probe)



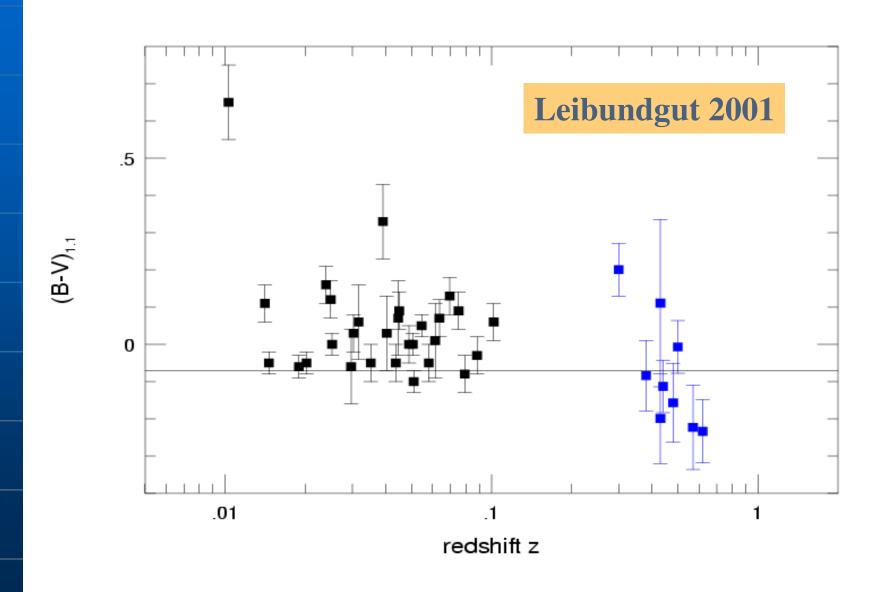


SNAP: "Supernova/Acceleration Probe"

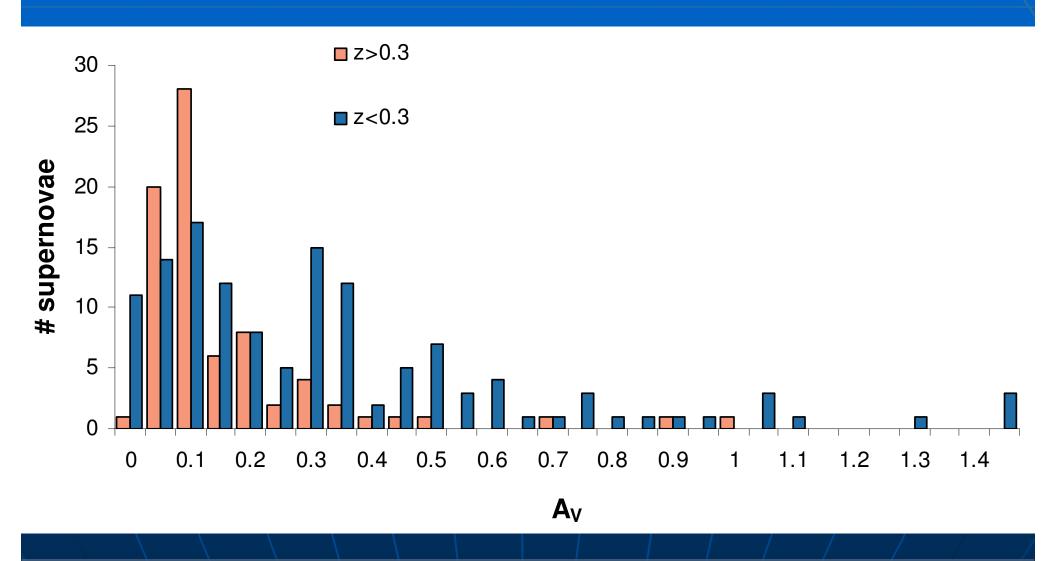
# What can still be wrong???

- q Systematic errors?
- q Pollution of high-Z samples?

# Is evolution a problem?

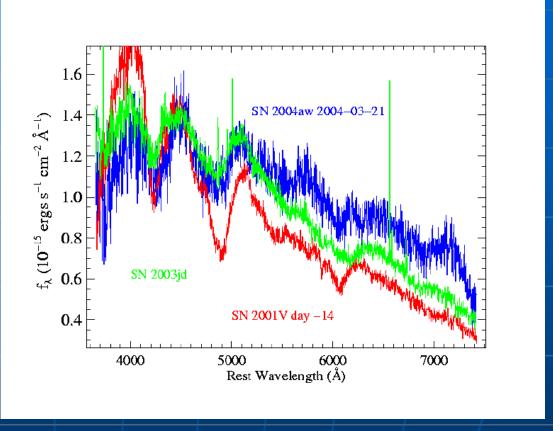


# Absorption distributions



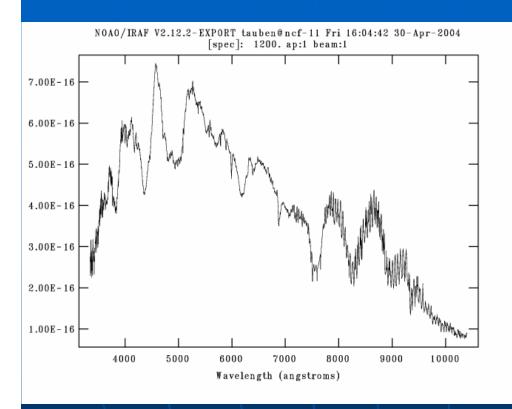
## Pollution of the samples: SN 2004aw?

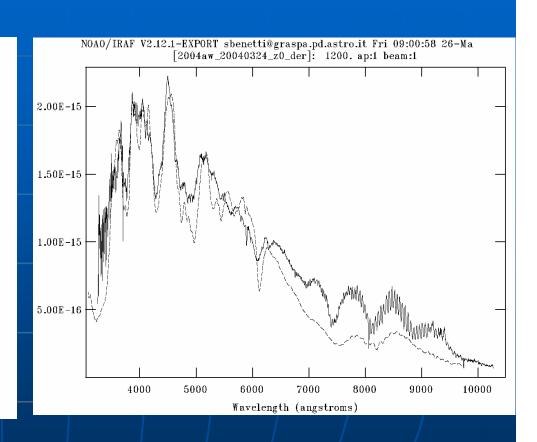




- Discovery: *March 19, 2004*
- Host galaxy: *NGC 3997*; V = 4771 km/s
- B-Maximum: ~ *April 5, 2004*

#### First classification





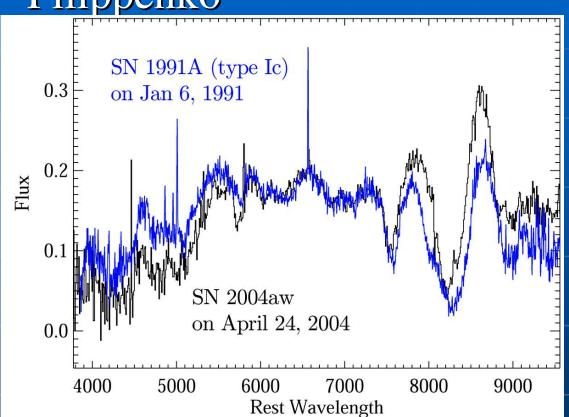
Early spectrum of SN 2004aw

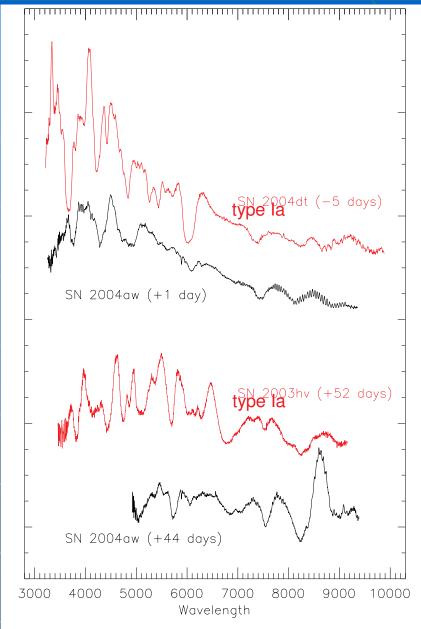
... and SN 1991T

(S. Benetti)

#### n After one month of observations:

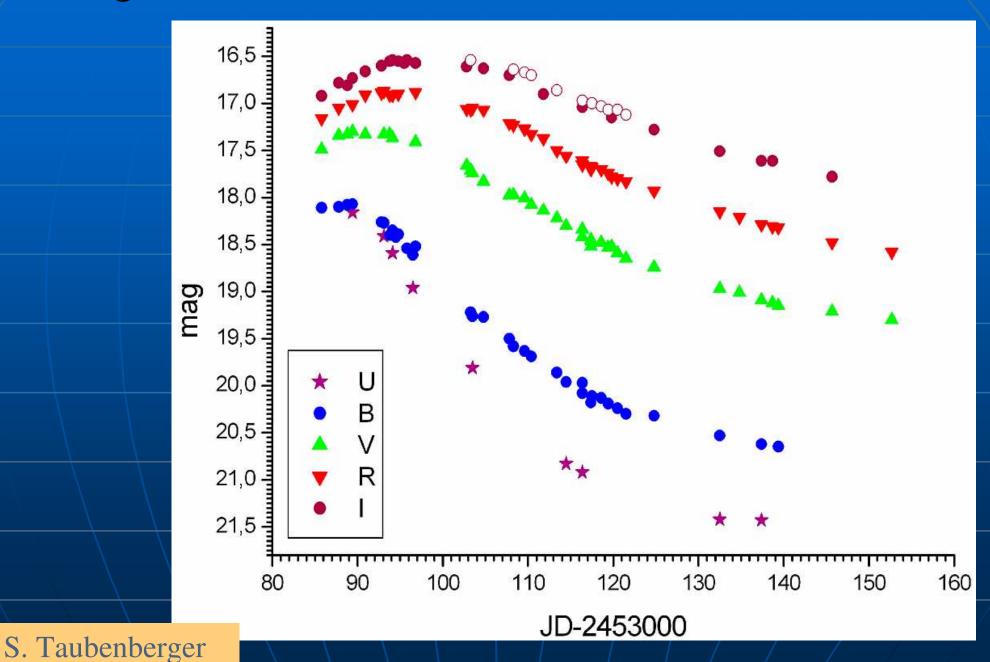
Reclassification by AlexFilippenko





It's a Ic!!!

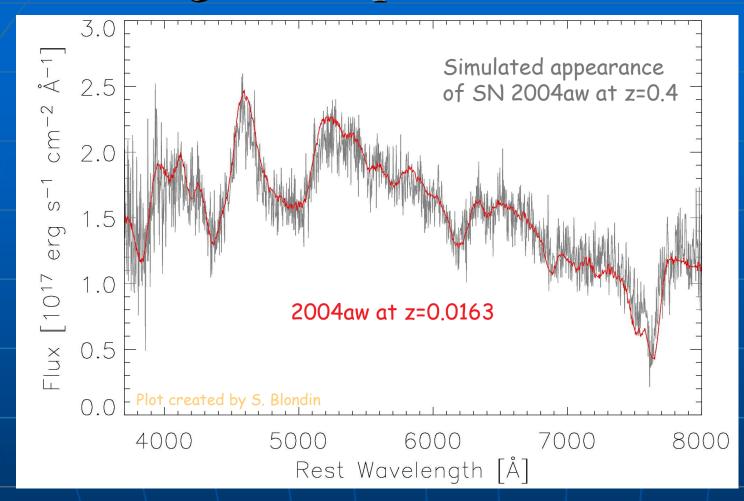
#### Light curves of SN 2004aw



## Absolute magnitudes

- <sup>n</sup> For 2004aw extinction and distance highly uncertain:
  - A(B) = 1.80 mag and A(V) = 1.36 mag from EW measurements
  - $\mu = 34.23$  from host galaxy recession velocity
- n Results for maximum brightness:
- B = -17.95 + -0.47 mag, V = -18.26 + -0.39 mag, ...(Errors almost arbitrary!)
- n For comparison:
  - peak luminosities of normal type Ia SNe B = -19.4 mag and V = -19.4 mag with scatter (+/- 0.3 mag)
  - For type Ic SNe:  $B = -16.2 \dots -19.4 \text{ mag}$  and  $V = -16.7 \dots -19.8 \text{ mag}$

## High-z sample contamination?



+1.5 days spectrum

(S. Taubenberger)

n Result of a classification code of S. Blondin (for the z = 0.4 case):

best match: 1991T (Ia pec) @ +17.4d

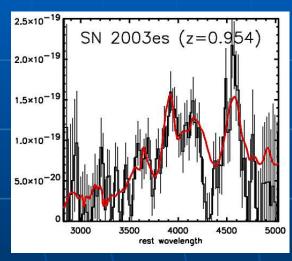
n

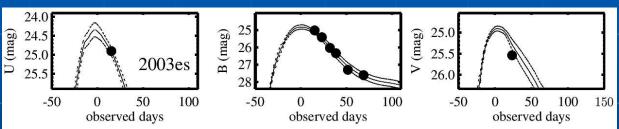
second best: 1992A (Ia) @ +9.0d

third best: 1995D (Ia) @ +8.1d

#### What does this mean?

- Assume a hypothetic 2004aw at z = 0.4 and UNREDDENED (otherwise too faint)
- Typical dataset obtained for such SNe:
   one spectrum + sparse photometry in 2 or 3 filters





Figures from Ries et al. 2004

- 2004aw would be classified as Ia (pec) according to the spectrum
- B-V color determined, result: 0.28 mag (instead of -0.05 mag for a typical type Ia)

- Mis-interpreted as extinction, correction for E(B-V) = 0.33 applied
- n This brings 2004aw to absolute magnitudes of -19.3 in B and V
- Similar to type Ia SNe, aligns rather well in the Hubble diagram

But: Only by chance!

- Of course the procedure is much more complex and sophisticated in reality.
- Nevertheless: the danger of contamination remains!

# Summary (Part III)

- n Type II supernovae are good distance indicatores out to a few Mpc.
- n They measure absolute distances without any calibration!
- n <u>Type Ia</u> supernovae are very good distance indicators in the <u>local Universe</u>.
- n They allow to measure relative distances very accurately (after calibration).
- n They provide the best distance indicators for cosmological distances if systematic errors can be controlled.