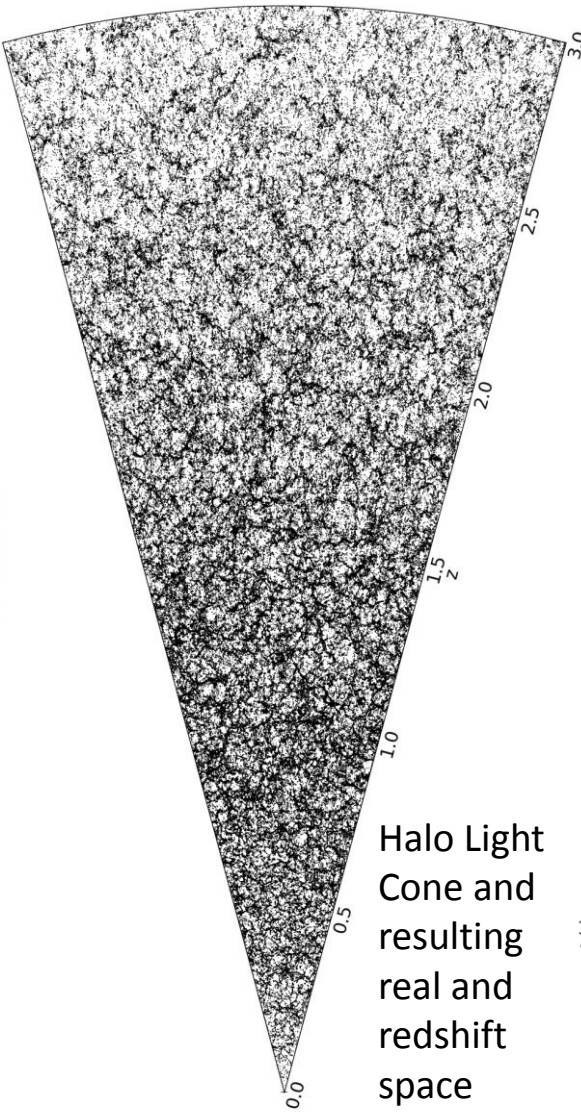


Large Scale Halo Light Cones

Shaun Cole, Alex Smith and Husni Almoubayyed
ICC, Durham



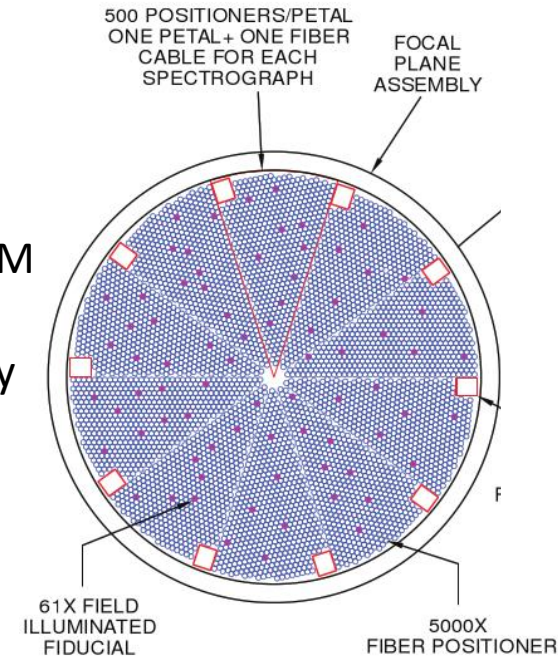
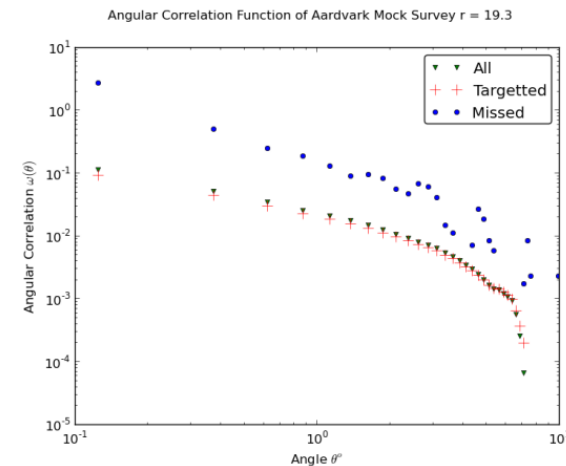
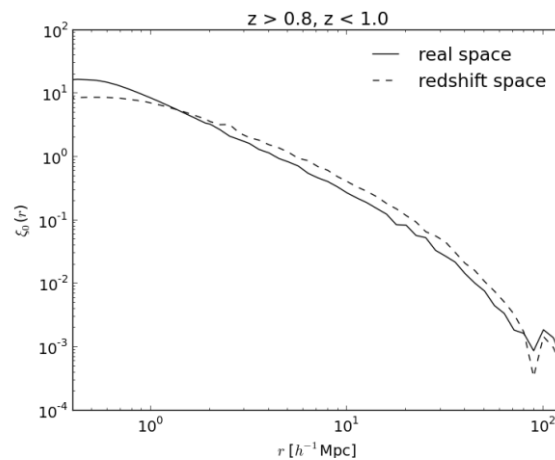
Halo Light Cone and resulting real and redshift space clustering

Subhalo merger trees from the 3 Gpc/h MXXL simulation used to construct halo lightcones with accurate clustering properties.

Can be populated with galaxies using HOD or SHAM

Example App: Design of DESI Bright Galaxy Survey (~10 million galaxies to $r \sim 19.5$)

Impact of fibre placement constraints on the clustering of targeted galaxies and development and testing of methods of mitigation.



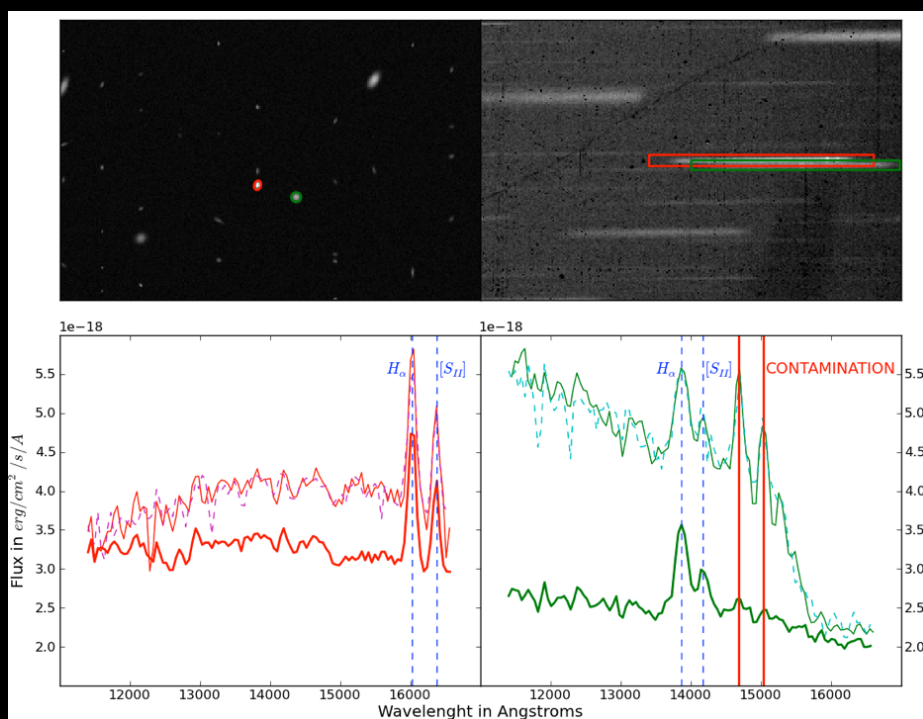
Fibres in the DESI focal plane and the effect on angular clustering



Simulating Spectroscopic Data for Euclid Galaxy Clustering

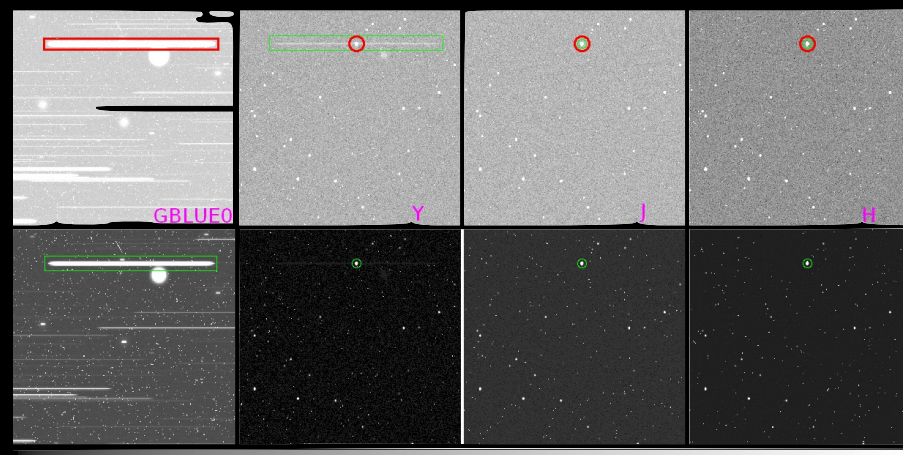
Katarina (Dida) Markovic, ICG Portsmouth, UK

- Euclid:
 - 10s of millions of spec-z
 - slitless spectroscopy

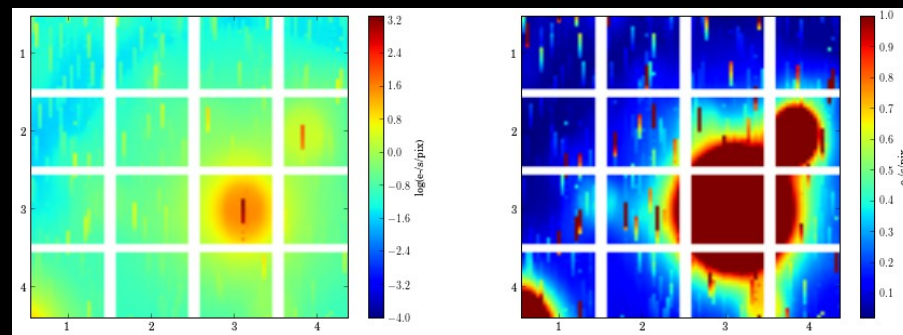


- contamination of adjacent spectra

- many pixel-level simulations on 15,000 deg^2 not feasible
- construct fast probabilistic simulations using Mangle \rightarrow see poster!
- with Sylvain de la Torre & Julien Zoubian, Marseille, FR (plots courtesy of JZ)



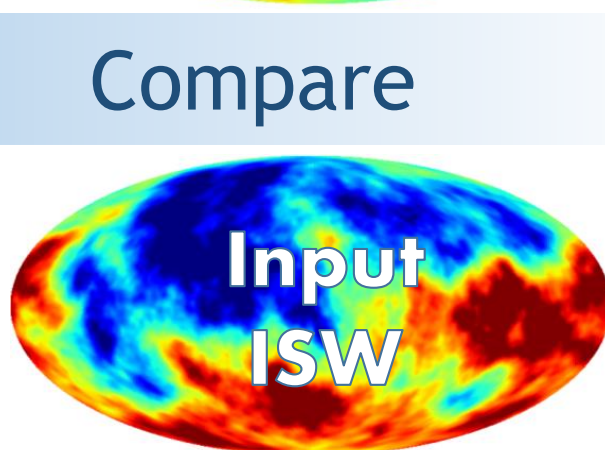
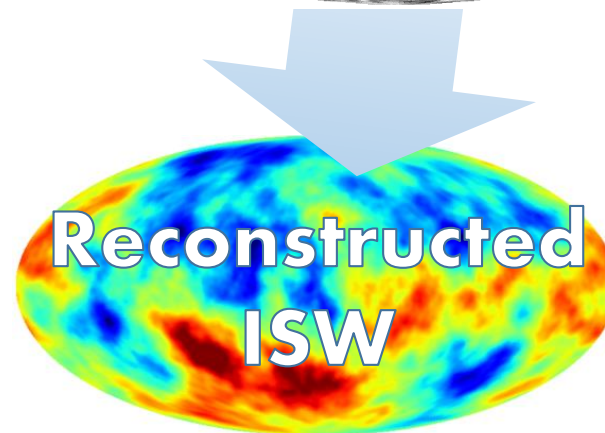
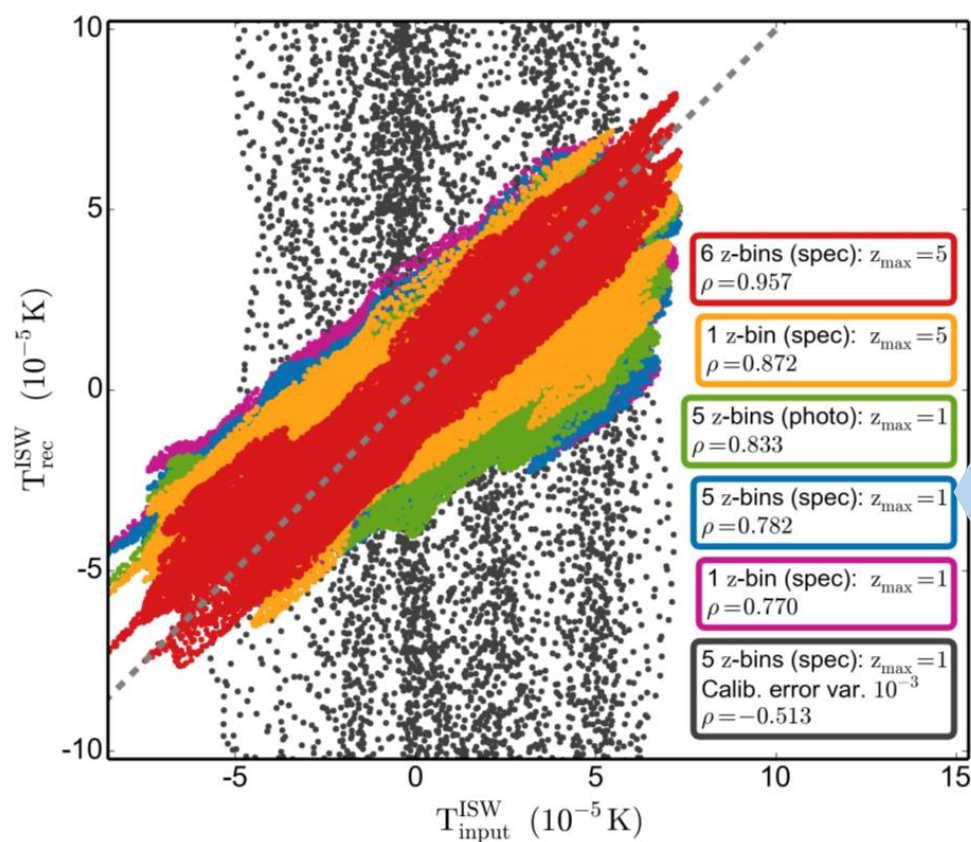
- contamination by noise effects: persistence



- contamination by noise effects: straylight

The impact of LSS survey systematics on ISW signal reconstruction

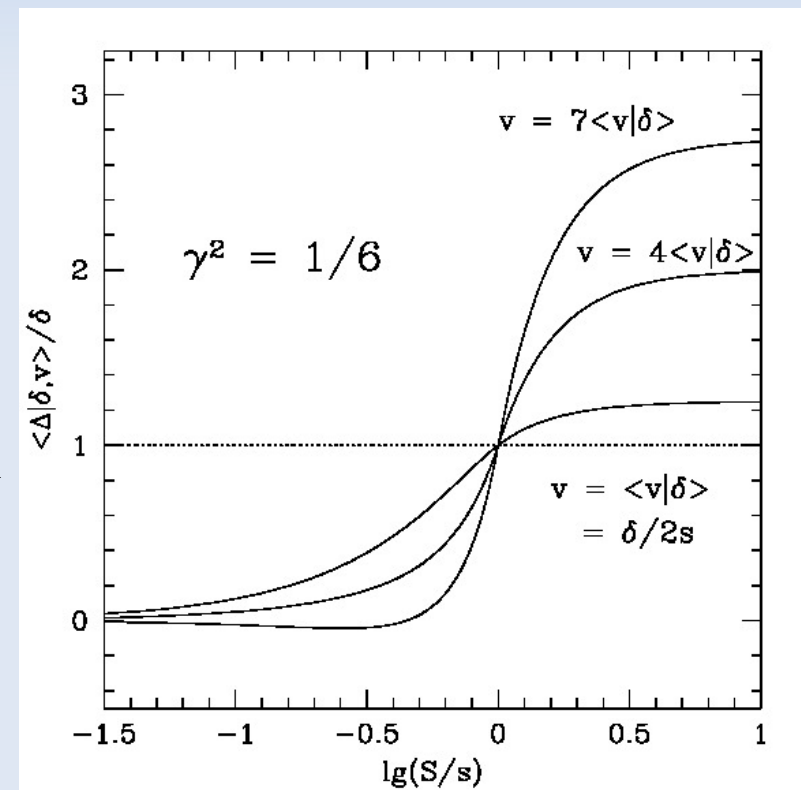
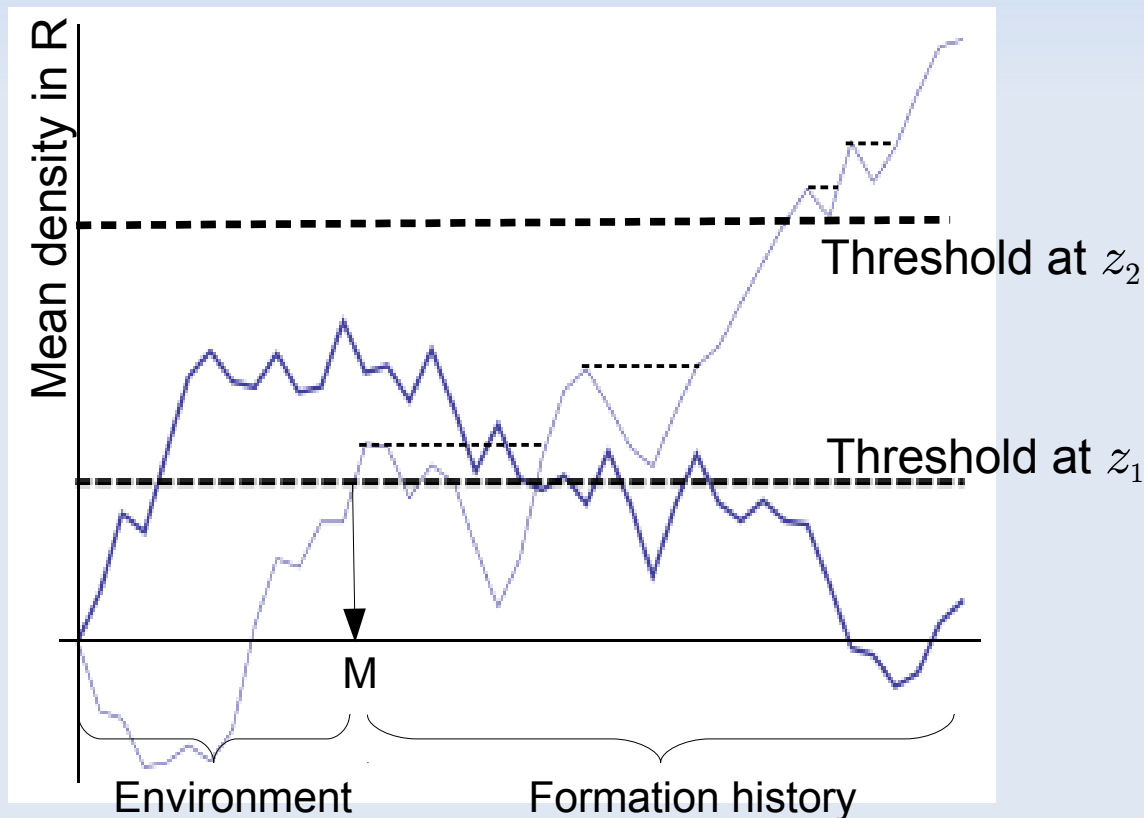
Jessica Muir and Dragan Huterer
University of Michigan



Compare

Lagrangian assembly bias

- ★ Small-scale clustering correlates with more than just large-scale mean density
- ★ Formation history correlates with environment
- ★ At same mass, **more concentrated** halos prefer **lower density** environment
- ★ Well described by just one extra variable: **steepness** of Lagrangian profile



Investigating the σ_8 tension by the cross-correlation of tSZ and cosmic shear

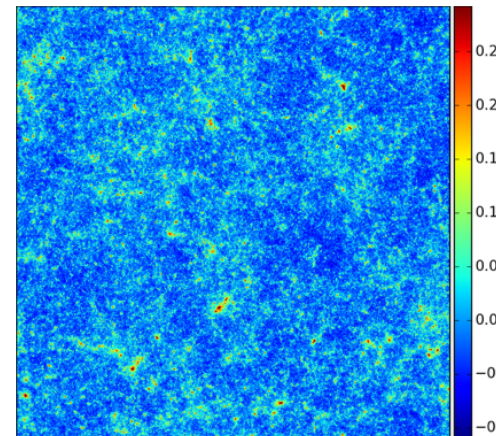
Ken Osato (Dept. of Physics, The Univ. of Tokyo)

- ✓ **Cosmic shear and thermal Sunyaev-Zel'dovich effect**
From our **baryonic N-body simulations**, mock cosmic shear and tSZ maps can be created.

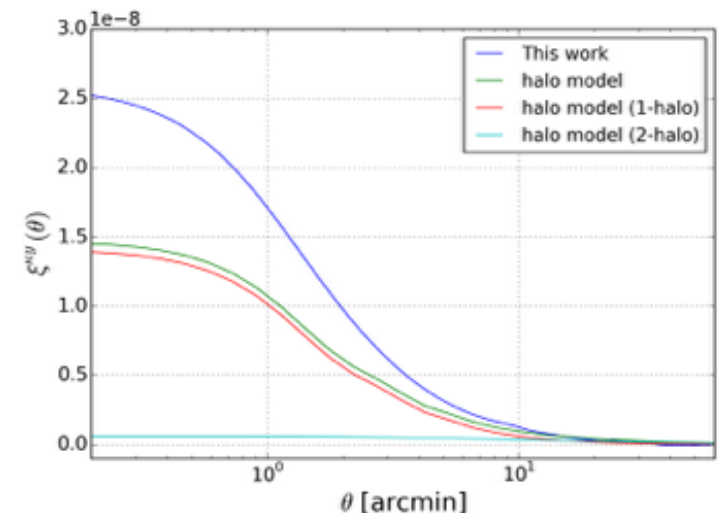
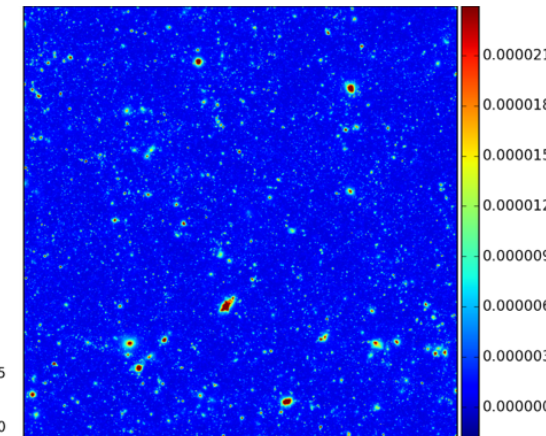
- ✓ **Cross-correlation**

Angular cross-correlation can be measured from two maps. This can be useful in evaluating cosmological parameters like σ_8 .

cosmic shear



tSZ



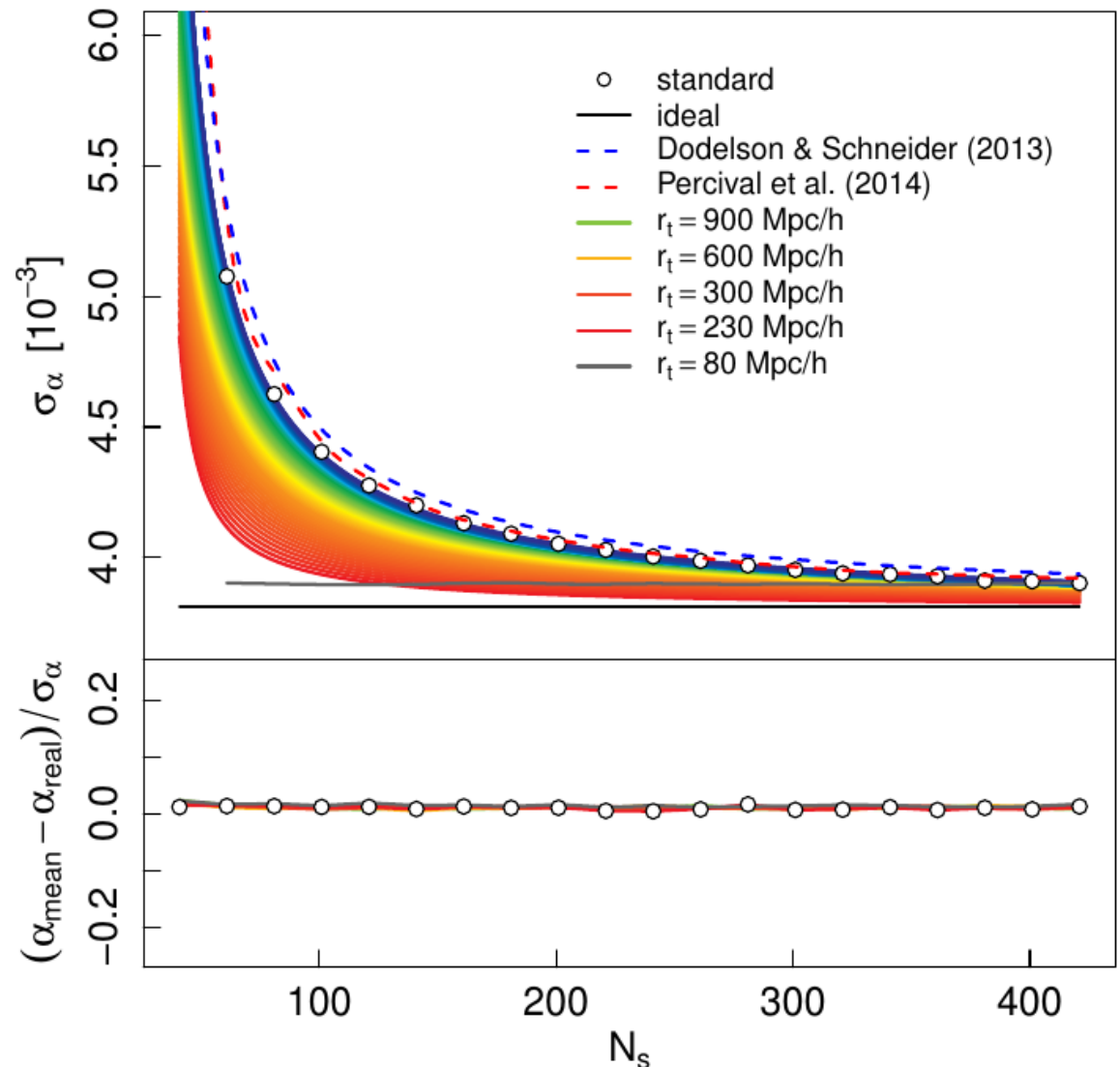
Improving the precision (matrix) in precision cosmology

By “tapering” the covariance matrix (\mathbf{C}) at large scales we get unbiased and less noisy estimation of the precision matrix (Ψ):

$$\mathbf{C}_t \equiv \hat{\mathbf{C}} \circ \mathbf{T}$$

$$\Psi^t = \mathbf{C}_t^{-1} \circ \mathbf{T}$$

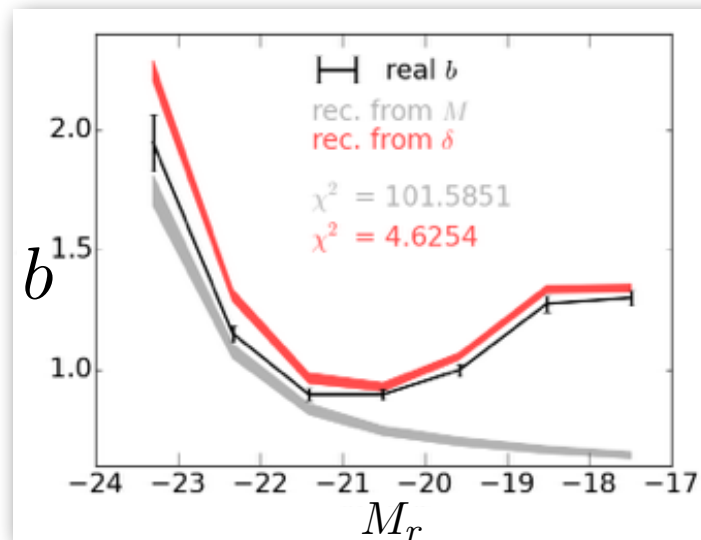
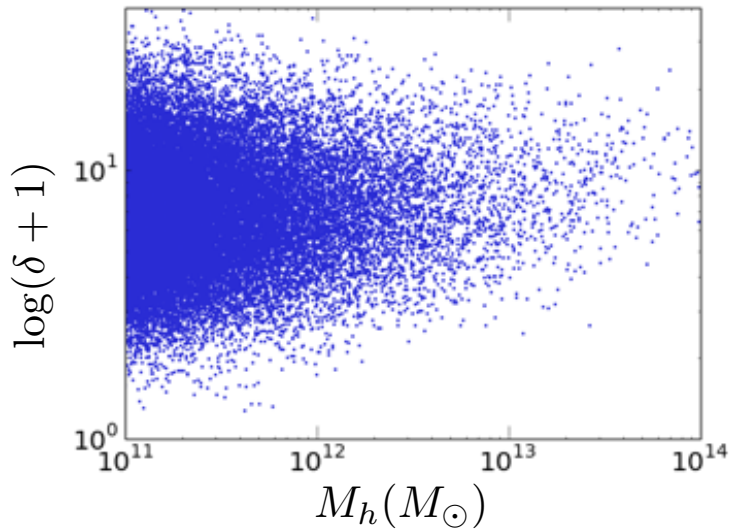
Error of the BAO peak scale (α) as function of the number of mocks used to compute Ψ . Using tapering method we obtain errors closer to the ideal expectations.



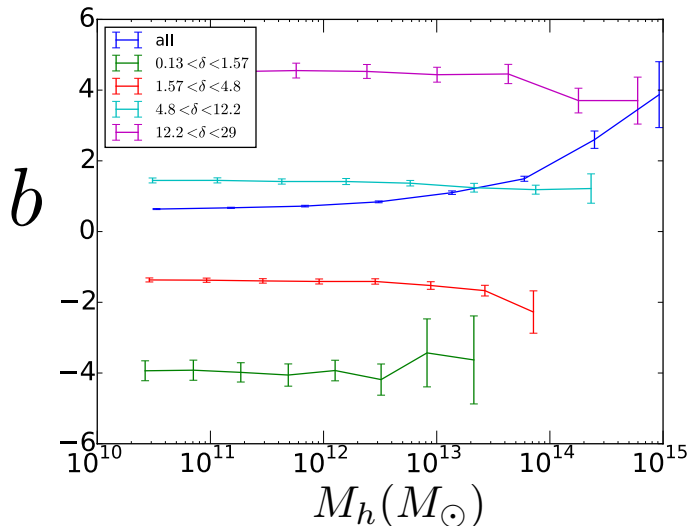
Halo Clustering: Environment vs Mass

Arnau Pujol
Enrique Gaztañaga

Institut de Ciències de l'Espai (IEEC-CSIC), Barcelona



HOD predictions can be affected by assembly bias. In this case, density (red) predicts much better clustering than mass (grey)



For fixed environmental density, Halo bias does not depend on Mass

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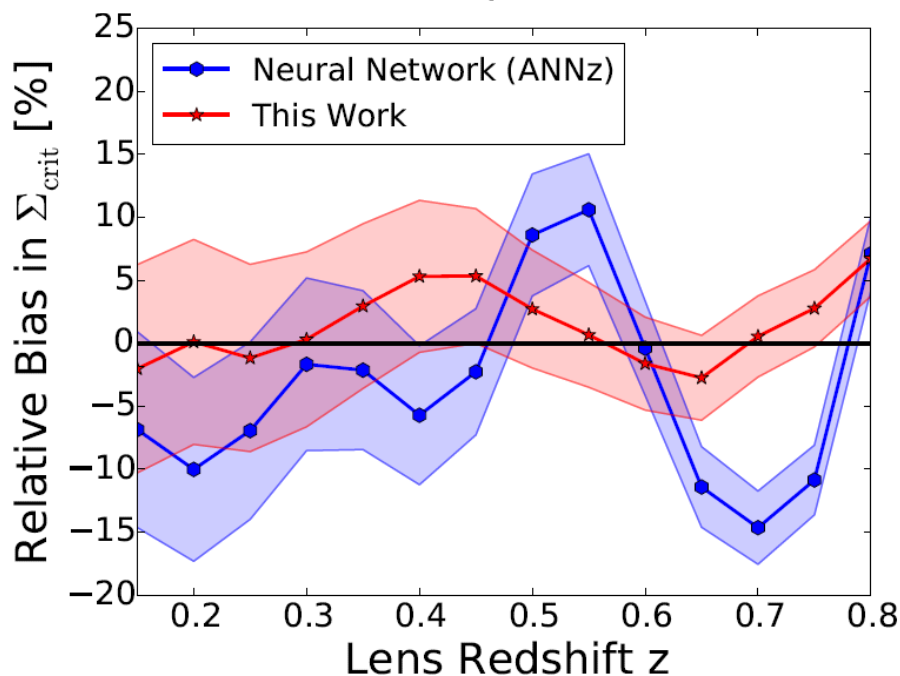
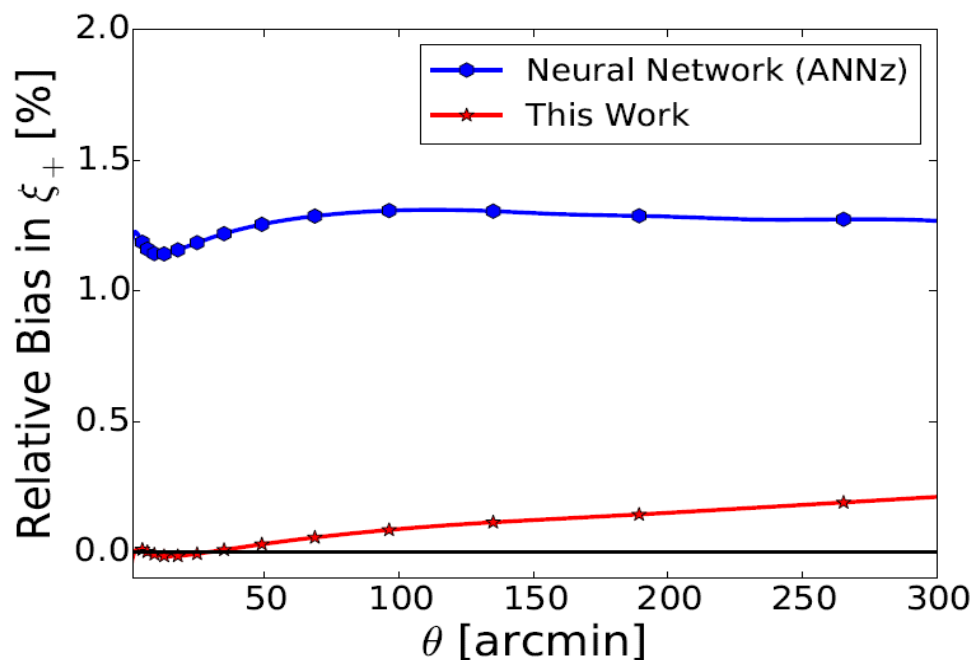
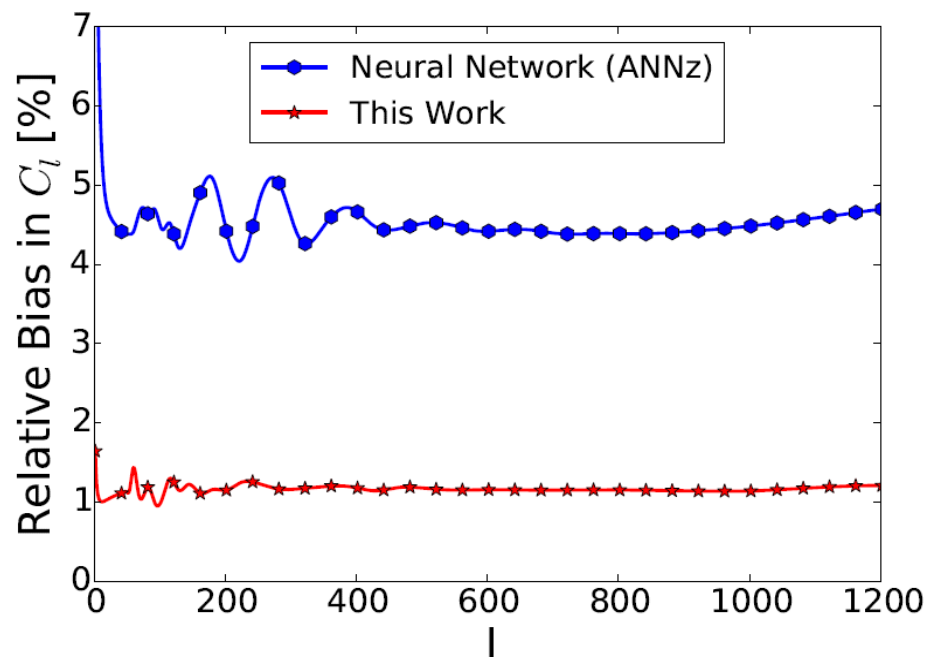
INSTITUTO DE
CIENCIAS DEL
ESPACIO

ICE

Accurate photometric redshift probability density estimation - method comparison and application

Markus Michael Rau, Stella Seitz, Fabrice Brimiouille, Eibe Frank, Oliver Friedrich, Daniel Gruen, Ben Hoyle

Accepted for publication in the MNRAS 2015 July 10



**Low PhotoZ Bias in
Weak Lensing and LSS**

**Markus Michael Rau et al.
Supv. Stella Seitz
arXiv:1503.08215**

GALAXY-GALAXY LENSING OBSERVABLES AND COSMOLOGICAL RESCALING

Malin Renneby[§], Excellence Cluster Universe, Ludwig-Maximilians-Universität (LMU)

(together with Stefan Hilbert, Excellence Cluster Universe, LMU, and Raul Angulo, Centro de Estudios de Física del Cosmos de Aragón)

Goal: Acquire an accurate, fast and flexible framework to predict large-scale structure observables in order to constrain cosmological parameters as well as galaxy formation models.

Initial step: Λ CDM-setup using the rescaling method detailed in Angulo & White (2010):

$$m_p \rightarrow m'_p = \left(\frac{\Omega'_m}{\Omega_m}\right) \left(\frac{H'}{H}\right)^2 \left(\frac{L'}{L}\right)^3 m_p$$

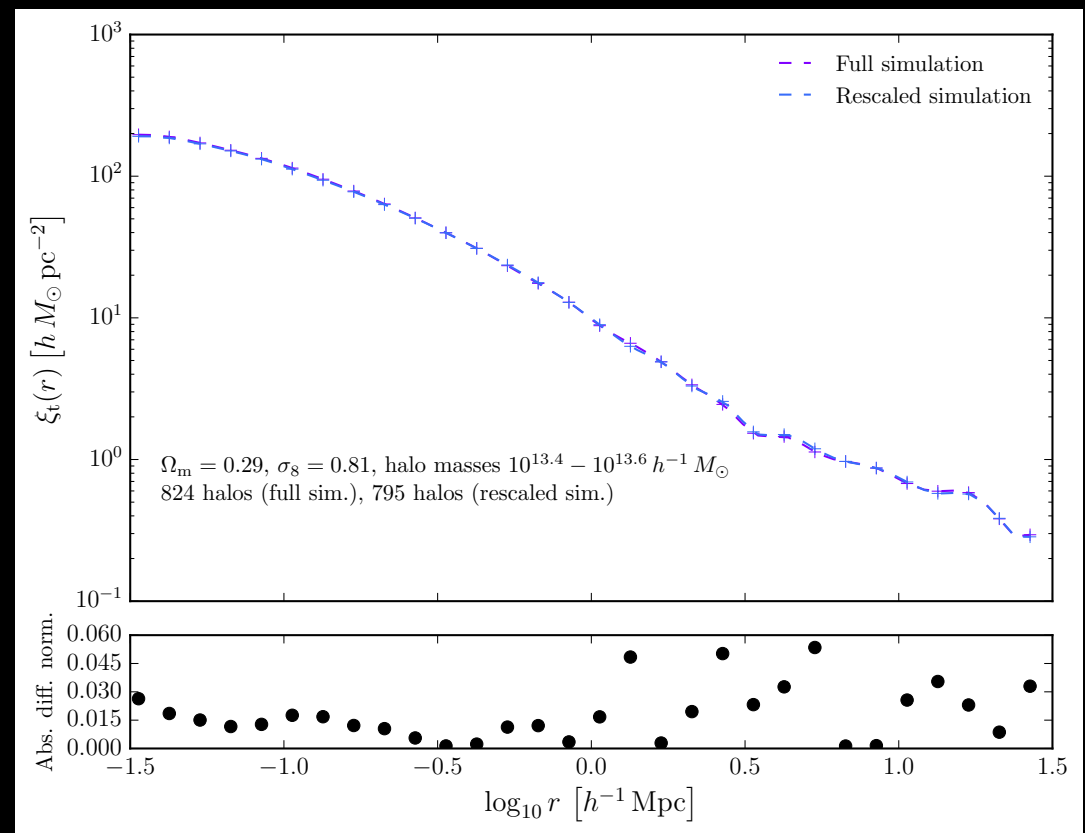
$$L \rightarrow L' = \alpha L$$

“Base cosmology”: Small Millennium-simulation: $250 (h^{-1} \text{ Mpc})^3$, $5 h^{-1} \text{ kpc}$ res., $1080^3 \approx 1.26$ billion DM particles $\rightarrow 8.861 \times 10^8 h^{-1} M_\odot$, $\Omega_m = 0.25$, $\Omega_b = 0.045$, $\Omega_\Lambda = 0.75$, $h = 0.73$, $\sigma_8 = 0.9$.

Target test cosmologies: $(\Omega_m = 0.15, \sigma_8 = 1.0)$,
 $(\Omega_m = 0.25, \sigma_8 = 0.6)$, $(\Omega_m = 0.29, \sigma_8 = 0.81)$,
 $(\Omega_m = 0.40, \sigma_8 = 0.7)$, $(\Omega_m = 0.80, \sigma_8 = 0.4)$.

Probes: Projected convergence and tangential shear profiles, halo mass function at $z = 0$.

Results: Percentage to sub-percentage deviations for the rescaled simulations with respect to the full simulations for the halo mass function except for the high mass tail, disagreements of the order of a few percent for $\Sigma(r)$ and $\xi_t(r)$ given certain halo masses and radial bin ranges (to appear in Renneby et. al. in prep.).

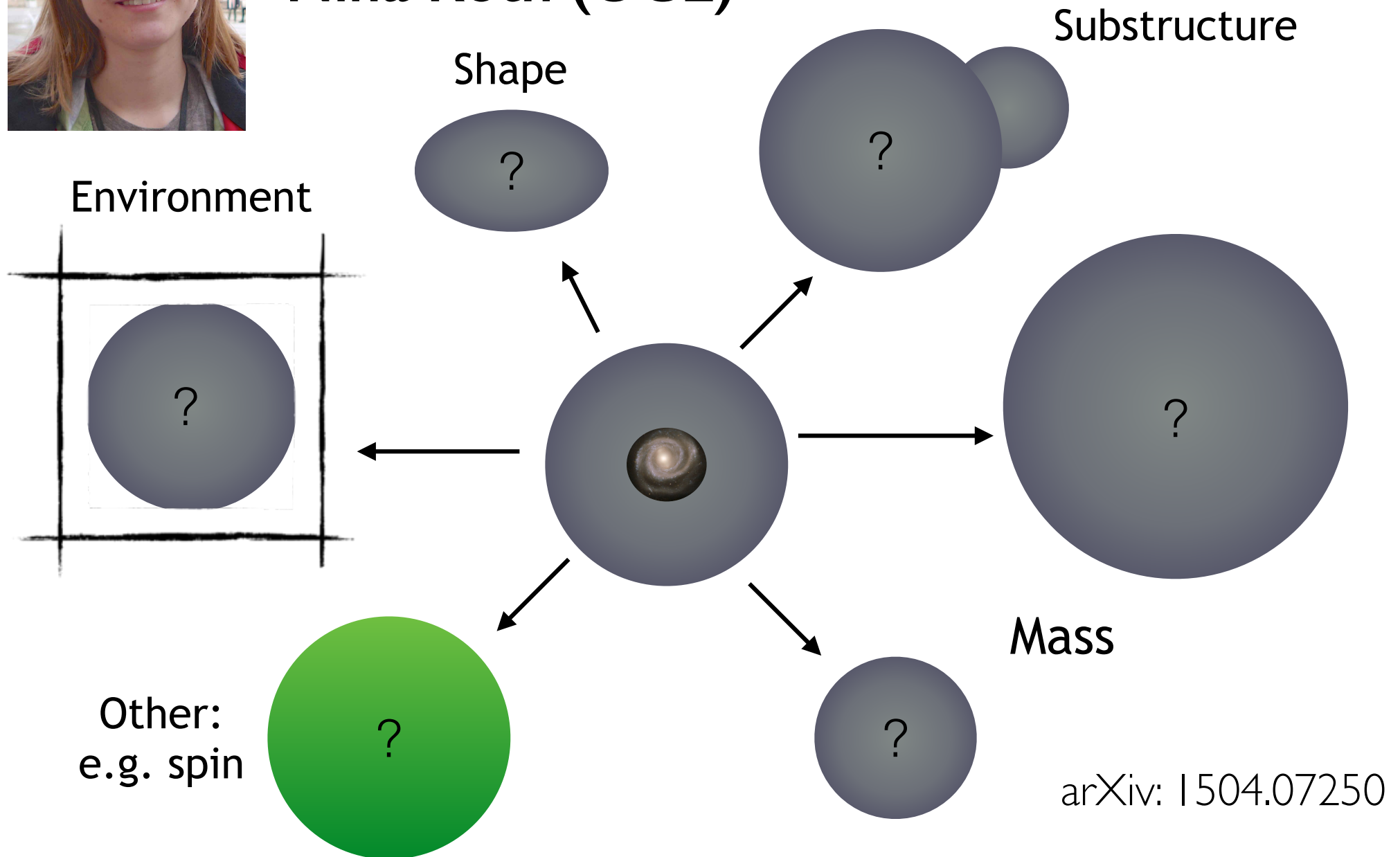


[§]malin.renneby@physik.uni-muenchen.de



Genetically modified halos

Nina Roth (UCL)

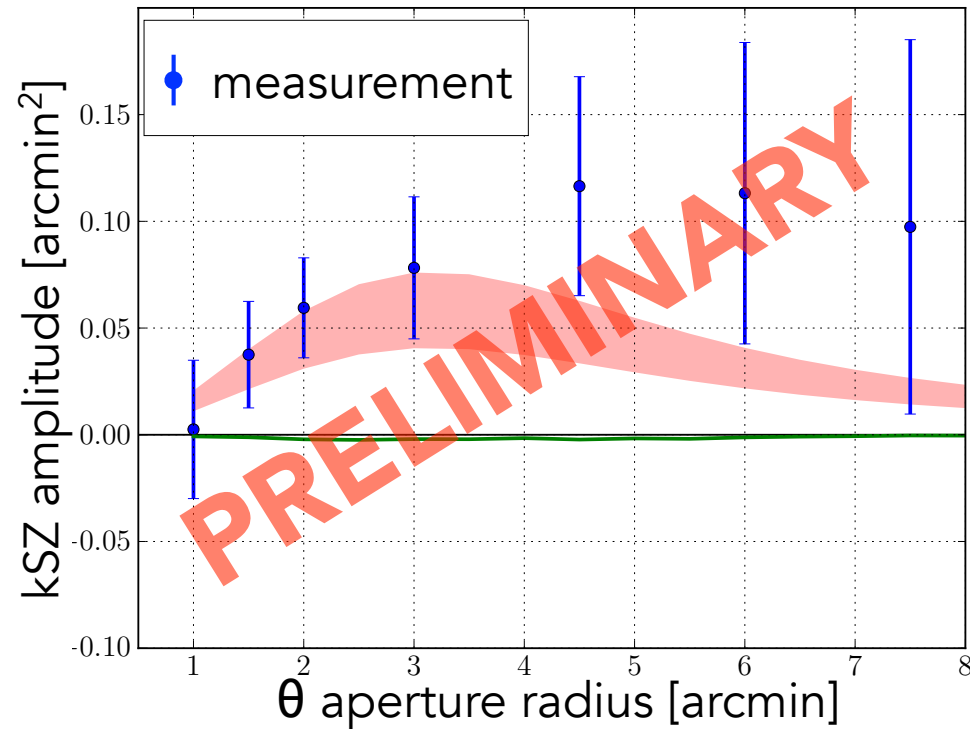
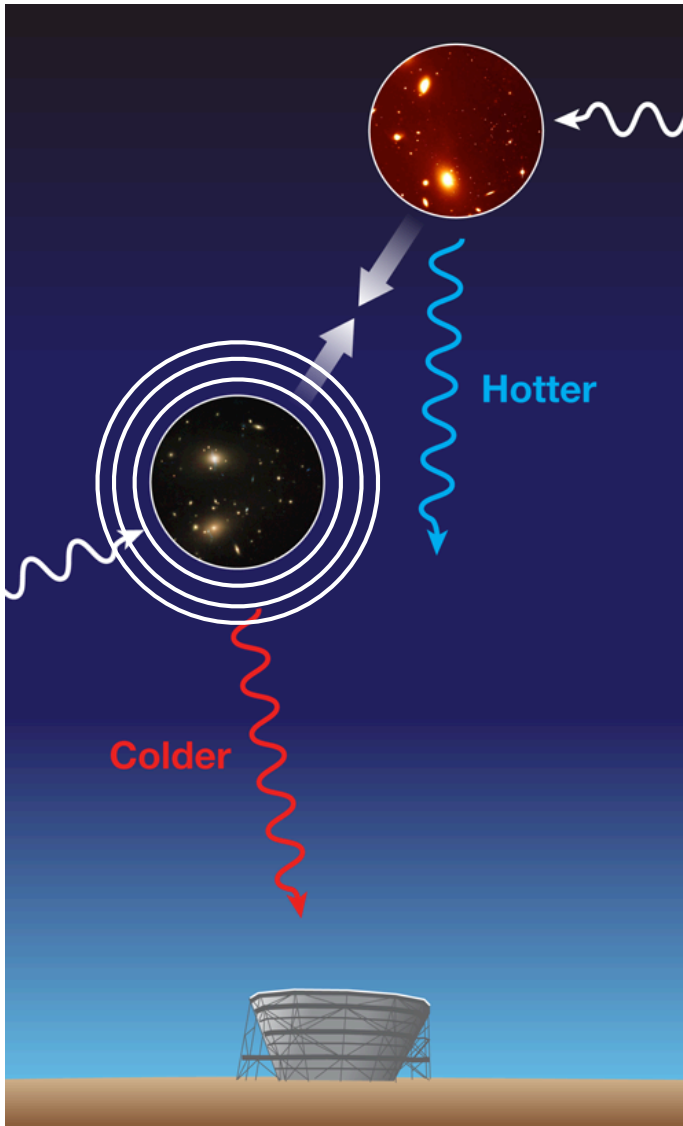


arXiv: 1504.07250



kSZ Effect & Missing Baryons

Emmanuel Schaan, S. Ferraro, M. Vargas-Magaña, K. Smith, S. Ho & ACTPol

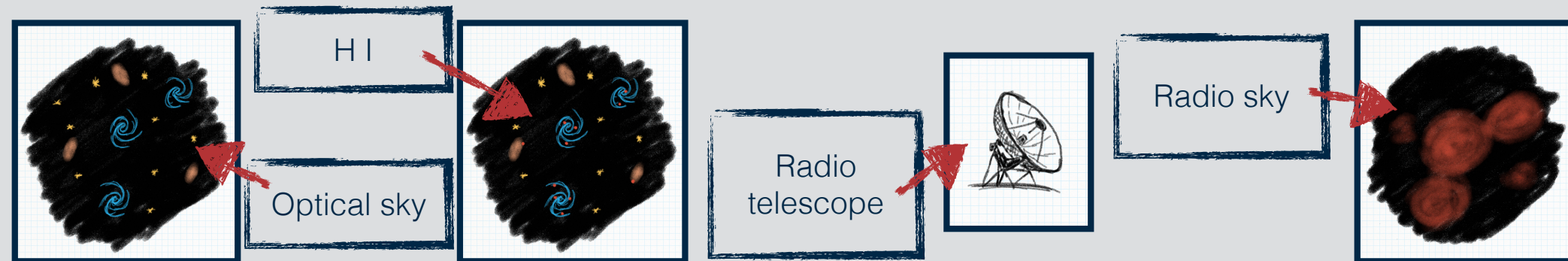


significant detection, implications for baryons and cosmology

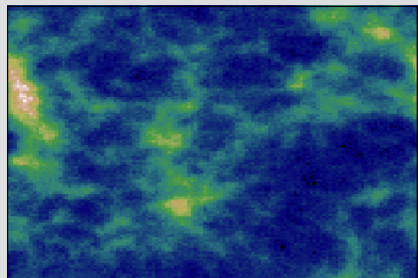
Simulations of H I Intensity Maps

S. Seehars¹, A. Paranjape^{1,2}, A. Witzemann¹, A. Refregier¹, A. Amara¹, R. Wechsler³, and M. R. Becker³

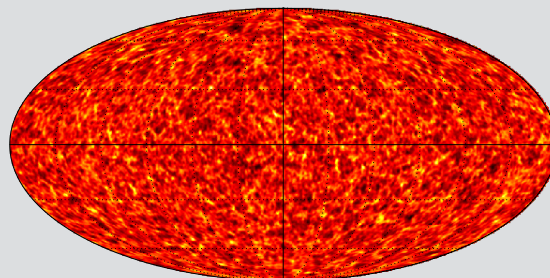
¹ETH Zurich, ²Inter-University Centre for Astronomy and Astrophysics, ³Stanford University



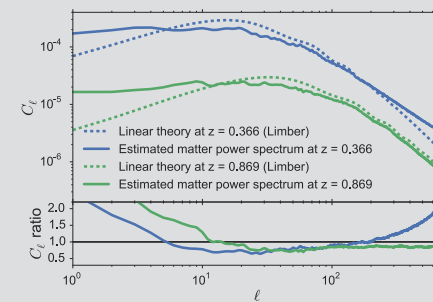
Dark Matter



Intensity Maps



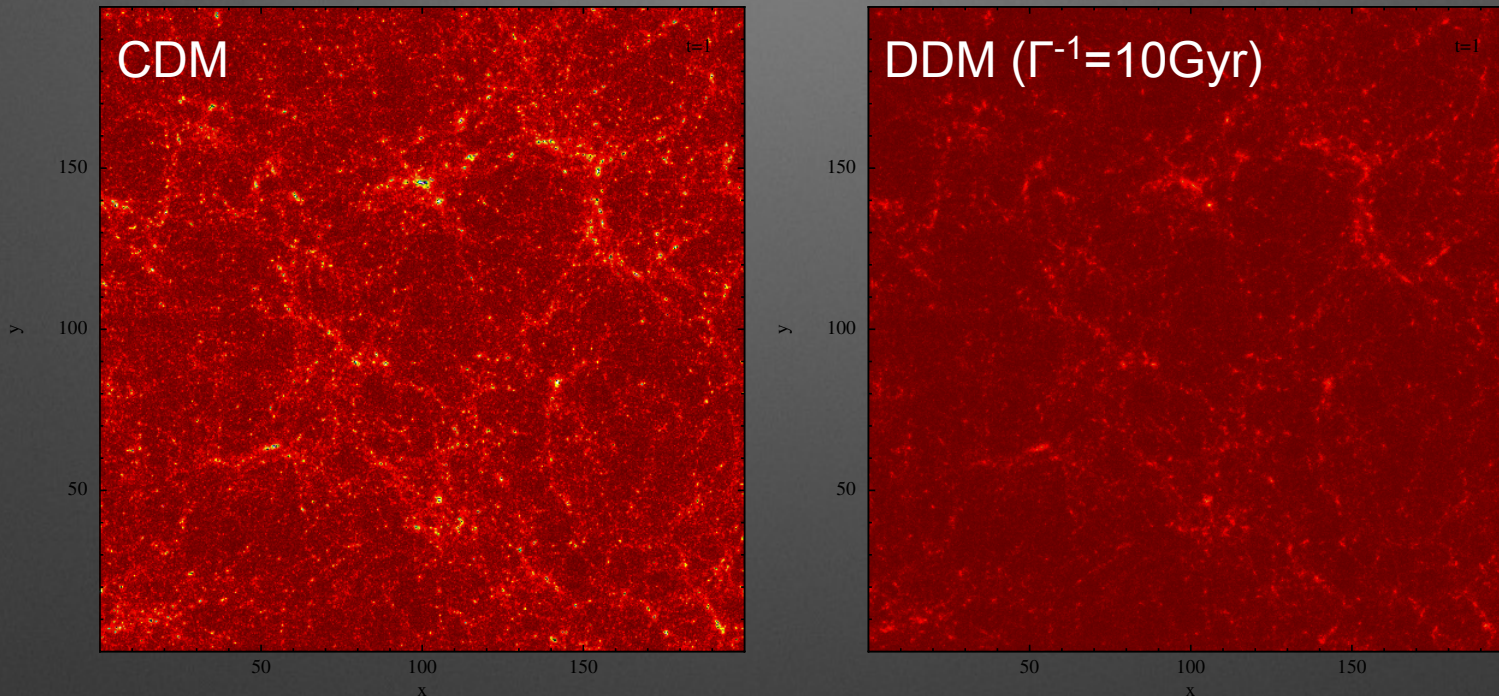
Angular Clustering



Structure formation in decaying dark matter model

Toyokazu Sekiguchi (Helsinki)

arXiv:1505.05511 with Kari Enqvist, Seshadri Nadathur & Tomo Takahashi



- ✓ N-body simulation in decaying DM model with $\Gamma^{-1} \sim O(1) t_{\text{age}}$
- ✓ Suppressed matter power spectrum & halo mass function
- ✓ Solution for tensions between Planck ΛCDM and recent observations of LSS including WL, cluster abundance

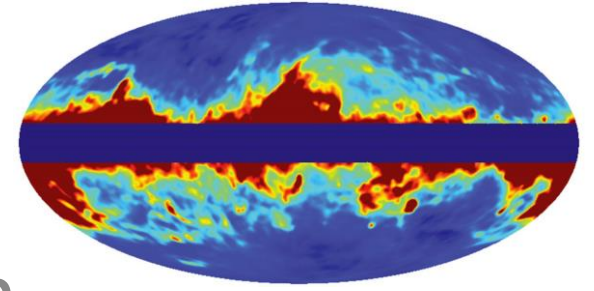
Multiplicative Errors in the Galaxy Power Spectrum

MNRAS 447, 2961 (2015)

arXiv:1410.0035

Daniel Shafer

University of Michigan

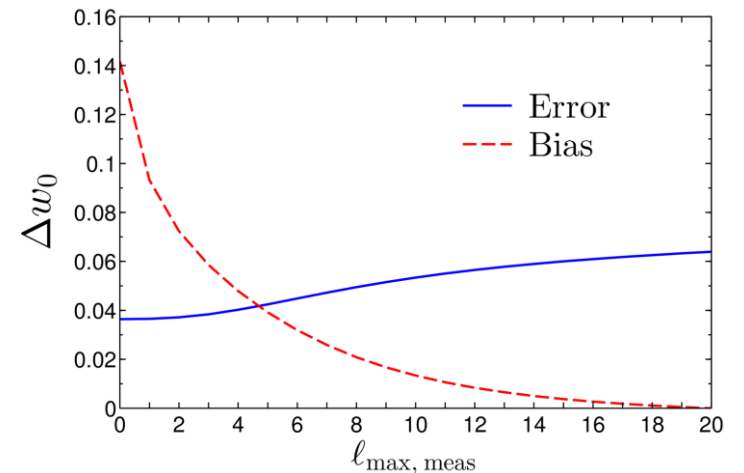


(In collaboration with
Dragan Huterer)

- Large-angle photometric calibration error must be removed from the galaxy power spectrum in order to probe the underlying cosmology.
- There is also a multiplicative effect, where large-angle error propagates to small scales:

$$T_\ell \simeq C_\ell + C_\ell^{\text{cal}} - \frac{1}{2\pi} C_\ell^{\text{cal}} C_\ell + \sigma_c^2 C_\ell$$

- We show that multiplicative errors are potentially dangerous for modern surveys, but they can be mitigated by using the large-angle power to measure the contamination directly.



Observational Evidence of the Bridge Effect of Void Filaments

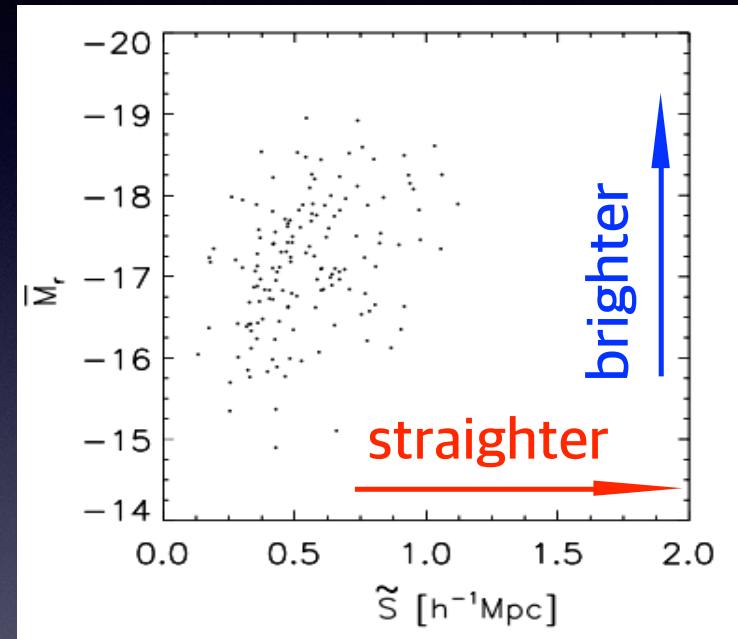
Junsup Shim, Jounghun Lee and Fiona Hoyle

The *bridge effect of void filaments* represents that **straighter void filament**, as a bridge, transports more efficiently the matter & gas from outer region into void galaxies. (Park & Lee 2009)

In this study, we aim to observationally confirm the *bridge effect of void filaments* using the void catalogue constructed by Pan et al. (2012) from SDSS DR 7.

We found that, on average, the void galaxies are **brighter** in the **straighter host void filament**.

It is also found that **richer void filaments** develop **stronger correlation** between the luminosity and straightness of the void filament.



(submitted to ApJ)

	$4 \leq N_{node} < 7$	$7 \leq N_{node} < 11$	$N_{node} \geq 11$
r	0.29 ± 0.10	0.50 ± 0.13	0.60 ± 0.13
\bar{z}	0.012 ± 0.005	0.014 ± 0.003	0.013 ± 0.004

—————▶ richer

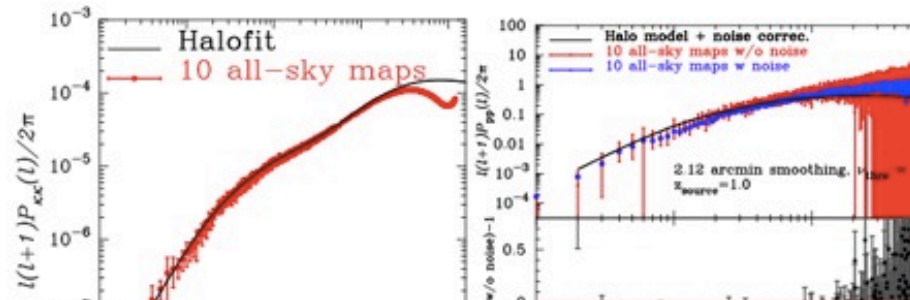
Probing Cosmology with Weak Lensing Selected Clusters

Masato Shirasaki (NAOJ)

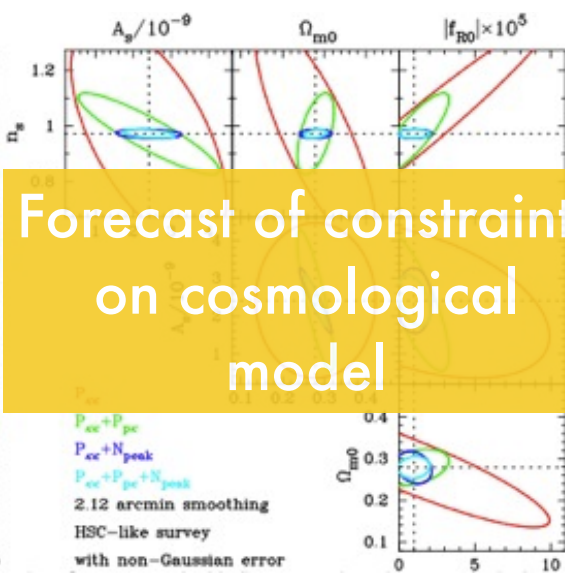


Full Sky Ray-tracing Simulation

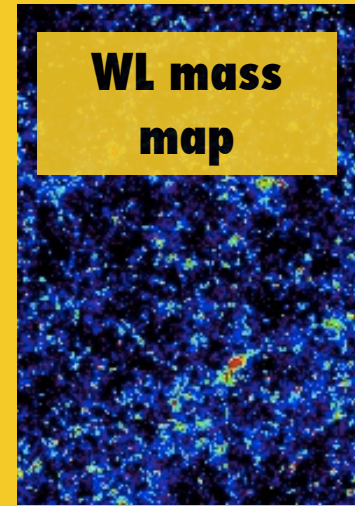
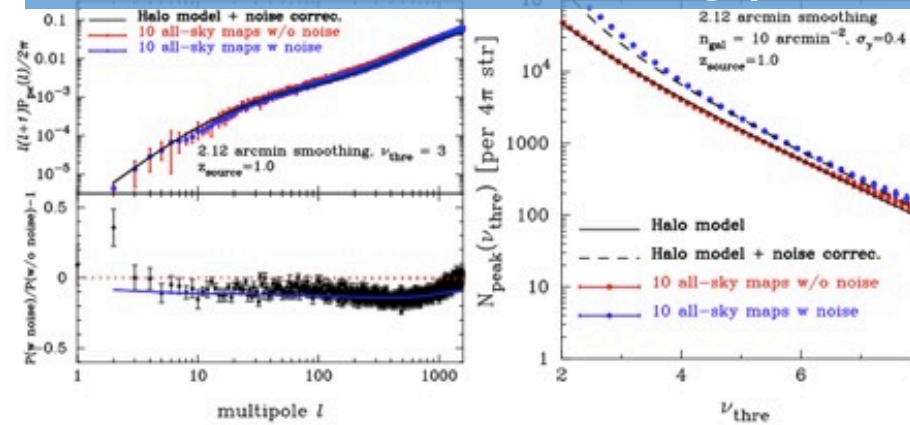
Weak lensing selected cluster = **peak** on WL mass map



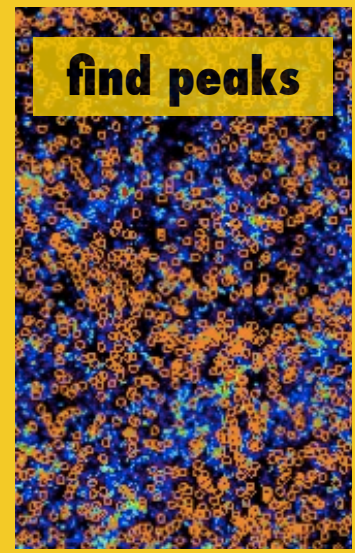
Detailed analyses of statistics for cosmic shear and lensing peaks



Forecast of constraints on cosmological model



WL mass map

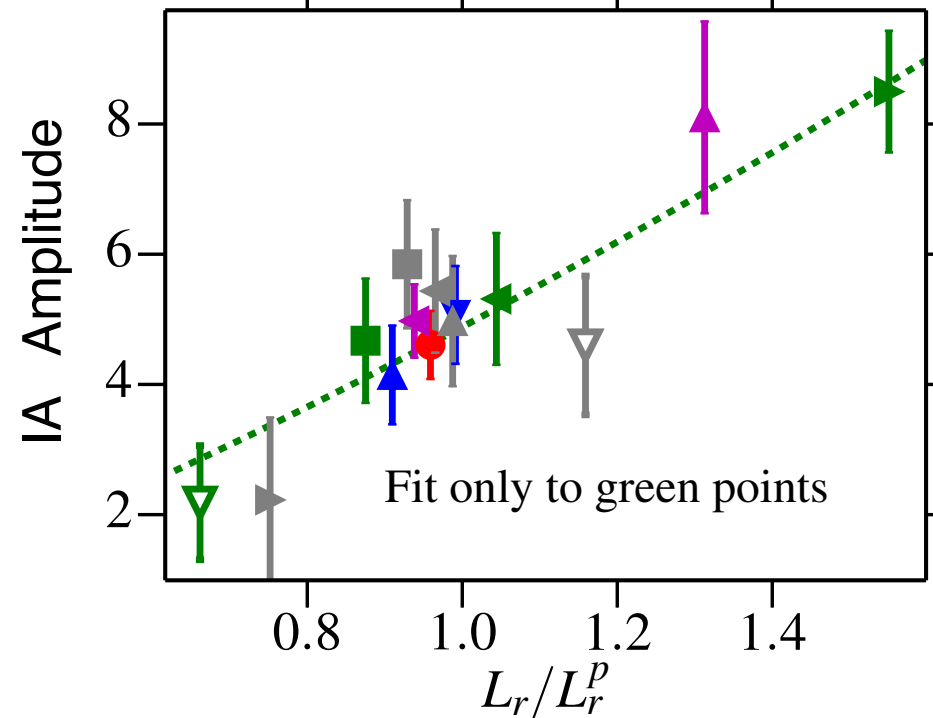
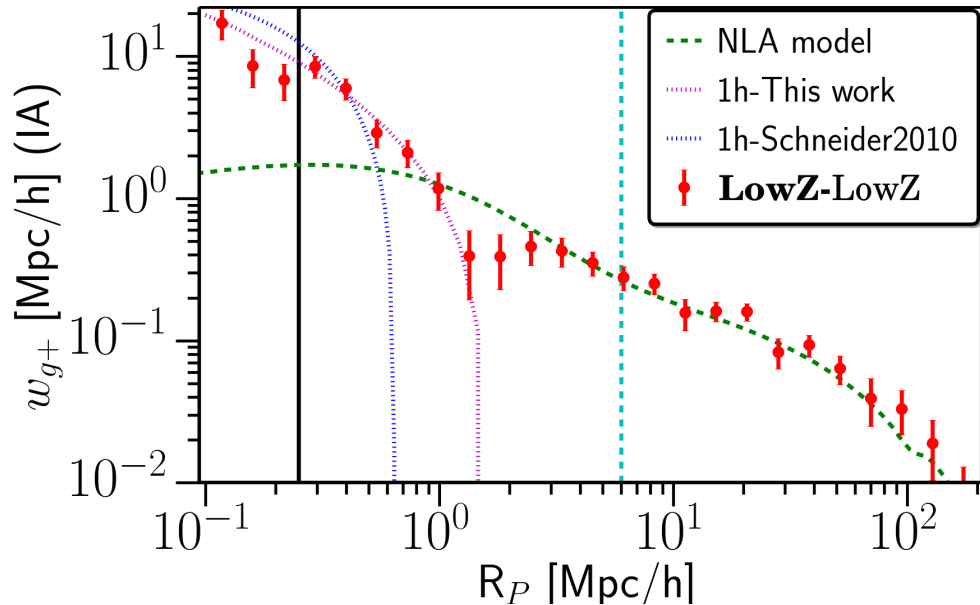


find peaks

Intrinsic Alignments of Galaxies

Sukhdeep Singh^{1*}, Rachel Mandelbaum¹

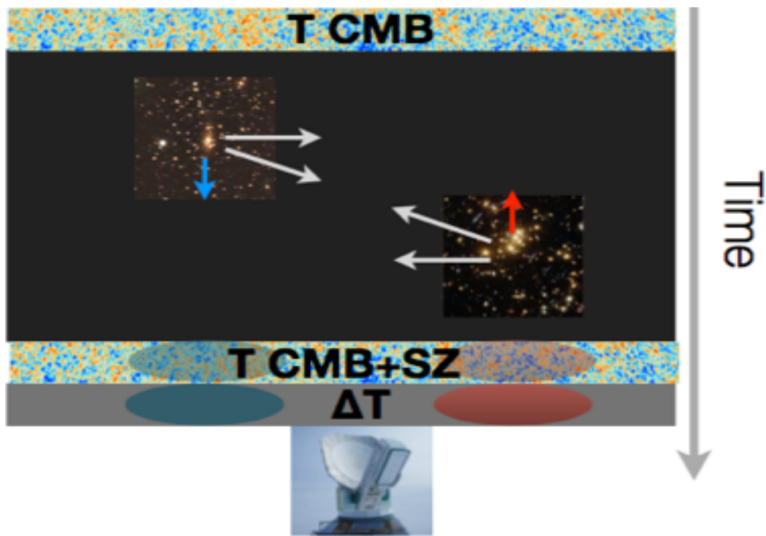
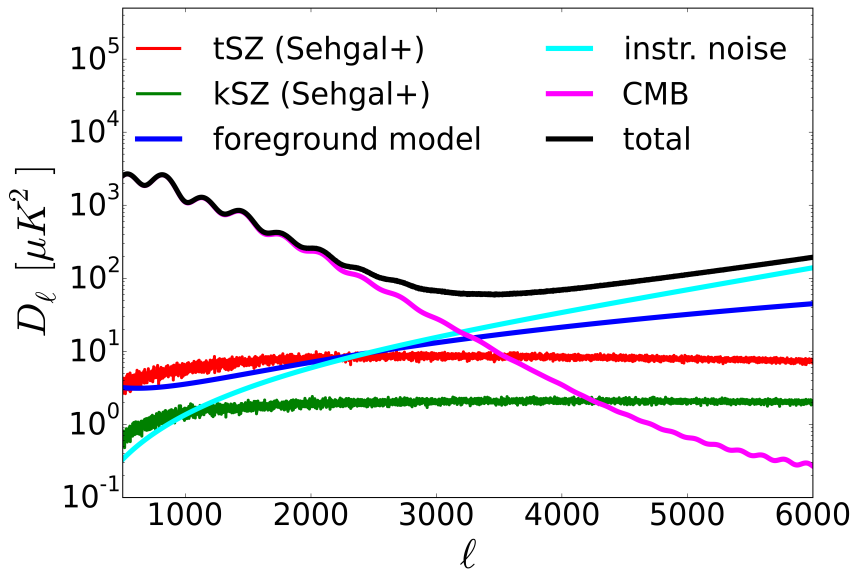
¹McWilliams Center for Cosmology, Carnegie Mellon University



- Intrinsic Alignments (IA) are contaminants in weak lensing studies.
- IA detected at high S/N in SDSS LOWZ galaxies. ([arxiv: 1411.1755](https://arxiv.org/abs/1411.1755))
- **Luminosity Dependence**: Brighter Galaxies have stronger IA
- **Environment Dependence**: Galaxies in more massive Halos show stronger alignments.

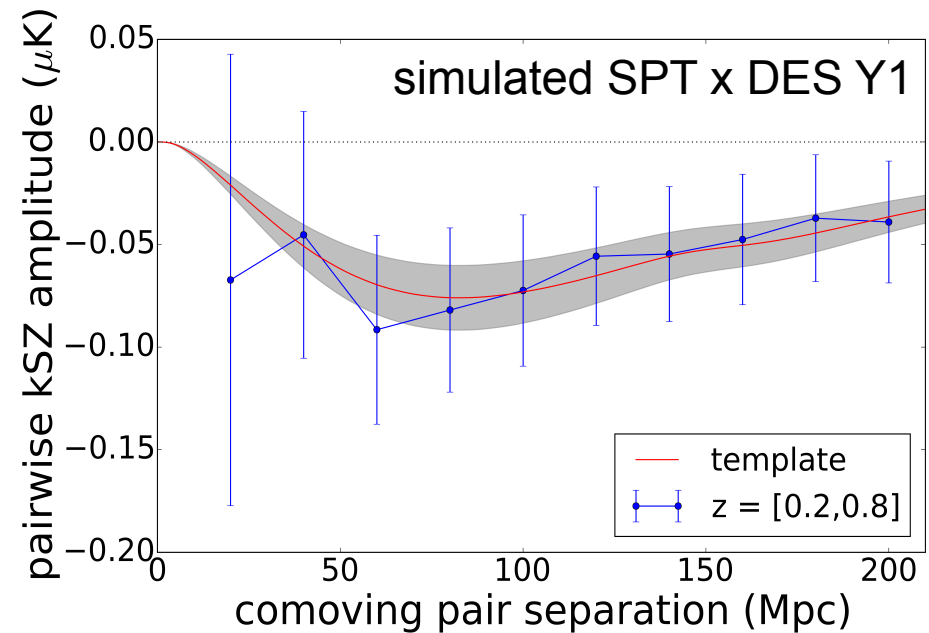
*sukhdeep@cmu.edu

Kinematic SZ from DES clusters & high-res. CMB data



Analytic template vs. simulations (Sehgal+):

- $>10\sigma$ with SPT x spec.-z cluster sample
- $\sim 5\sigma$ with SPT x DES Y1 (incl. photo-z errs.)



With real data:

- Testing on DESxACT: indication of a signal (1.7σ)
- now DESxSPT joint analysis, results soon (SPT contacts: K. Story, S. Flender)



Leibniz-Institut für
Astrophysik Potsdam

Unterstützt von / Supported by



Alexander von Humboldt
Stiftung/Foundation

Cosmicflows observations give us **CLUES** to matter distribution

Jenny Sorce

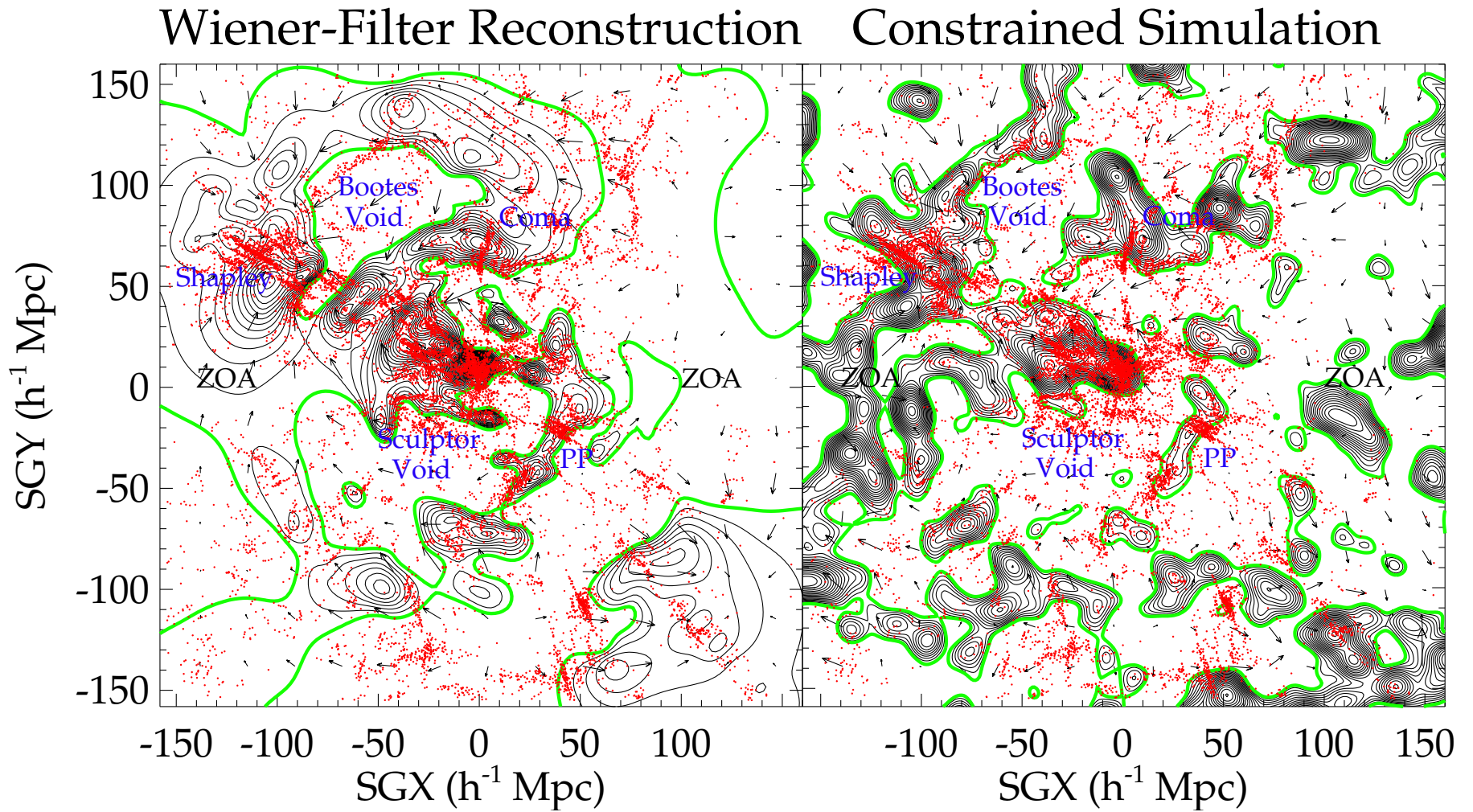
Theoretical and Observational Progress on Large-scale Structure of the Universe
MPA/ESO/MPE/Excellence Cluster Universe Joint Conference 2015

Garching, July 20th, 2015

AIP / Leibniz-Institut für Astrophysik Potsdam

Observed, Reconstructed and Simulated

Local Large Scale Structure



Analytical approach to density-weighted velocity correlation function of Large-Scape Structure

Naonori Sugiyama

(Collaborators: T.Okumura and D.N.Spergel)

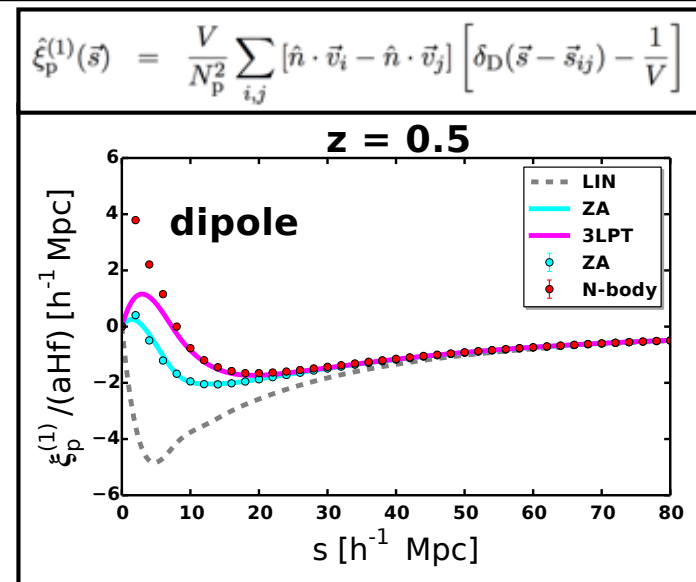
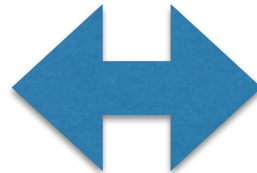
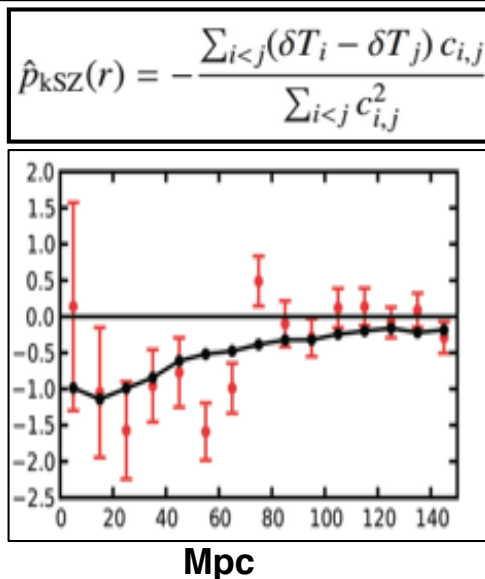
In anticipation of accurate measurement of kSZ effect in the future,
an analytical mode of pair-wise velocity correlation function is needed.

Simple relation between pair-wise velocity and density power spectra.

$$\hat{P}_p^{(n)}(\vec{k}, \hat{n}) = \left(i \frac{aHf}{\vec{k} \cdot \hat{n}} \right)^n \frac{\partial^n}{\partial^n f} \hat{P}_m(D, f, \vec{k}, \hat{n}).$$

$\tilde{P}_{\text{kSZ}} (\mu\text{K})$

Hand et al.
(2012)



An introduction to the post-Friedmann approach

The post-Friedmann approach [1] is based on a post-Newtonian type expansion in inverse powers of the speed of light, c , but designed for a Λ CDM FRW cosmology rather than an isolated object such as the solar system. The post-Friedmann formalism is based on a Poisson gauge metric (see below).

The leading order of the expansion is equivalent to taking the weak field, quasi-static, low velocity limit of the Einstein equations. Importantly, in this limit, the density contrast is not required to be small. Equivalently, this limit can be thought of as keeping the potentials small, but allowing their derivatives to be large.

The advantage of the expansion is that it allows equations beyond the exact limiting case to be derived, for example the equation for the vector potential in the metric.

Post-Friedmann metric

$$g_{00} = -\left(1 - \frac{2U_N}{c^2} + \frac{1}{c^4}(2U_N^2 - 4U_P)\right)$$

$$g_{0i} = -\frac{B_i}{c^3}$$

$$g_{ij} = a^2 \left[\left(1 + \frac{2V_N}{c^2} + \frac{2V_N^2 + 4V_P}{c^4}\right) \delta_{ij} + \frac{h_{ij}}{c^4} \right]$$

Leading Order Gravitational Equations

Using the post-Friedmann expansion, we can extract the leading order (c^{-2} and c^{-3}) gravitational equations. In addition to the standard Newton-Poisson equation, we get an equation setting the two scalar potentials to be equal, as expected in a Newtonian regime.

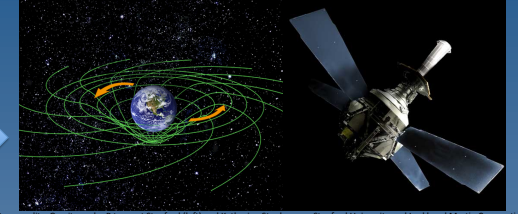
There is also an equation for the vector potential in the metric, which is sourced by the vector part of the momentum field. In the Newtonian limit, i.e. where the approximations hold exactly, this term vanishes, hence the requirement to use an expansion rather than just applying the approximations to the full Einstein equations.

Scalar equations

$$\frac{1}{c^2} \nabla^2 V_N = -\frac{4\pi G a^2 \rho_b}{c^2} \delta; \quad \frac{2}{c^2 a^2} \nabla^2 (V_N - U_N) = 0$$

Vector equation

$$\frac{1}{c^3} \nabla^2 B_i = -\frac{16\pi G a^2 \rho_b}{c^3} (1+\delta) v_i + \frac{2}{c^3} (\dot{a} U_{N,i} + a \dot{V}_{N,i})$$



Picture credits: Gravity probe B team at Stanford (left) and Katherine Stephenson, Stanford University and Lockheed Martin Corporation (right).

What is the vector potential?

The vector potential in the metric represents the ubiquitous relativistic effect of frame-dragging. A rotating massive body “drags” space-time around itself, causing test particles to spiral around the massive body. Gravity probe B [2] has measured this frame dragging effect for the Earth, finding it to be consistent with the value expected for General Relativity.

The non-linear relativistic universe

Dan B Thomas^{1,2}, Marco Bruni¹, David Wands¹
¹ICG, Portsmouth, UK; ²University of Cyprus, Nicosia, Cyprus

Why is Newtonian theory not enough?

Newtonian theory is widely used (typically using N-body simulations) to study non-linear scales in the universe. However,

- We would like to be able to test the approximations built into the Newtonian limit.
- There exist quantities in General Relativity that don't exist in Newtonian gravity, such as gravitational waves, the second scalar potential and the vector gravitational potential.

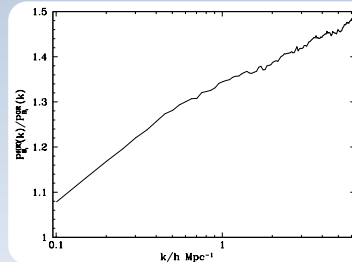
➔ *To completely characterise the universe on non-linear scales, we need to go beyond Newtonian theory*

f(R) gravity

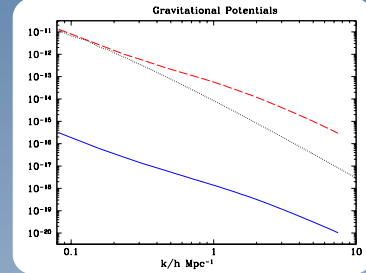
(in collaboration with Kazuya Koyama, Baojiu Li and Gong-bo Zhao)

The post-Friedmann formalism can be applied to other metric theories of gravity, not just General Relativity. We have applied it to Hu-Sawicki [8] f(R) gravity and derived the leading order gravitational equations. These include an equation for the vector potential, which we then extracted from f(R) N-body simulations, run using the ECOSMOG code [9].

As for GR, the small magnitude of the vector potential suggests that the approximations used to derive the equations used in the f(R) N-body simulations are valid, and also that future observations are only sensitive to the leading order scalar potentials.



The ratio of the vector power spectra in f(R) gravity and in General Relativity, as extracted from N-body simulations with the same initial conditions.



The power spectra of the scalar (red) and vector (blue) gravitational potentials, as extracted from N-body simulations. The grey curve is the linear theory scalar potential power spectrum for comparison.

Extraction from N-body simulations

In the leading order regime, the velocity and density fields obey the standard Newtonian equations of motion. Thus, by extracting the momentum field from N-body simulations (run using GADGET-2 [3] and analysed using DTFF [4]), the power spectrum of the vector potential can be calculated [5,6].

We found that the power spectrum of the vector potential is of order 10^5 times smaller than the power spectrum of the standard scalar gravitational potential. This small value suggests that

- The approximations inherent to the Newtonian regime are highly accurate
- Future observations (such as weak lensing, see over) are insensitive to the presence of the vector potential, even on non-linear scales

Conclusions

- We have examined the non-linear universe beyond Newtonian gravity, in particular focussing on the vector potential, which is the first correction to Newtonian gravity.
- We have extracted the vector potential from Λ CDM N-body simulations and found its power spectrum to be of order 10^5 times smaller than that of the scalar potential.
- This small value of the vector potential suggests that the approximations in the Newtonian regime are highly accurate.
- We have calculated the complete weak-lensing deflection angle on non-linear scales and shown that the additional relativistic contributions are small, thus justifying the standard ray-tracing approach.
- We have also examined f(R) gravity, again showing the vector potential to be small. This has similar consequences to the Λ CDM case.

Weak lensing

We have derived the complete weak lensing deflection angle up to c^{-4} order on fully non-linear scales [7]. This deflection angle is shown below, where the blue term on the first line is the standard weak lensing formula.

The blue terms on the third line comprise the Born correction and lens-lens coupling, which are the “first-order squared” terms that are included when ray tracing is performed. The extra terms show the additional relativistic terms that should be included at this order for consistency: the red terms are the contributions from the vector potential, the green term is from the tensor perturbations and the purple term is from the difference in the scalar potentials.

We have performed an order of magnitude estimate of the lensing affects from the vector potential, the lowest order correction to the standard ray tracing. We have shown that its contributions to both the E- and B-modes of cosmic shear are negligible for the next generation of weak-lensing surveys, thus justifying the standard ray-tracing approach.

$$\alpha^i = \int_0^{\chi_s} d\chi \left(\frac{\chi - \chi_s}{\chi_s} \right) \left[\begin{array}{l} \frac{1}{2a} (U_N + V_N)^i - \frac{1}{c^2} B_{\chi}^i \\ + \frac{2}{c^4} (U_P + V_P)^i - \frac{a}{c^4} \dot{B}^i + \frac{1}{2c^4} h_{\chi\chi}^i \\ + \frac{1}{c^4} (U_N + V_N)^i_j \int_0^{\chi} d\chi' (\chi - \chi') (U_N + V_N)^j \end{array} \right]$$

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Dusty Starbursts and AGNs in a $z=3.1$ Protocluster revealed by ALMA

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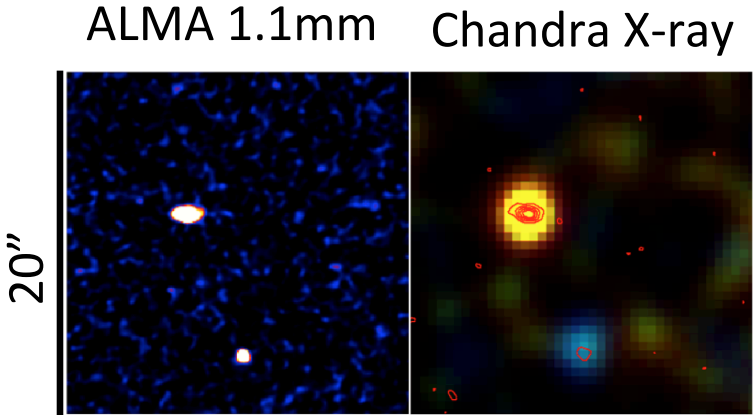
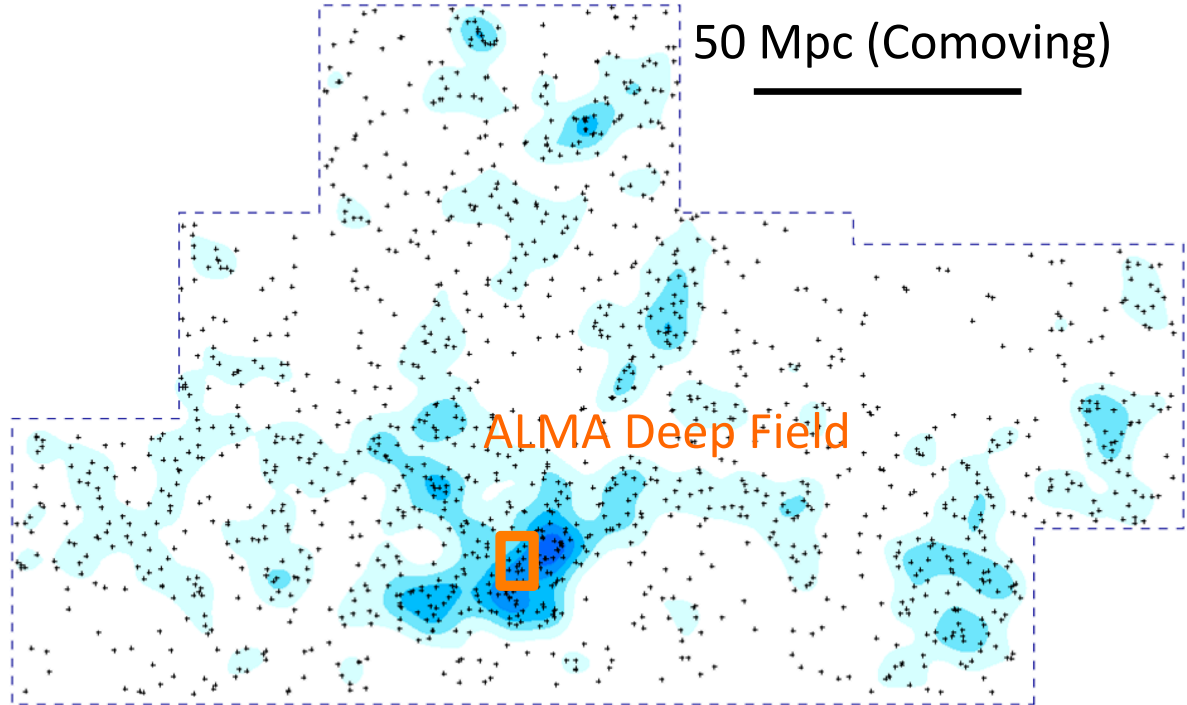
Theme: The most active populations and Cosmic Structure

Observation: ALMA Deep Field at the node of Cosmic Web at $z=3.09$

Results:

- Excess of SMGs (30 times)
- A number of multiple SMGs
- High AGN fraction (60%)

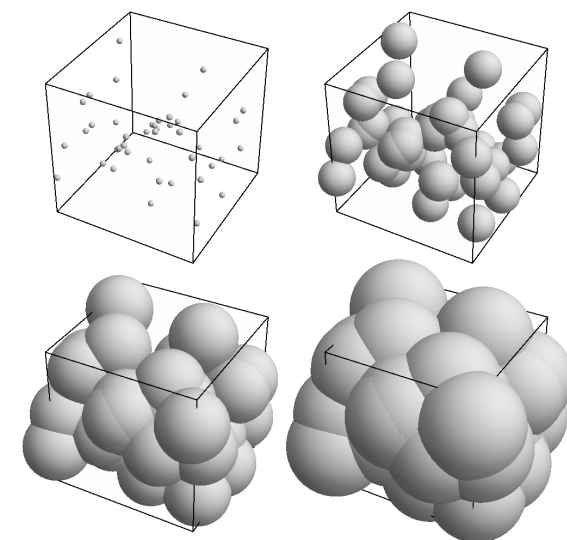
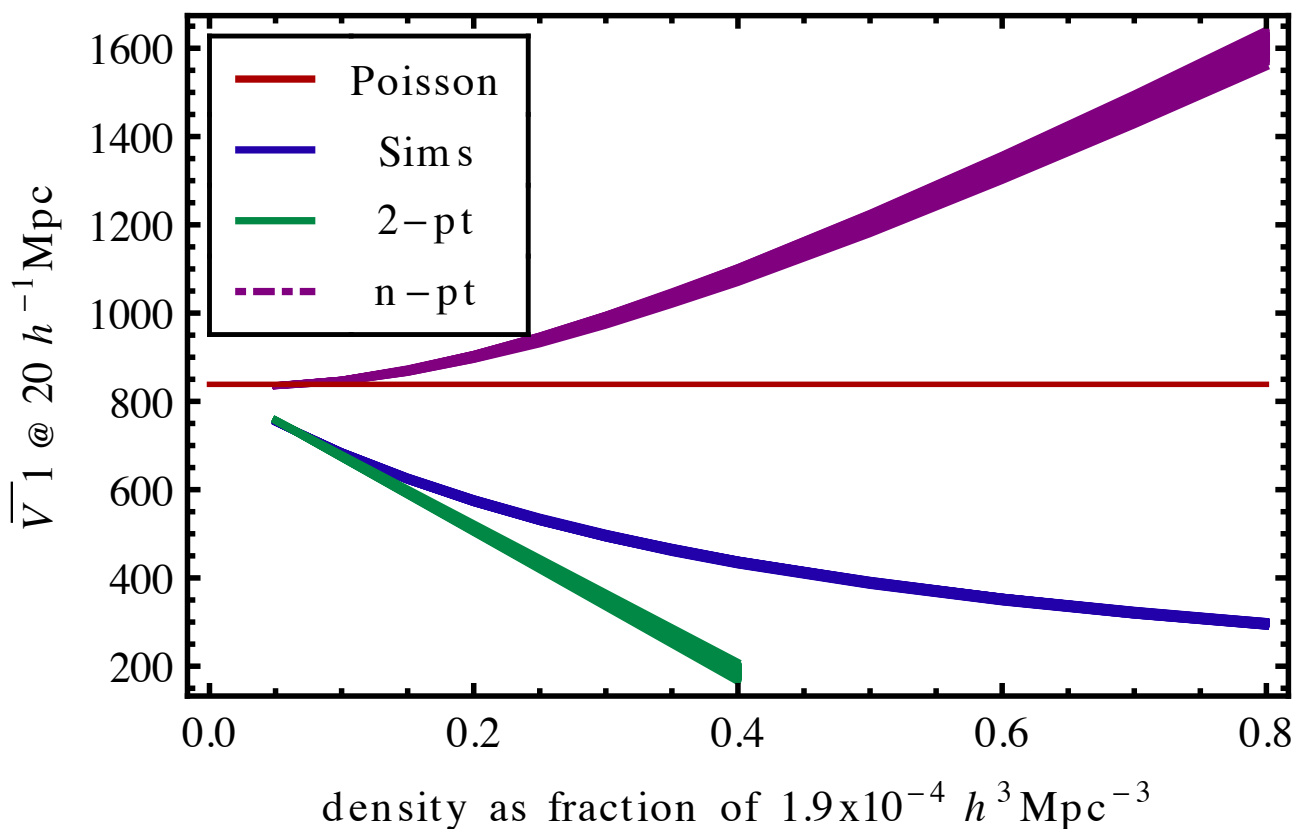
=> vigorous star-formation and SMBH growth at the node ($\sim 3\text{Mpc}$ scale).



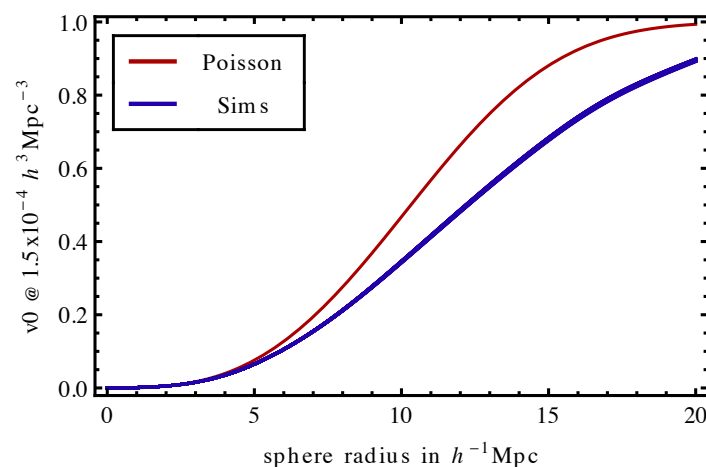
Example of AGN-host SMGs

$z=3.09$ Large Scale Structure traced by LAEs (Yamada+2012)

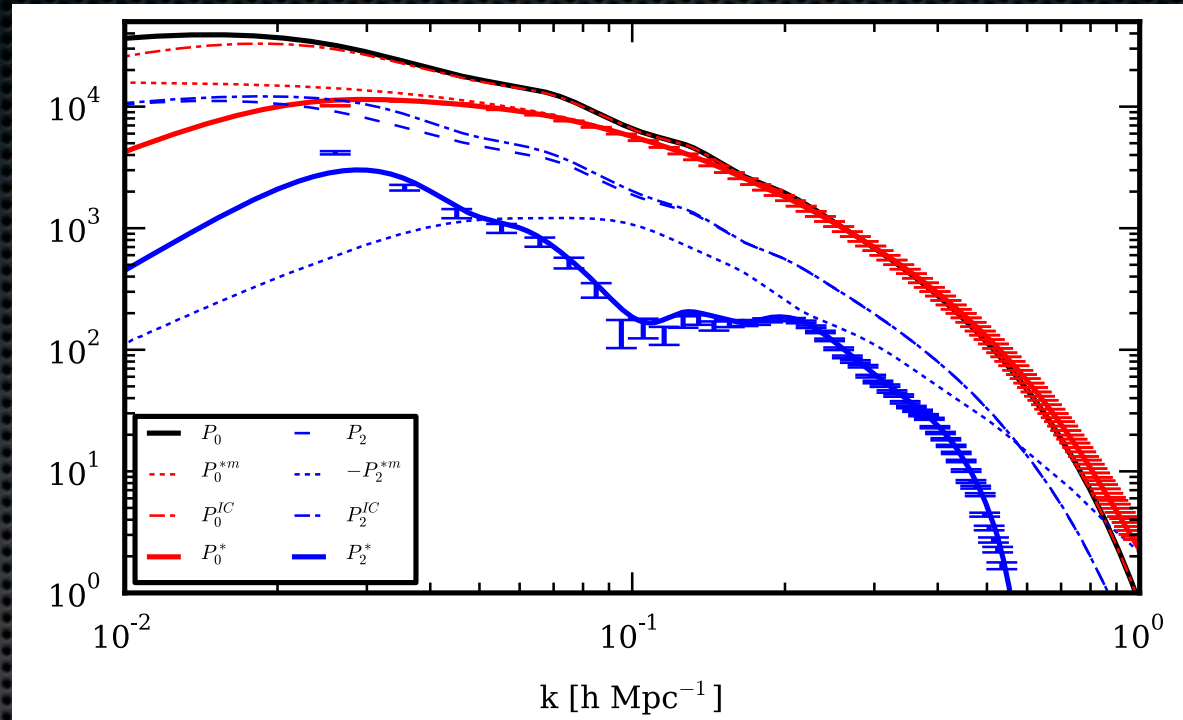
Minkowski functionals as a tracer of cosmic inhomogeneity



- Full n-point hierarchy ξ_{n+1} probed
- Highly significant detection of non-Gaussian correlations



Clipping the wings of non-linear structure

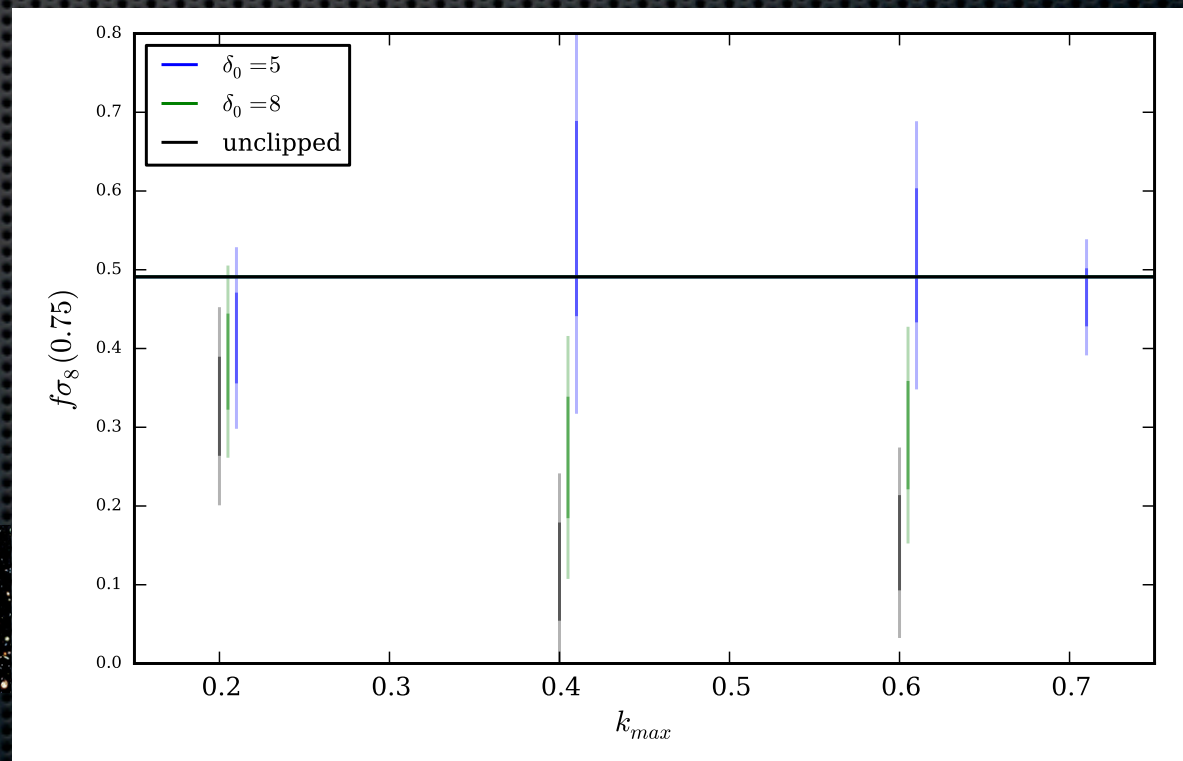


Michael Wilson

Prof. J. Peacock,
Dr S. de la Torre,
Prof. A. Taylor



VIMOS PUBLIC EXTRAGALACTIC REDSHIFT SURVEY

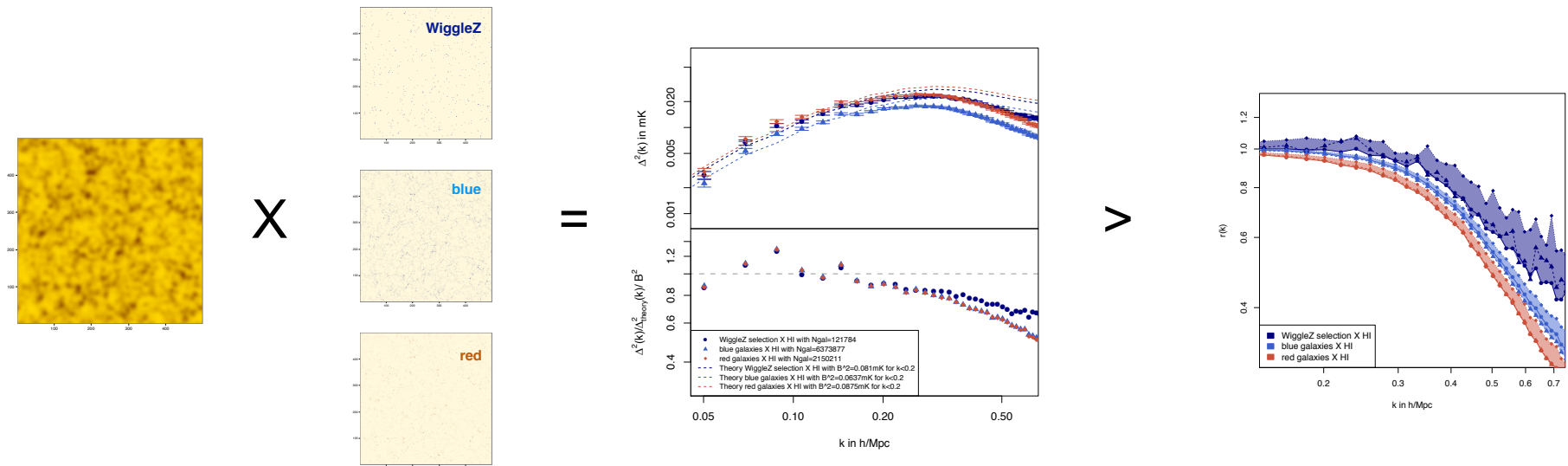


Intensity Mapping Cross-Correlations:

Connecting the Largest Scales to Galaxy Evolution

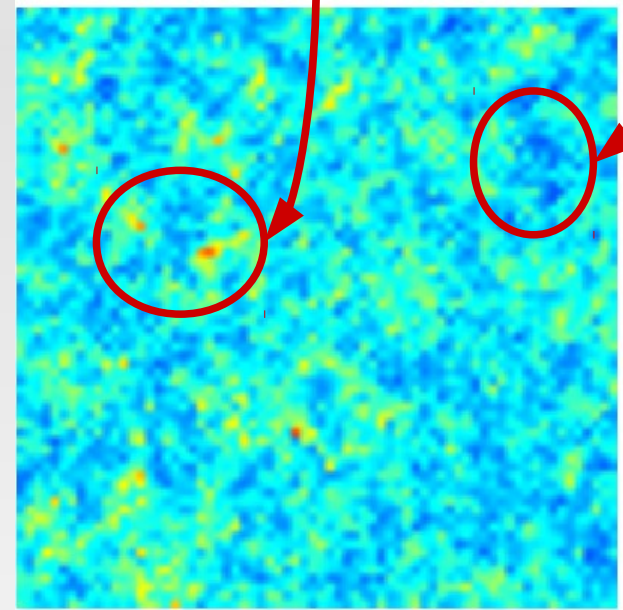
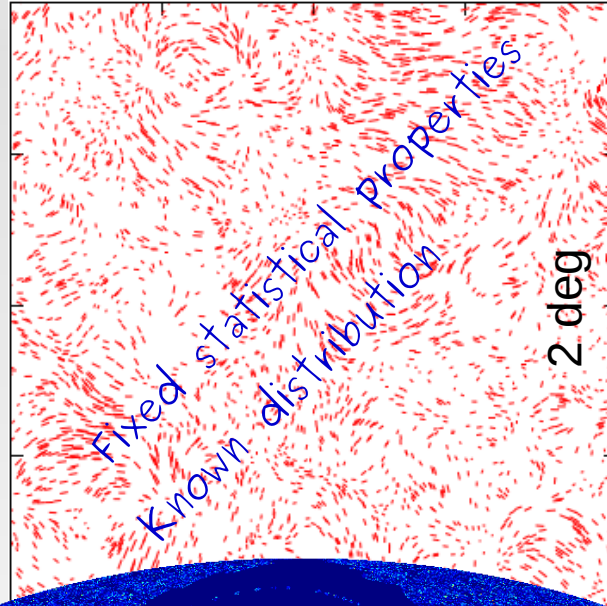
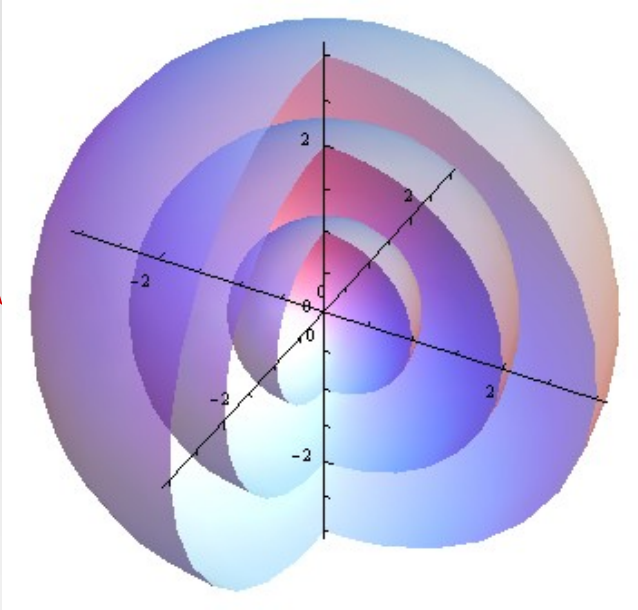
LAURA WOLZ
University of Melbourne

In Collaboration with Chiara Tonini, Chris Blake and Stuart Wyithe



Full-sky realisations of correlated lognormal fields

$$C_l^{\delta\kappa}(z, z') \delta_l^{l'} \delta_m^{m'} \equiv \langle \delta_{lm}(z) \kappa_{l'm'}^*(z') \rangle$$



Spherical symmetry

Redshift Space Distortions
Evolution
Selection effects
Wide area surveys

Applications

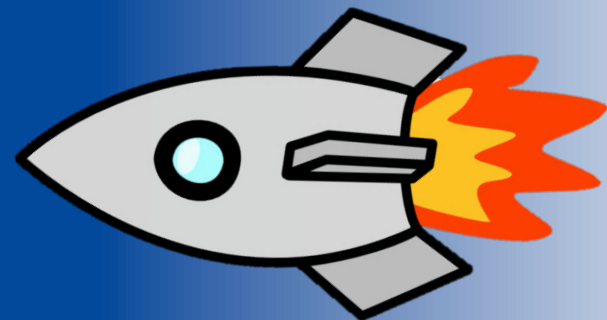
Cov. Matrix estimation
Tests for systematics
Estimator validation

Kinematic Dipole Detection with Galaxy Surveys

Mijin Yoon

Dragan Huterer

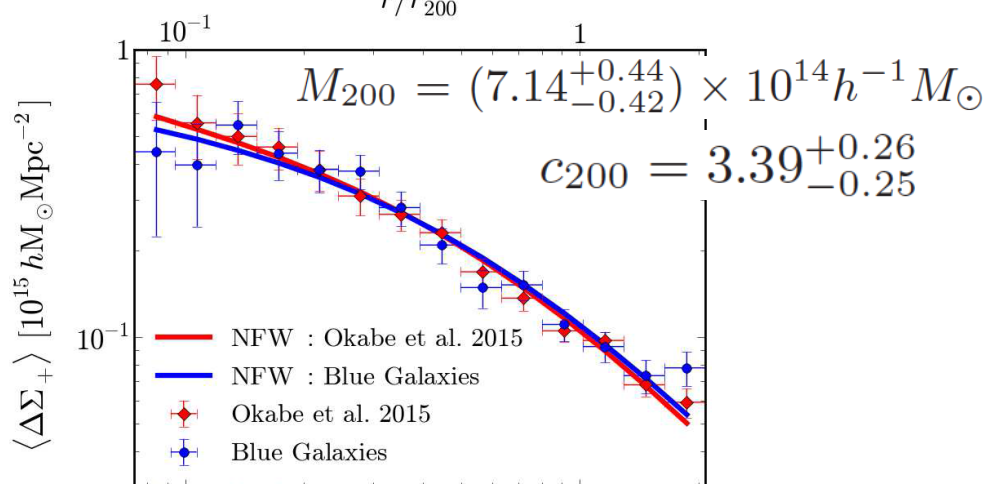
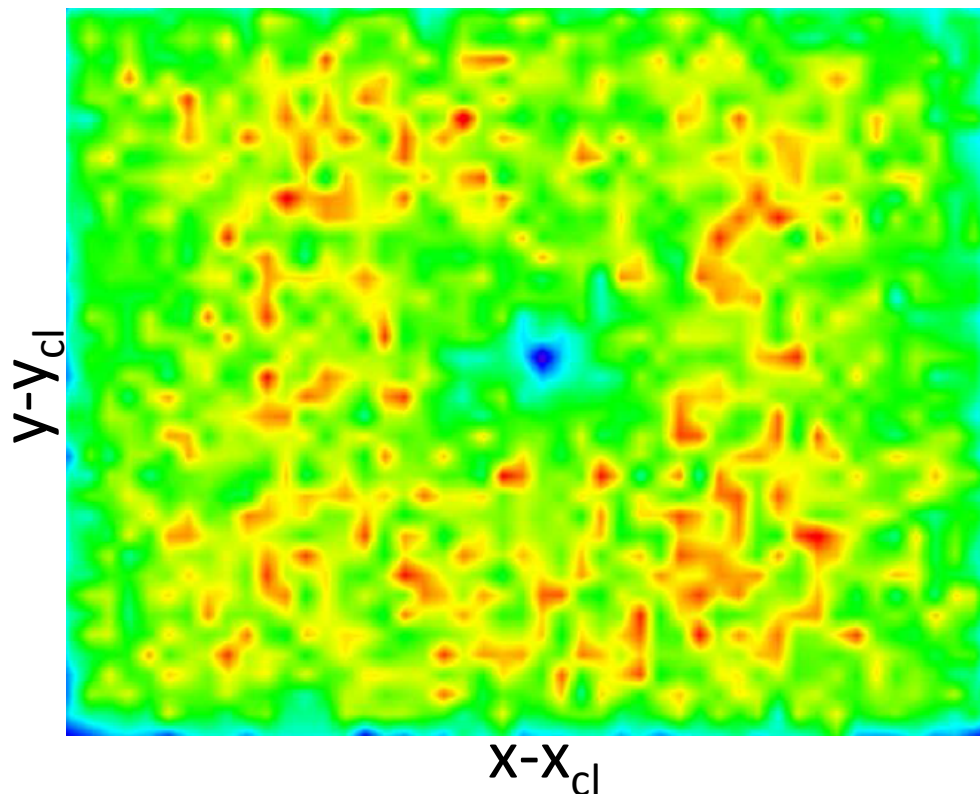
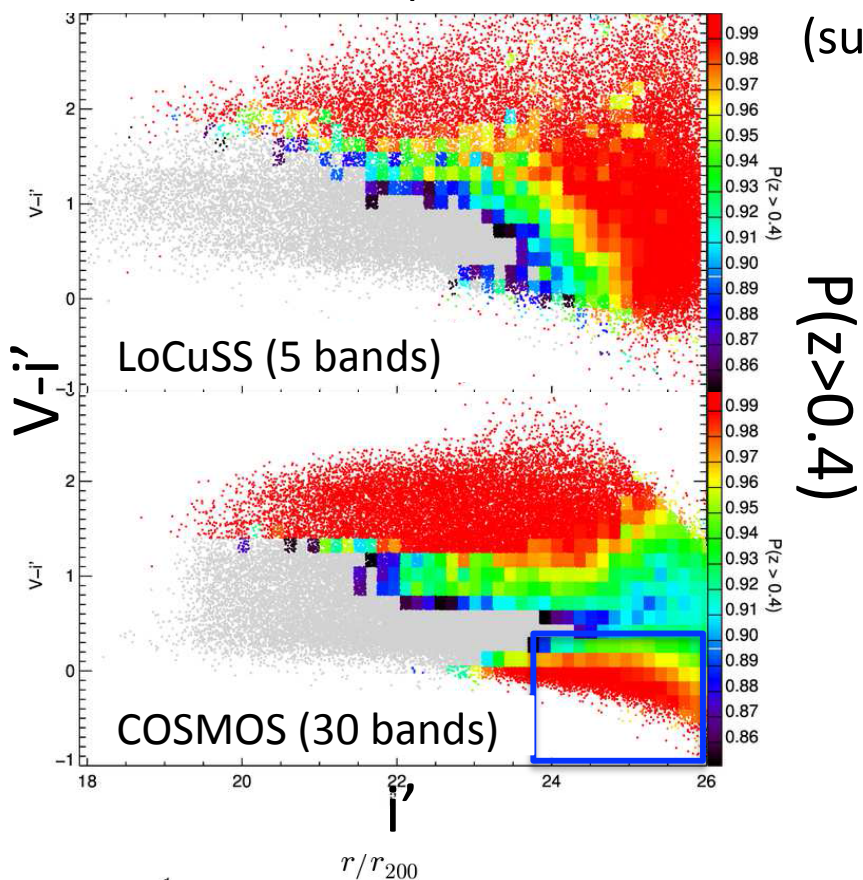
University of Michigan



LoCuSS: Exploring the selection of blue background galaxies for cluster weak-lensing

F. Ziparo, G. P. Smith, N. Okabe, C. P. Haines, M. J. Pereira, E. Egami

(submitted to MNRAS)



LoCuSS and COSMOS combined to select blue background galaxies. We obtain:

- Stacked mass and concentration consistent with red galaxy selection (1% contamination)
- Number density profile consistent with magnification lens bias (not flat)